

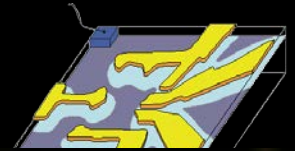
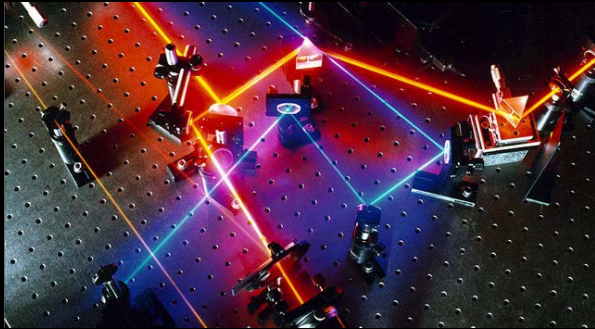
Quantum interactions with radiation that moves

Klaus Mølmer
Aarhus University
Denmark



Quantum Science Seminar





Quantum Optics

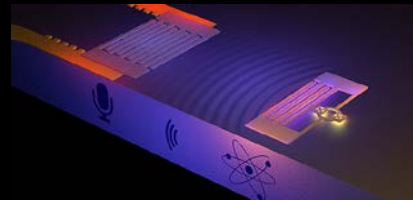
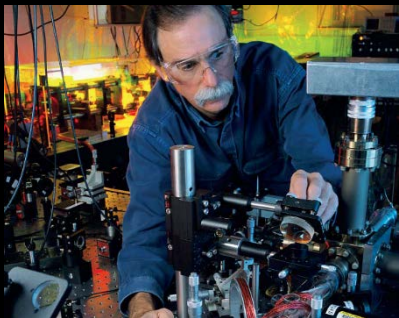
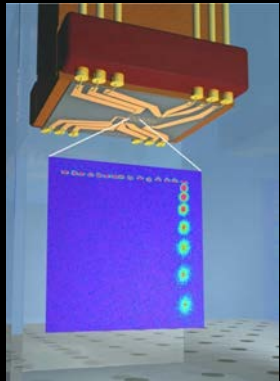
Atoms
Ions
Photons
Cavities
Travelling fields

...

Bits and Pieces

Quantum dots
Superconductors
Magnons
Cantilevers
Microwaves
Bulk and surface waves

...



The quantum theory of light

Maxwell's equations

$$\begin{aligned}\nabla \cdot \vec{B} &= 0 \\ \nabla \times \vec{E} + \partial \vec{B} / \partial t &= 0 \\ \nabla \cdot \vec{D} &= \rho \\ \nabla \times \vec{H} - \partial \vec{D} / \partial t &= \vec{J}\end{aligned}$$

electron coordinates are
quantum observables

Quantum mechanics is a "virus"

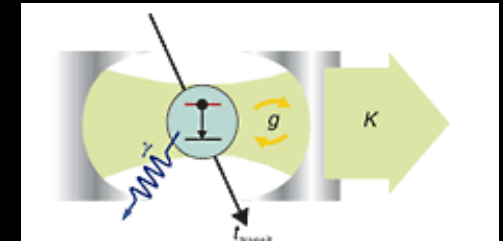
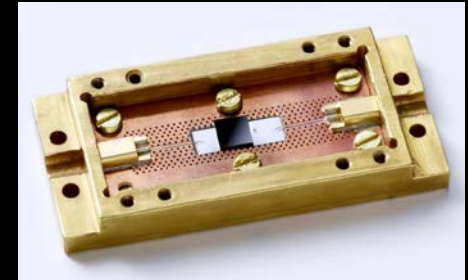
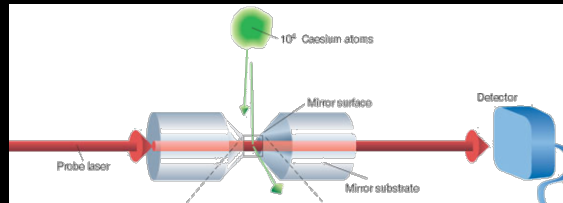
Fields \rightarrow quantum observables

Maxwell's Eqs \rightarrow Heisenberg Eqs of motion for field *operators*

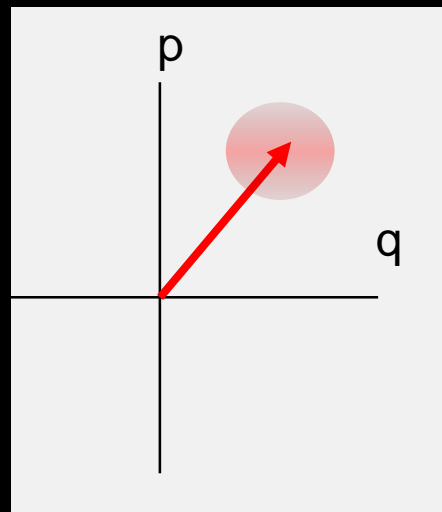
Quantum states of light

Annihilation and creation a, a^\dagger
number operator $n=a^\dagger a$,

Fock or number states, $|n\rangle$



Coherent state, $|\alpha\rangle$

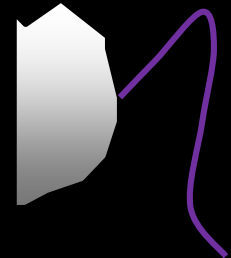
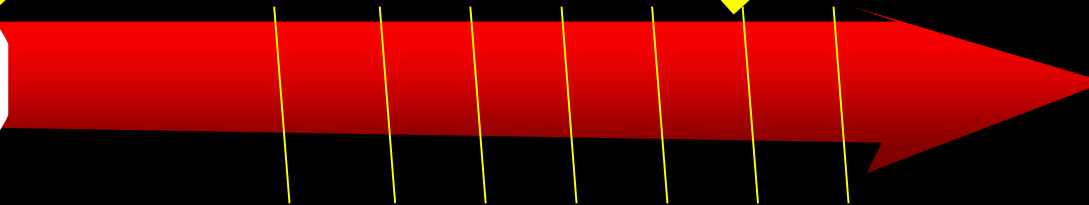
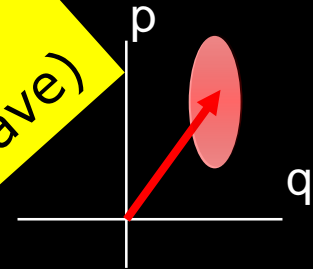


The small print in quantum optics textbooks

Schrödinger picture
Can I have my cake and eat it too ?



Single mode
(standing wave)



Schrödinger picture (expansion on number states)
is practically impossible.

Heisenberg picture (field observables)
yields mean values, correlation functions.

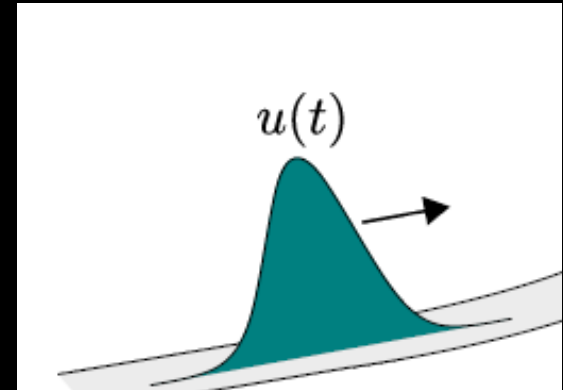
(Source) master equation: $\frac{d}{dt}\rho = \frac{1}{i\hbar} [H, \rho] - \frac{1}{2}(L^\dagger L \rho + \rho L^\dagger L) + L\rho L^\dagger$
Emitted field $\propto L = \sqrt{\gamma} \sigma$

The state of a *pulse* of light (microwave, SAW, ...)

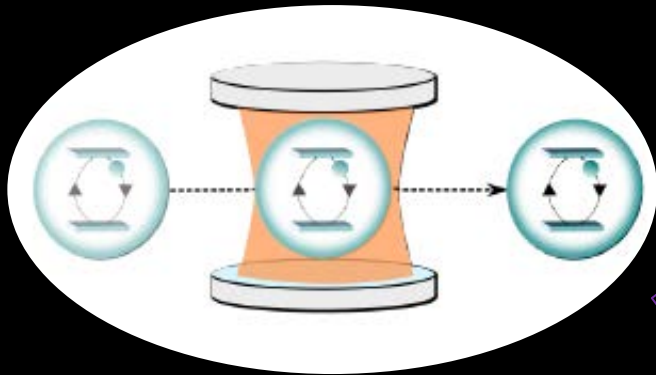
Wave packet:
solution of wave equation

Second quantization:
 $|n\rangle$ Fock state
or superposition state $\sum_n c_n |n\rangle$

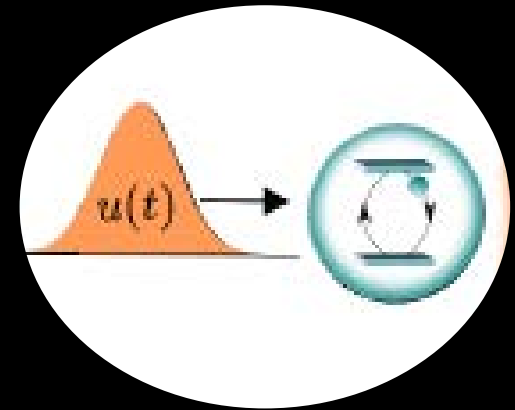
Such pulses may drive quantum systems,
may work as flying qubits, may probe quantum systems,
may transport pure or mixed states, transport energy ...



How does a quantum pulse interact with a qubit ?



Flying atom, fixed mode:
coupling $g \rightarrow g(t)$



Flying mode, fixed atom:
coupling $g \rightarrow u(t)$?

The same ?

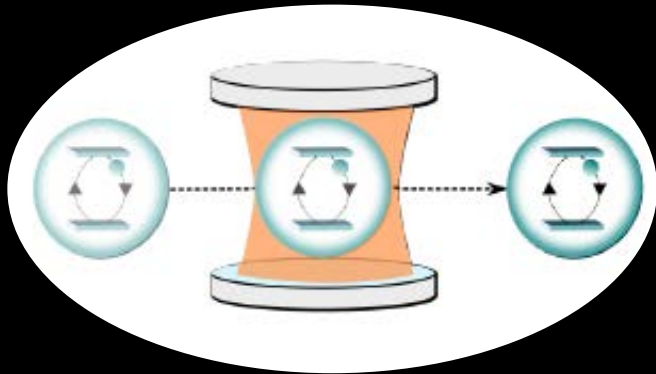
NO !

→ Exchange of quanta between emitter and field

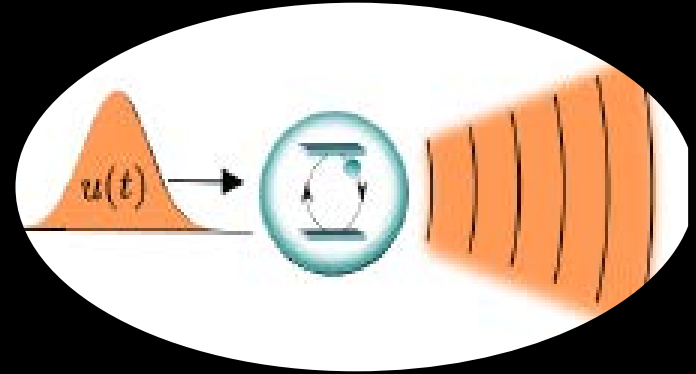
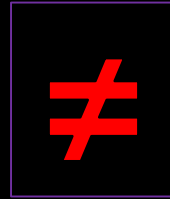
→ Distortion of the pulse (mode continuum)

mix of the two: genuine multi-mode theory

How does a quantum pulse interact with a qubit ?



Flying atom, fixed mode:
coupling $g \rightarrow g(t)$



One atom & mode **continuum**
Open quantum system

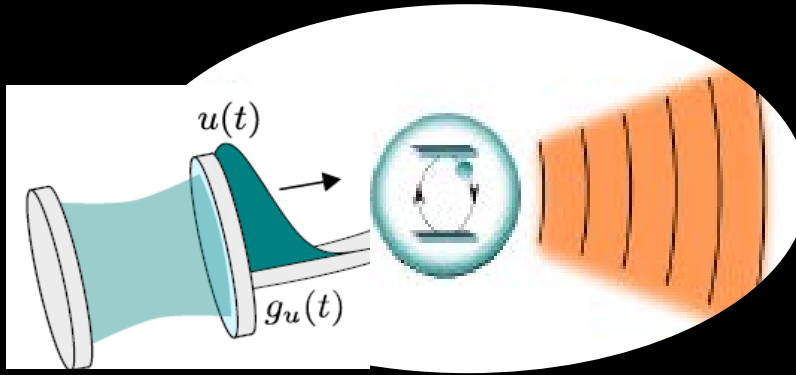
Cascaded system master equation
(Gardiner 1993, Carmichael 1993)

See also:

B. Q. Baragiola, et al (J. Combes),
“n-photon wave packets interacting
with an arbitrary quantum system,”
Phys. Rev. A 86, 013811 (2012).

How does a quantum pulse interact with a qubit ?

A trick !



$$g_u(t) = \frac{u^*(t)}{\sqrt{1 - \int_0^t dt' |u(t')|^2}}$$

Single-mode cavity and an atom

Jaynes-Cummings Hamiltonian:

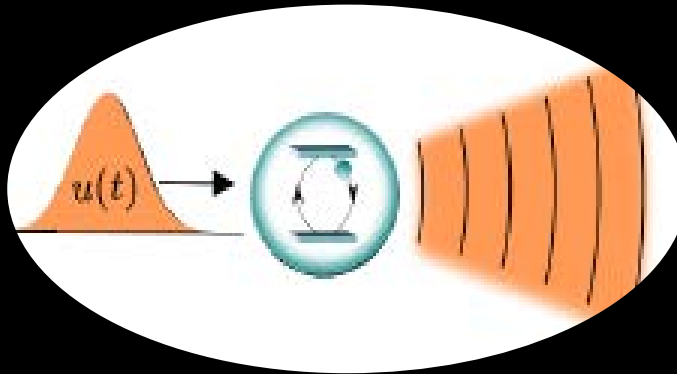
$$H = \frac{i\sqrt{\gamma}}{2} (g_u^*(t)a_u^+ \sigma - g_u(t)a_u \sigma^+)$$

Damping (Lindblad) operator:

$$L = g_u(t)a_u + \sqrt{\gamma} \sigma$$

Alexander Holm Kiilerich and Klaus Mølmer
Input-Output Theory with Quantum Pulses
Phys. Rev. Lett. **123**, 123604 (2019).

How does a quantum pulse interact with a qubit ?



$$H = \frac{i\sqrt{\gamma}}{2} (g_u^*(t) a_u^\dagger \sigma - g_u(t) a_u \sigma^+)$$

$$L = g_u(t) a_u + \sqrt{\gamma} \sigma$$

Master equation:

$$\begin{aligned} \frac{d}{dt} \rho &= \frac{1}{i\hbar} [H, \rho] - \frac{1}{2} (L^\dagger L \rho + \rho L^\dagger L) + L \rho L^\dagger \\ &= \sqrt{\gamma} \{ g_u(t) (a_u \rho \sigma^+ - a_u \sigma^+ \rho) + g_u^*(t) (\sigma \rho a_u^\dagger - \rho a_u^\dagger \sigma) \} \\ &\quad + D[\sqrt{\gamma} \sigma] \rho + D[g_u(t) a_u] \rho \end{aligned}$$

Chiral "Hamiltonian"
Excitations: " \rightarrow "

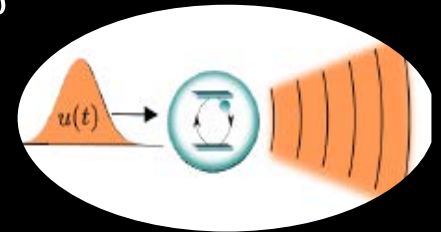
$\rho_F \rightarrow \text{vacuum state}$

Alexander Holm Kiilerich and Klaus Mølmer
Input-Output Theory with Quantum Pulses
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How does a quantum pulse interact with a qubit ?

Field + Atom Master Equation:

$$\begin{aligned} \frac{d}{dt}\rho = & \sqrt{\gamma} \{ g_u(t) (a_u \rho \sigma^+ - a_u \sigma^+ \rho) + g_u^*(t) (\sigma \rho a_u^\dagger - \rho a_u^\dagger \sigma) \} \\ & + D[\sqrt{\gamma} \sigma] \rho + D[g_u(t) a_u] \rho \end{aligned}$$



If input "cavity" field is in a *coherent state*: $|\alpha\rangle\langle\alpha| \rightarrow |\alpha(t)\rangle\langle\alpha(t)|$

Atom Master Equation :

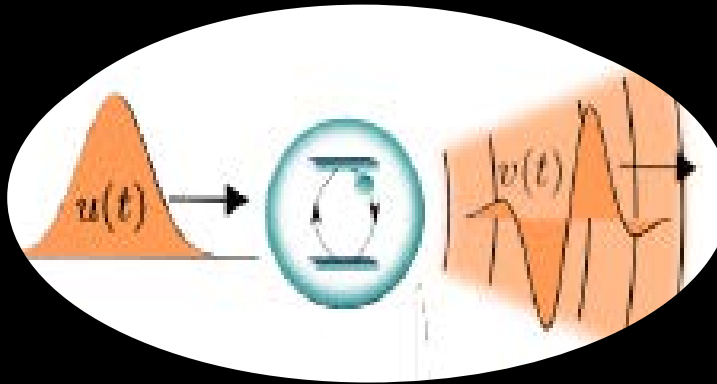
$\rho_F \rightarrow \text{vacuum state}$

$$\frac{d}{dt}\rho = \underbrace{\sqrt{\gamma} [u(t)\alpha^*(0) \sigma - u^*(t)\alpha(0)\sigma^+]}_{\text{classical drive}} \rho_A + \underbrace{D[\sqrt{\gamma} \sigma] \rho_A}_{\text{atomic decay}}$$

Input Fock state is more difficult: solve $\rho_{FA}(t)$

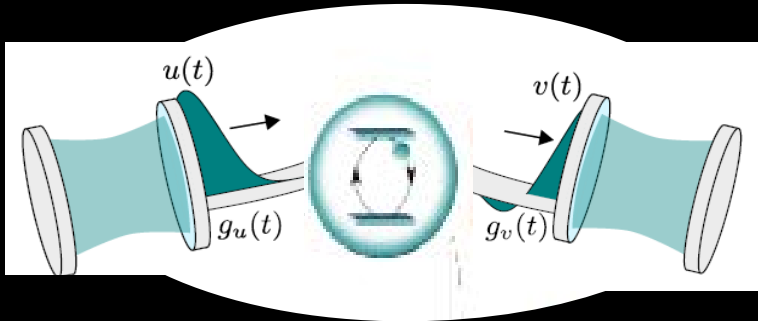
Opposite to Jaynes-Cummings model: Fock state easy.

How does a quantum pulse interact with a qubit ?



What about the state of the pulse after the interaction ?

$\rho_F \rightarrow \text{vacuum state}$



State contents of pulse $v(t)$ after the interaction ?

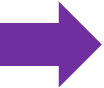


Cascaded Master Equation for $u(t)$ -cavity + qubit + $v(t)$ -cavity

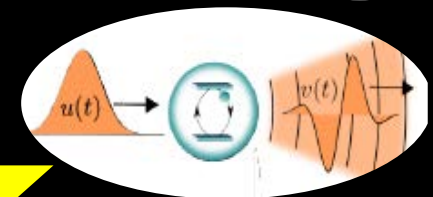
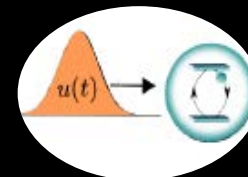
$$g_v(t) = - \frac{v^*(t)}{\sqrt{\int_0^t dt' |v(t')|^2}}$$

More general "scatterer": $H_s \{L_i\}$

How does a quantum pulse interact with a qubit ?

More general "scatterer": $H_s \{L_i\}$

	Quantum input pulse (u)	Quantum scatterer (s)	Quantum output pulse (v)
	\times	\times	$\circ \quad \rho_{us}$
	\circ	\times	$\times \quad \rho_{sv}$
	\times	\times	$\times \quad \rho_{usv}$



$$H = -\frac{i}{2} \left(g_u(t) \sqrt{\gamma} a_u \sigma^+ + g_u(t) g_v^*(t) a_u a_v^+ + \sqrt{\gamma} g_v(t) \sigma a_v^+ - h.c. \right)$$

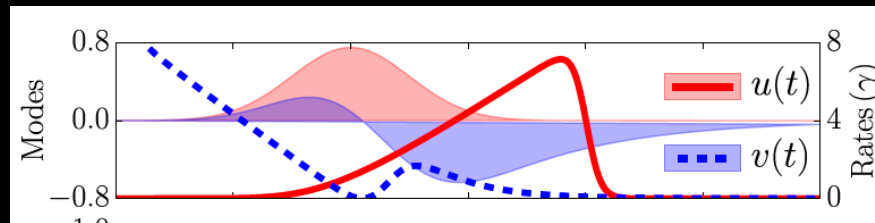
$$L = g_u(t) a_u + \sqrt{\gamma} \sigma + g_v(t) a_v$$

Chiral "Hamiltonian"
Excitations: " \rightarrow "

Examples

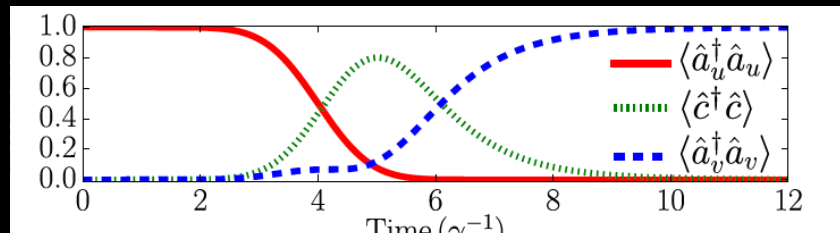
Single photon scattering on an empty cavity ($\sigma \rightarrow c$).

Input wave packet $u(t)$, $g_u(t)$ \rightarrow output wave packet $v(t)$, $g_v(t)$



Result:

Occupation of input, cavity and output:



Examples

Scattering of a coherent state scattering on an empty cavity ($\sigma \rightarrow c$).

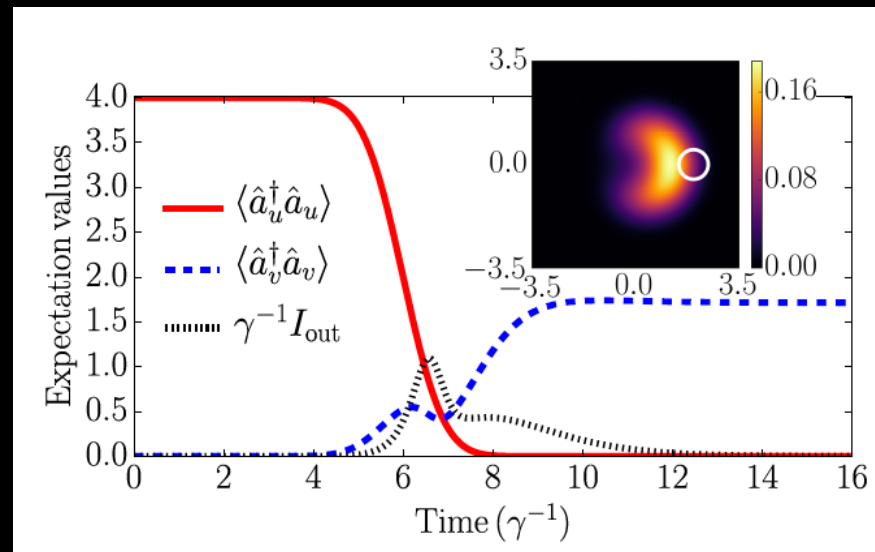
Input wave packet $u(t)$, $g_u(t)$

Cavity with phase noise (shaking mirror) \rightarrow Output is multi-mode

Dominant output mode

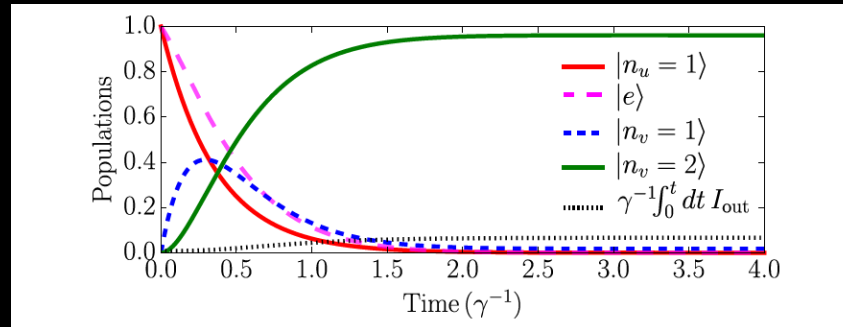
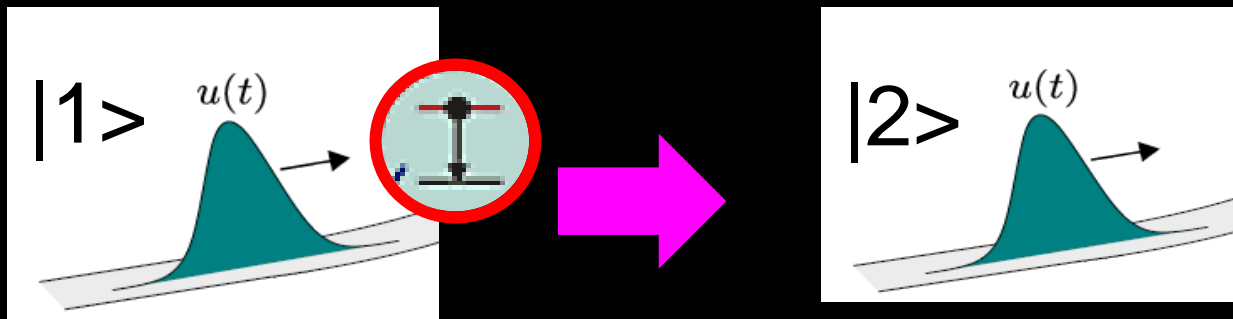
$|\alpha=2\rangle$ coherent input state

\rightarrow dephased and damped $W(q,p)$



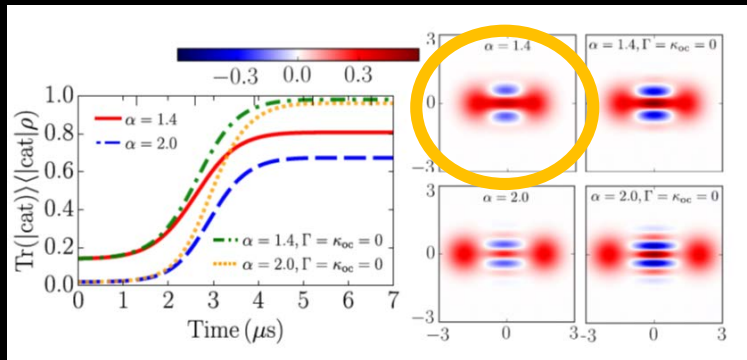
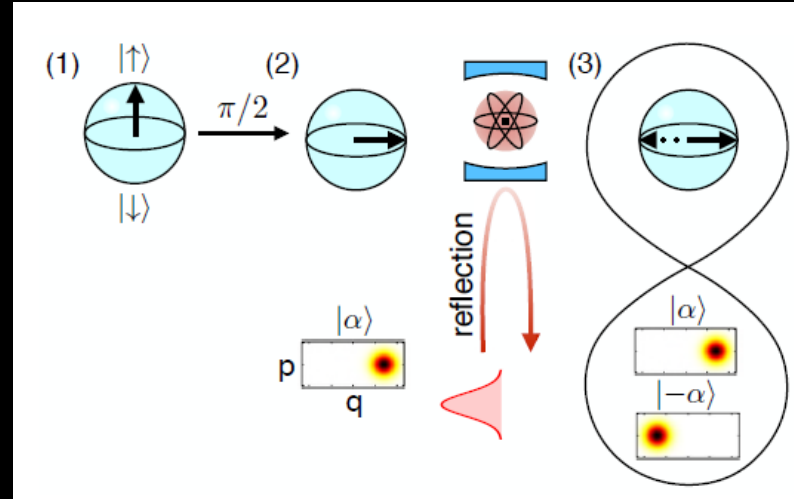
Examples

Stimulated emission (same mode)



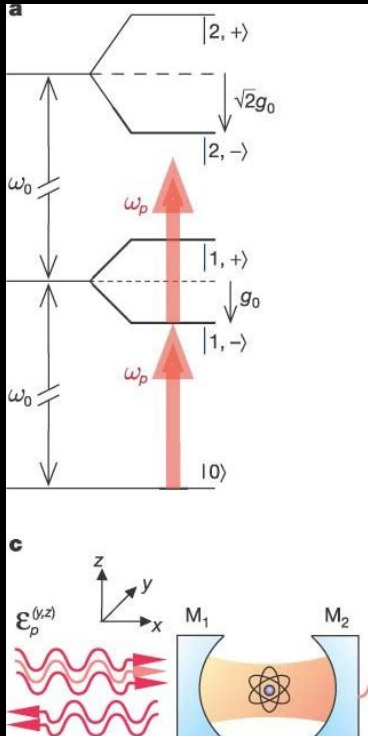
Examples

Schrödinger's cat



Hacker et al (Rempe group),
Nature Photonics 2019.

The "photon bandwidth dilemma"

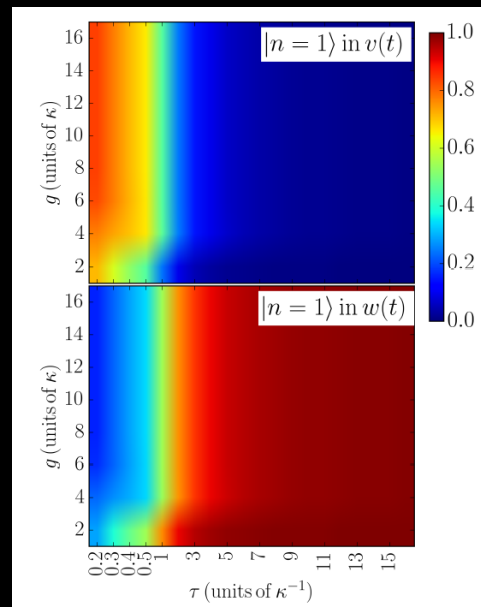


Photon blockade in cavity QED

Wave packet incident on cavity with a single atom may be fully transmitted for 1 photon (resonant with eigenstate) and reflected for 2 or more photons (non resonant).

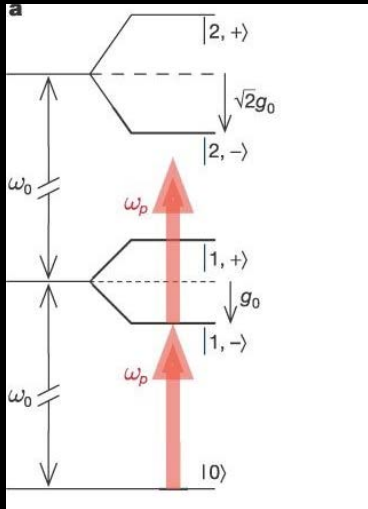
One photon,
reflected

One photon,
transmitted



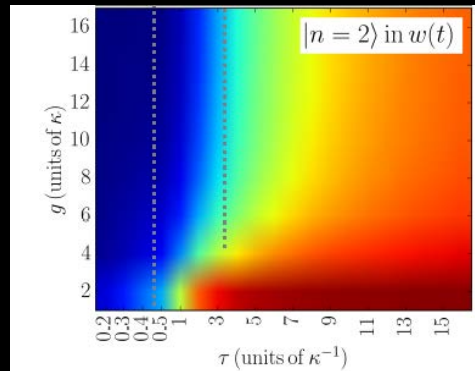
[arXiv: 2003.04573](https://arxiv.org/abs/2003.04573)
Quantum interactions
with pulses of radiation
[A. Kiilerich](#), KM

The "photon bandwidth dilemma"



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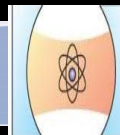
But, two photons in a pulse are also transmitted ???

Short pulse =
broad band:
always reflects

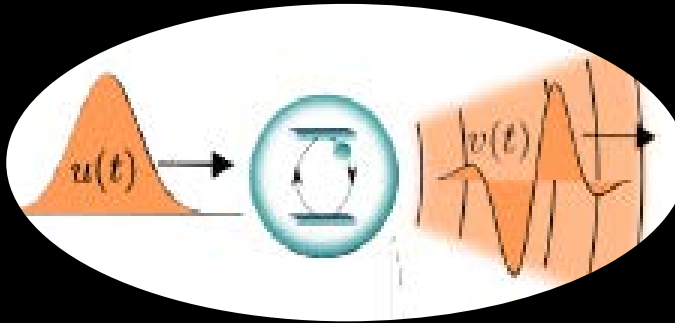


Narrow band = long pulse
small photon overlap
sequential transmission !

[arXiv:2003.04573](https://arxiv.org/abs/2003.04573)
Quantum interactions
with pulses of radiation
[A. Kiilerich](#), KM



What is the output mode $v(t)$?

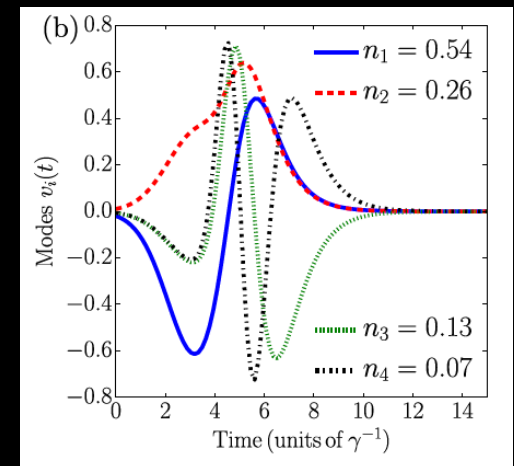
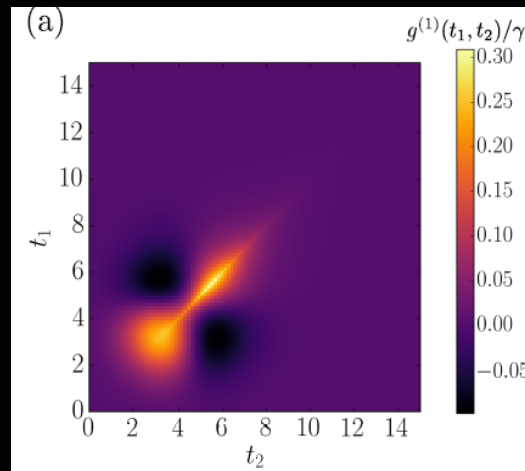


Output field:

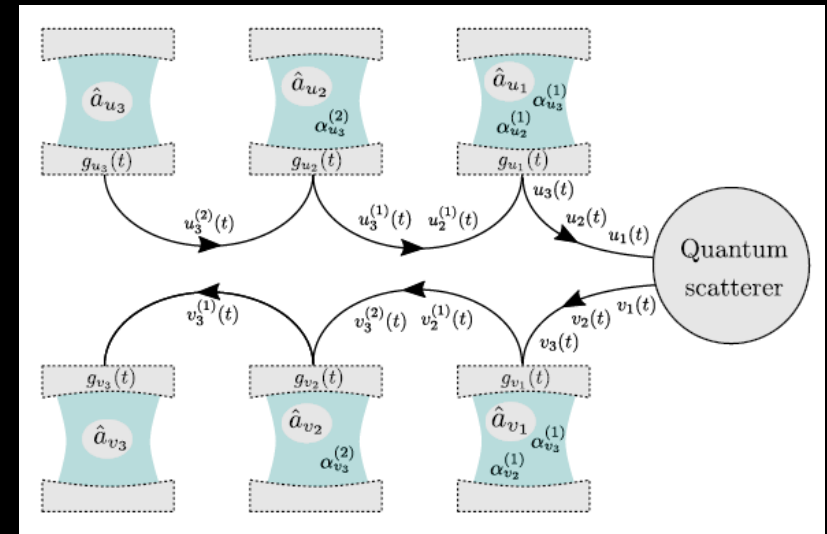
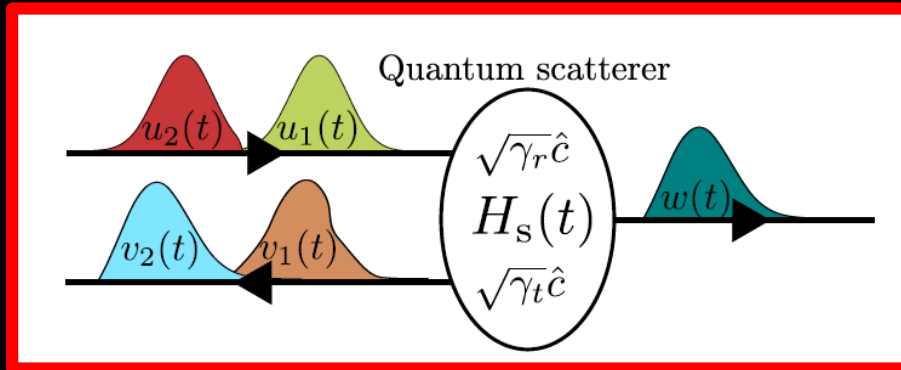
$$L = g_u(t)a_u + \sqrt{\gamma} \sigma$$

$$g^1(t, t') = \langle L^+(t)L(t') \rangle \\ = \sum_i n_i v_i^*(t)v_i(t')$$

Coherent state on cavity with phase noise (shaking mirror)



Multiple input and output modes



[arXiv:2003.04573](https://arxiv.org/abs/2003.04573)

Quantum interactions with pulses of radiation

Alexander Holm Kiilerich, Klaus Mølmer

Conclusion

Sometimes photons move, and that is what we like about them ; =)

Moving photons occupy a continuum of modes, and their quantum states and dynamics are non-trivial.

Quantum information protocols rely on precise handling of the modes.

If we may restrict to few incident and outgoing travelling modes (solutions of classical wave equation), we can apply usual master equation theory.

Theory applies to any wave (but assumes Markov approximation and linear dispersion)