

# Quantum computing over the rainbow: the quantum optical frequency comb as a platform for measurement-based universal quantum computing

Olivier Pfister  
University of Virginia



Quantum Fields and Quantum Information website:  
<https://sites.google.com/view/qfqi>

# Quantum computing over the rainbow: the quantum optical frequency comb as a platform for measurement-based universal quantum computing

OR

## How I learned to stop worrying and love the comb

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QSS#13 07/16/2020

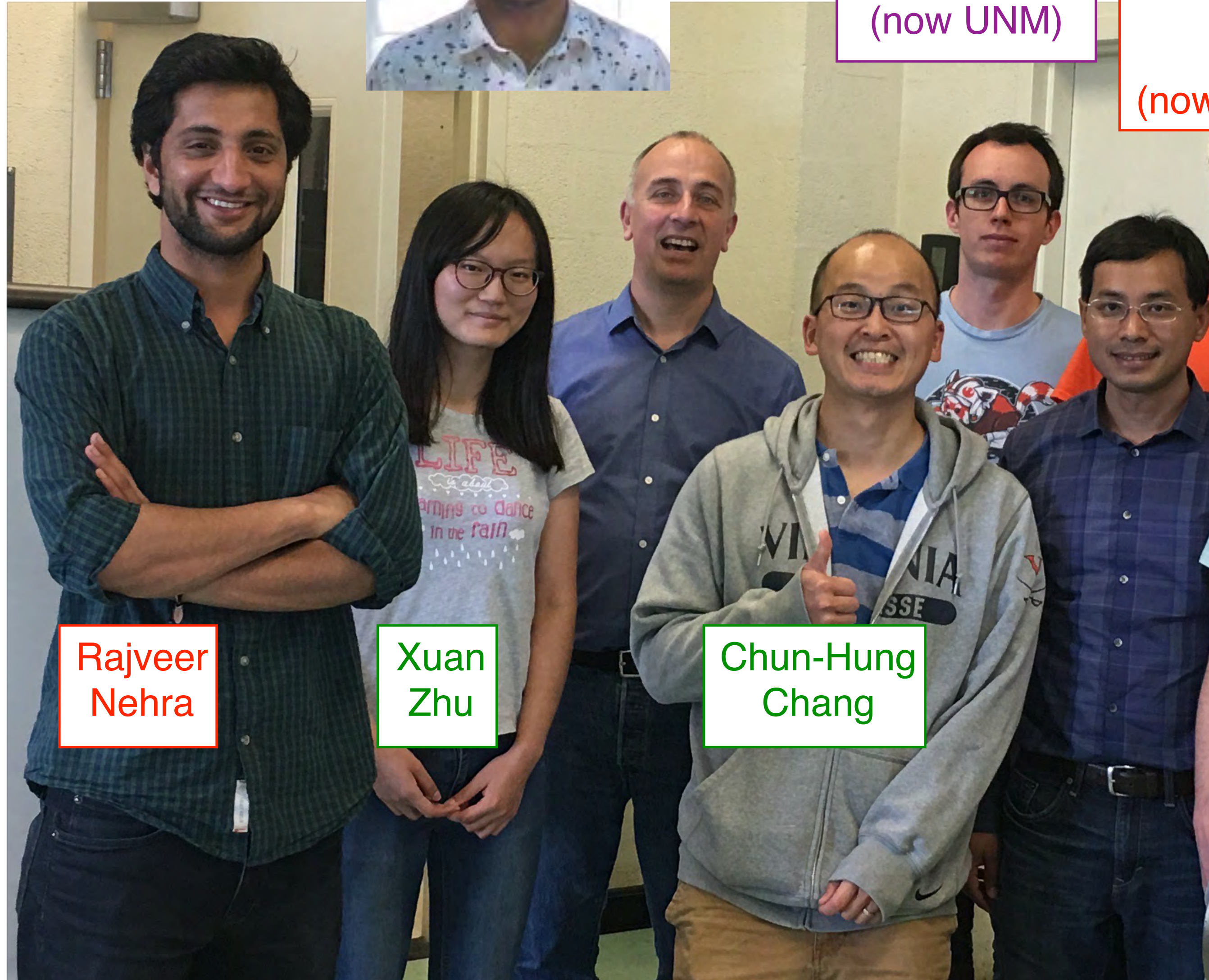


Carlos González Arciniegas



Rafael  
Alexander  
(now UNM)

Aye  
Win  
(now U.OK.)



Rajveer  
Nehra

Xuan  
Zhu

Chun-Hung  
Chang

## The QFQI group



Miller  
Eaton



# Collaborators



Avi  
Pe'er  
Bar-Ilan U.



Paulo  
Nussenzveig  
U. São Paulo



Marcelo  
Martinelli  
U. São Paulo



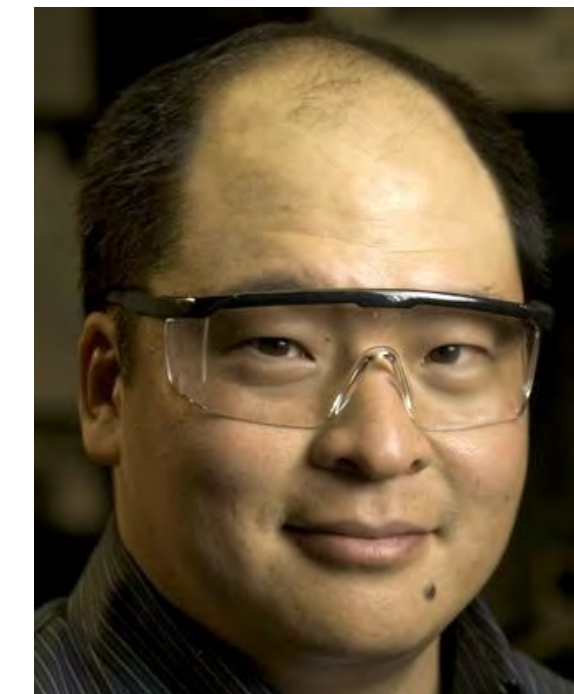
Andreas  
Beling  
UVA ECE



Joe  
Campbell  
UVA ECE



Xu  
Yi  
UVA ECE



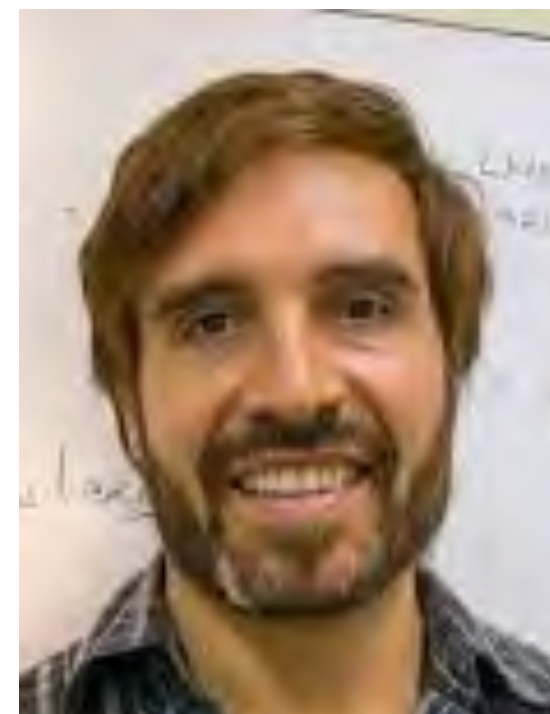
Sae Woo  
Nam  
NIST



Thomas  
Gerrits  
NIST



Myungshik  
Kim  
Imperial  
College



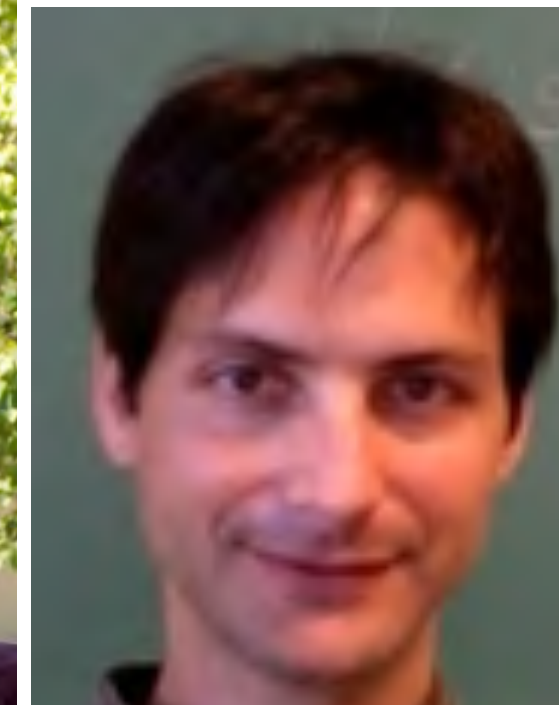
Leandro  
Aolita  
U. Fed. Rio  
Janeiro



Nick  
Menicucci  
RMIT



Raf  
Alexander  
UNM



Israel  
Klich  
UVA



Simonetta  
Liuti  
UVA



# Current projects

- Large-scale cluster-state entanglement
- Quantum state engineering and characterization (quantum error correction encodings)
- Quantum photonics on chip. Collab. w/ UVA ECE: Campbell, Beling, Yi
- Quantum simulation of nuclear physics
- Quantum simulation of condensed matter physics

Today I will address the first two topics.

# Building a continuous-variable quantum computer with light

## Gaussian quantum optics (fields)

- Large-scale entanglement (squeezing)
- No postselection



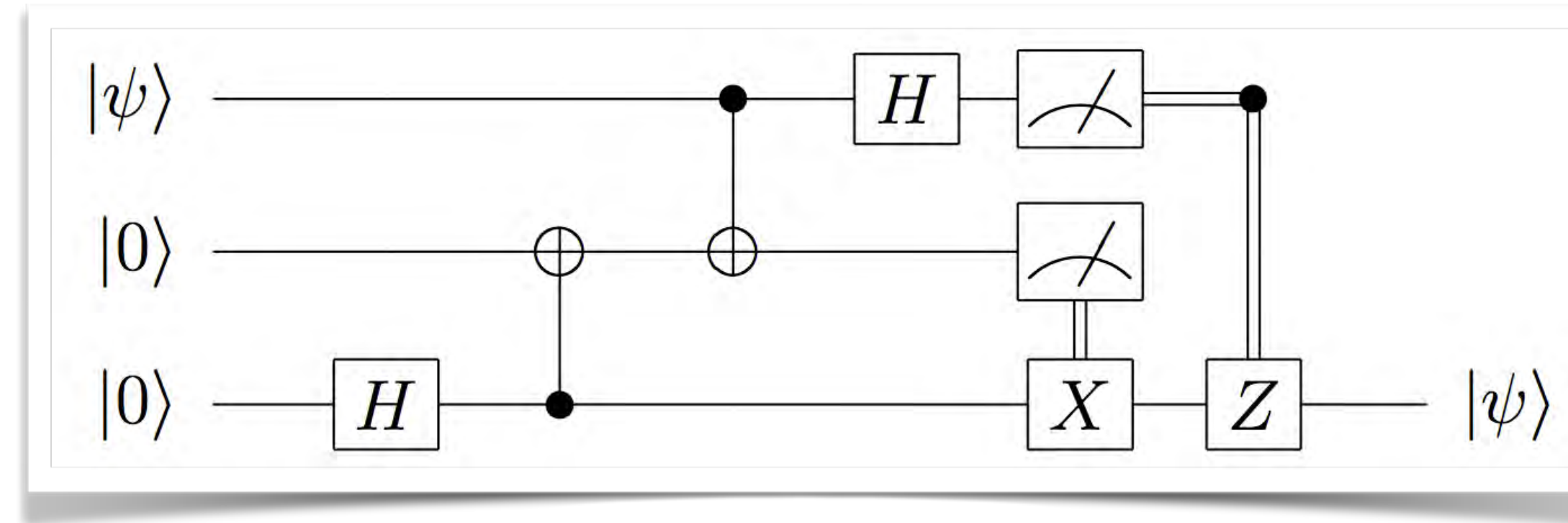
## Non-Gaussian quantum optics (photons)

- Exponential speedup
- Quantum error correction

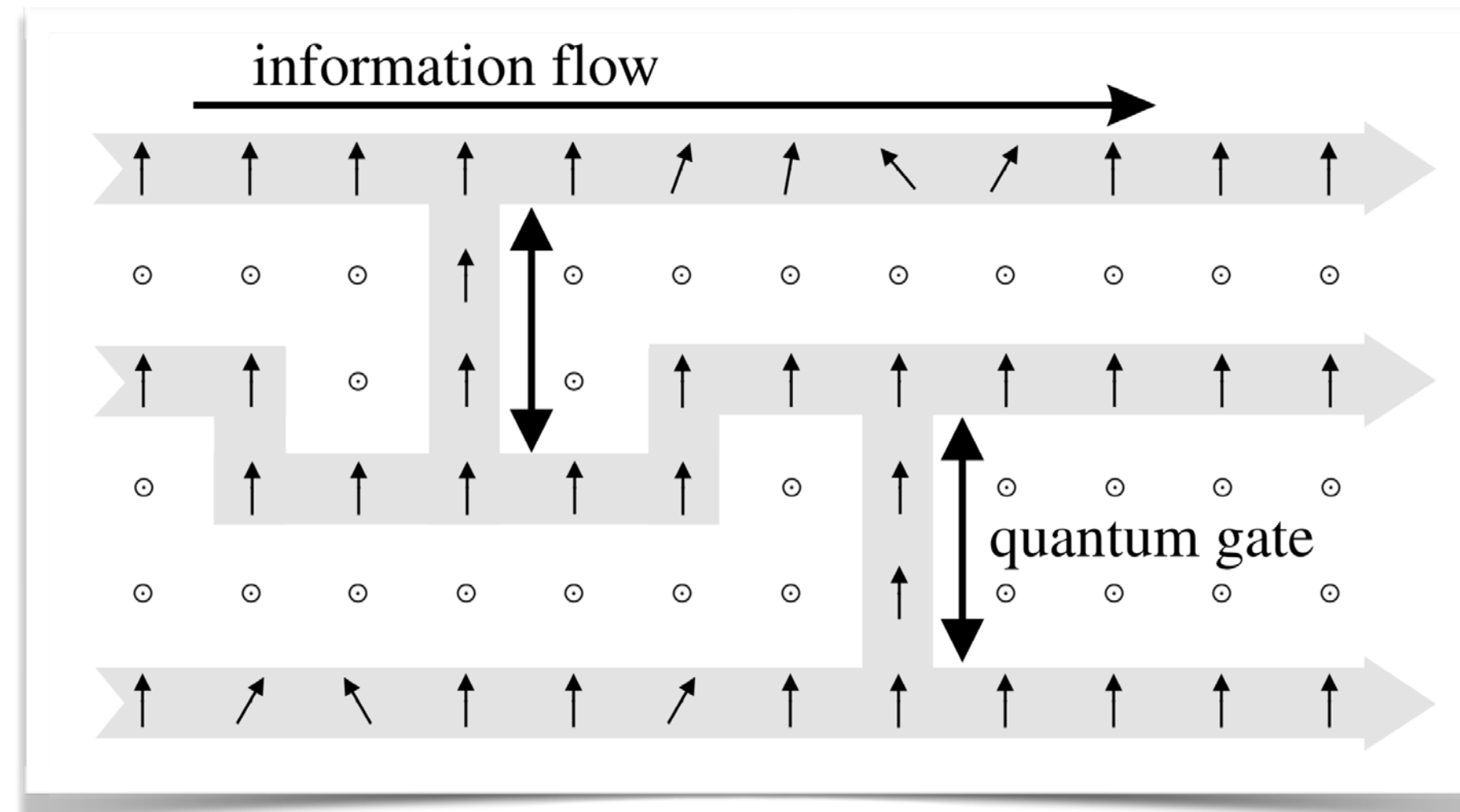


# Two main flavors of universal quantum computing

## 1. Circuit-based



## 2. Measurement-based



VOLUME 86, NUMBER 22

PHYSICAL REVIEW LETTERS

28 MAY 2001

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Robert Raussendorf and Hans J. Briegel

PRL 97, 110501 (2006)

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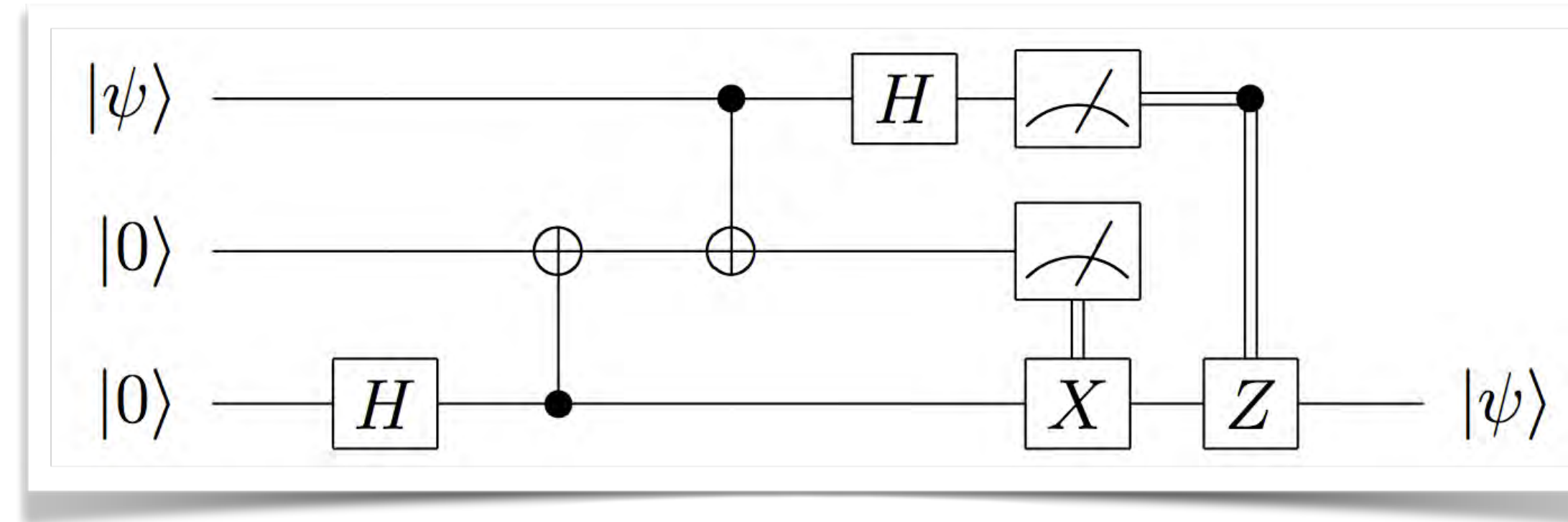
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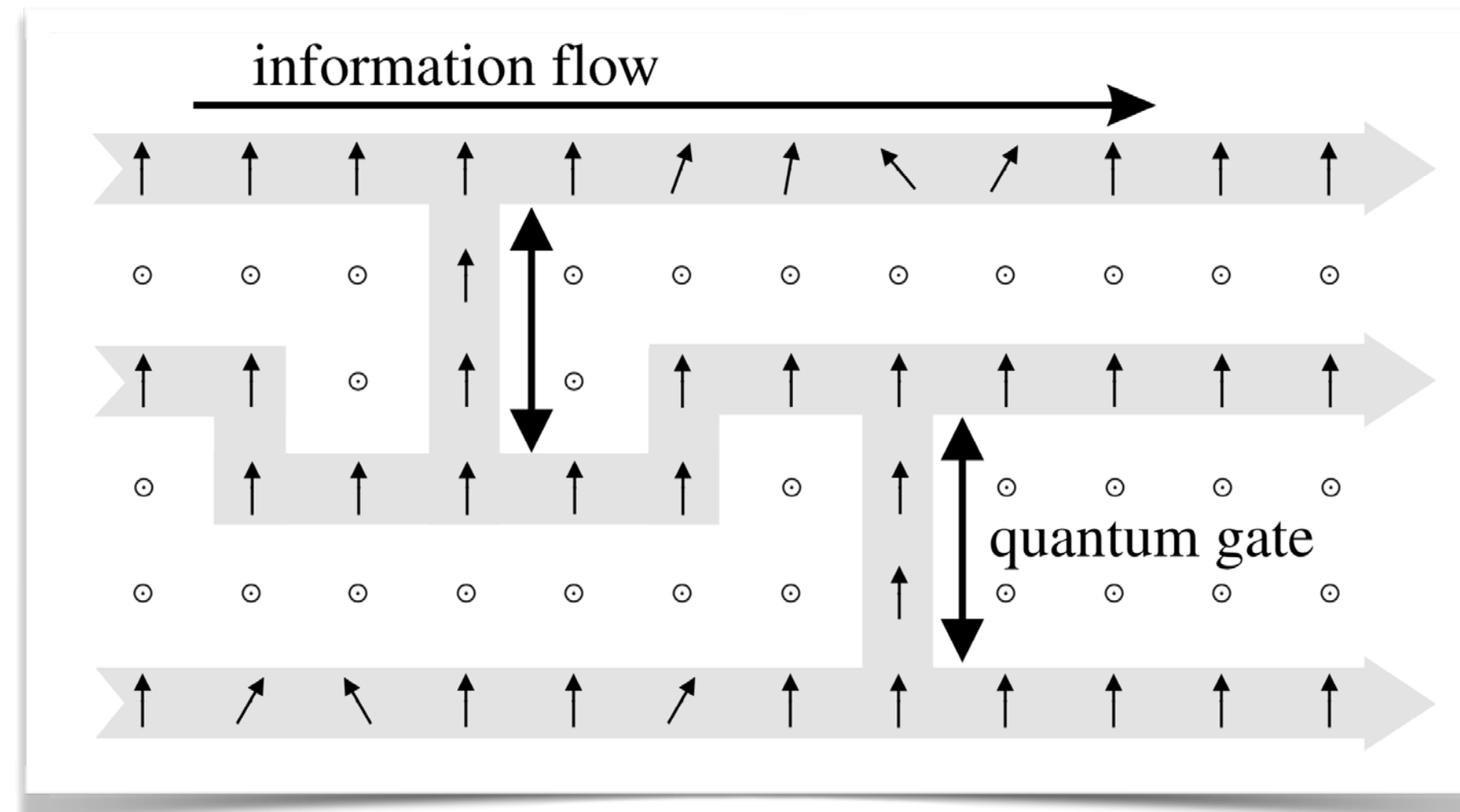
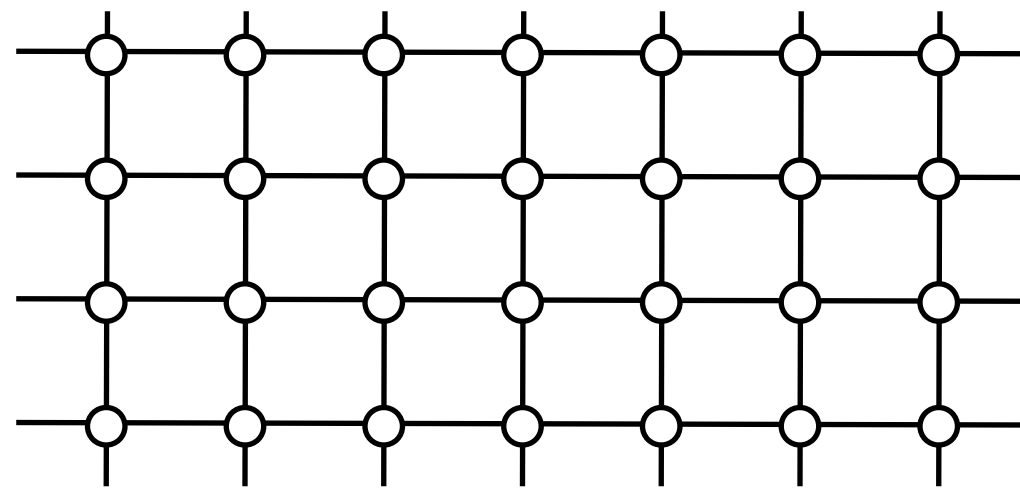


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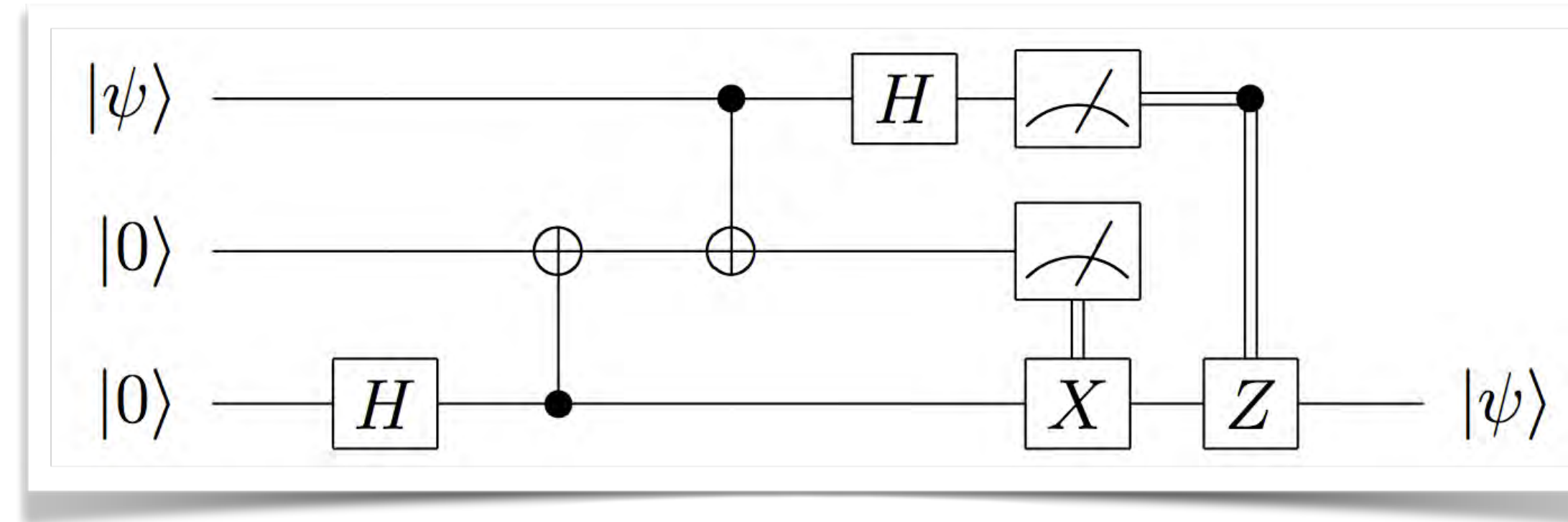
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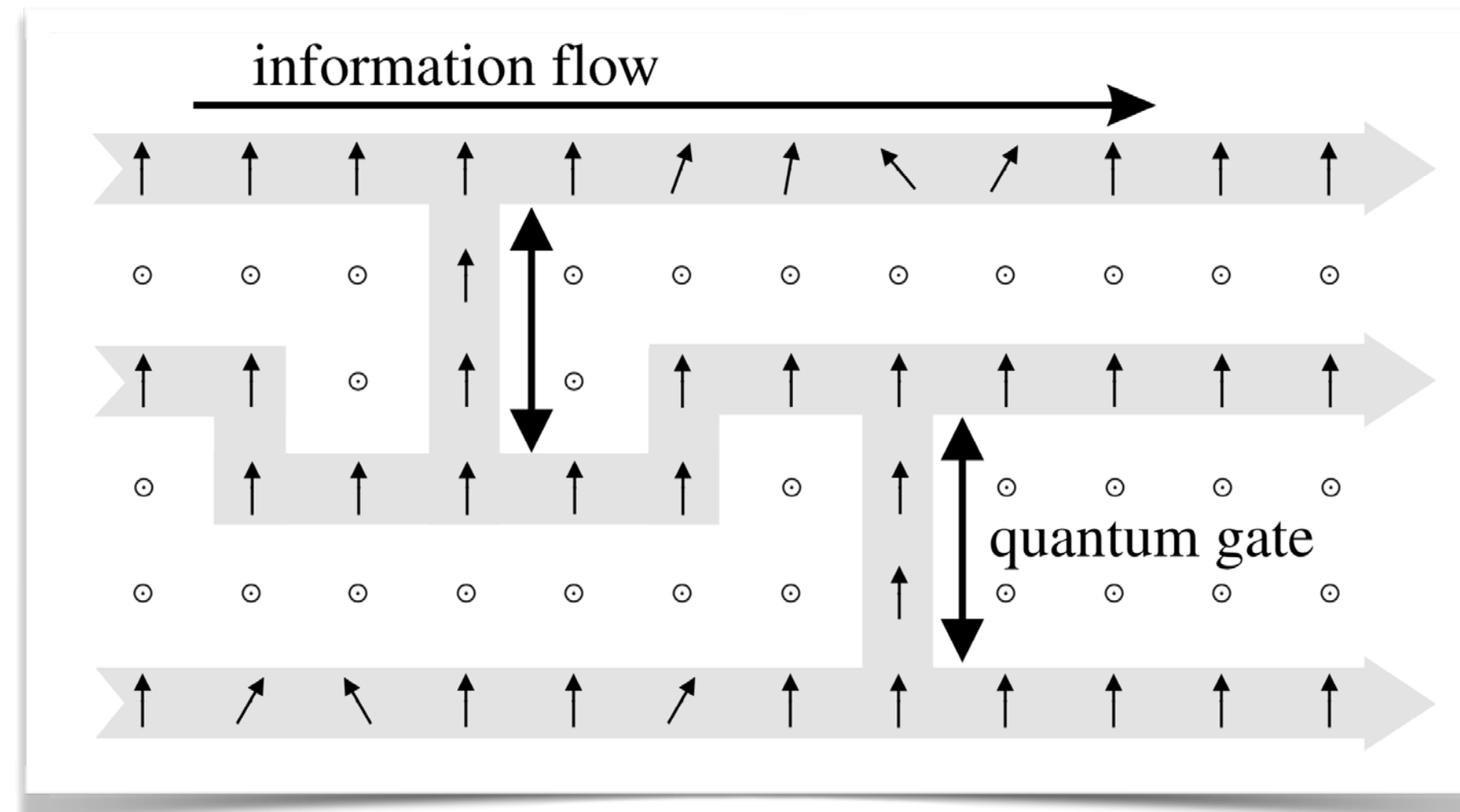
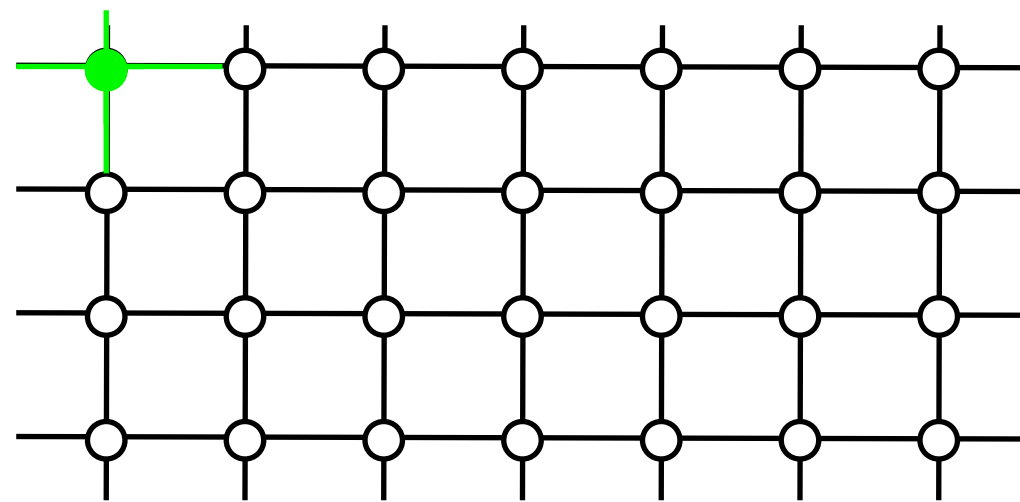


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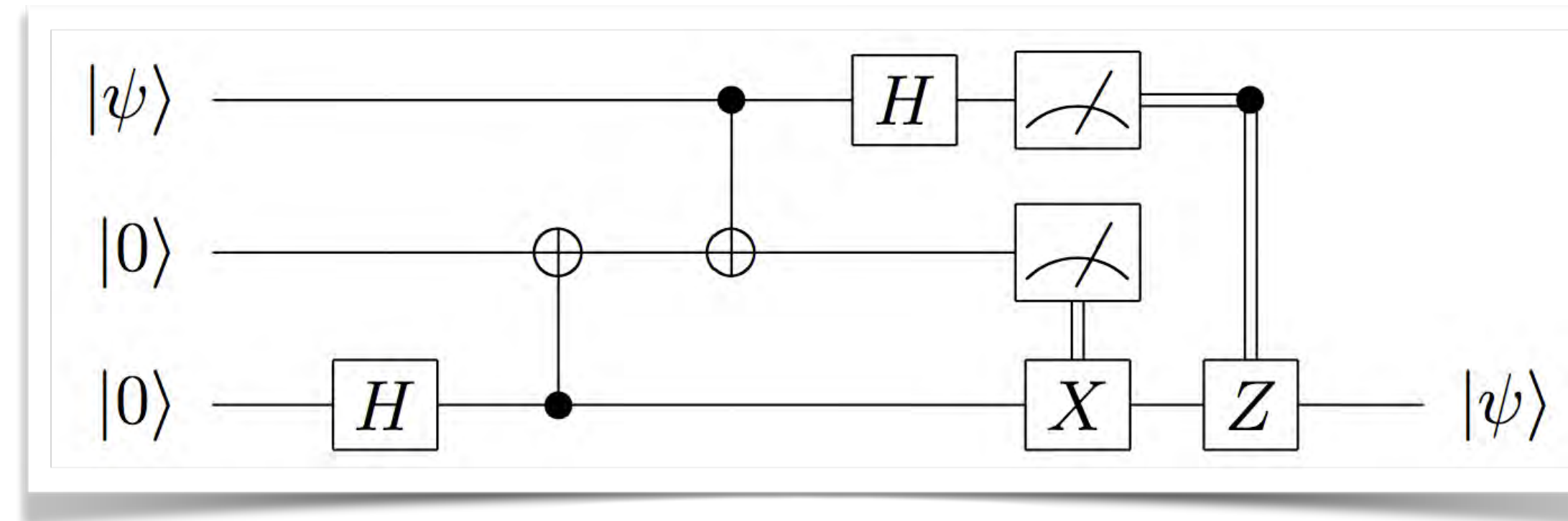
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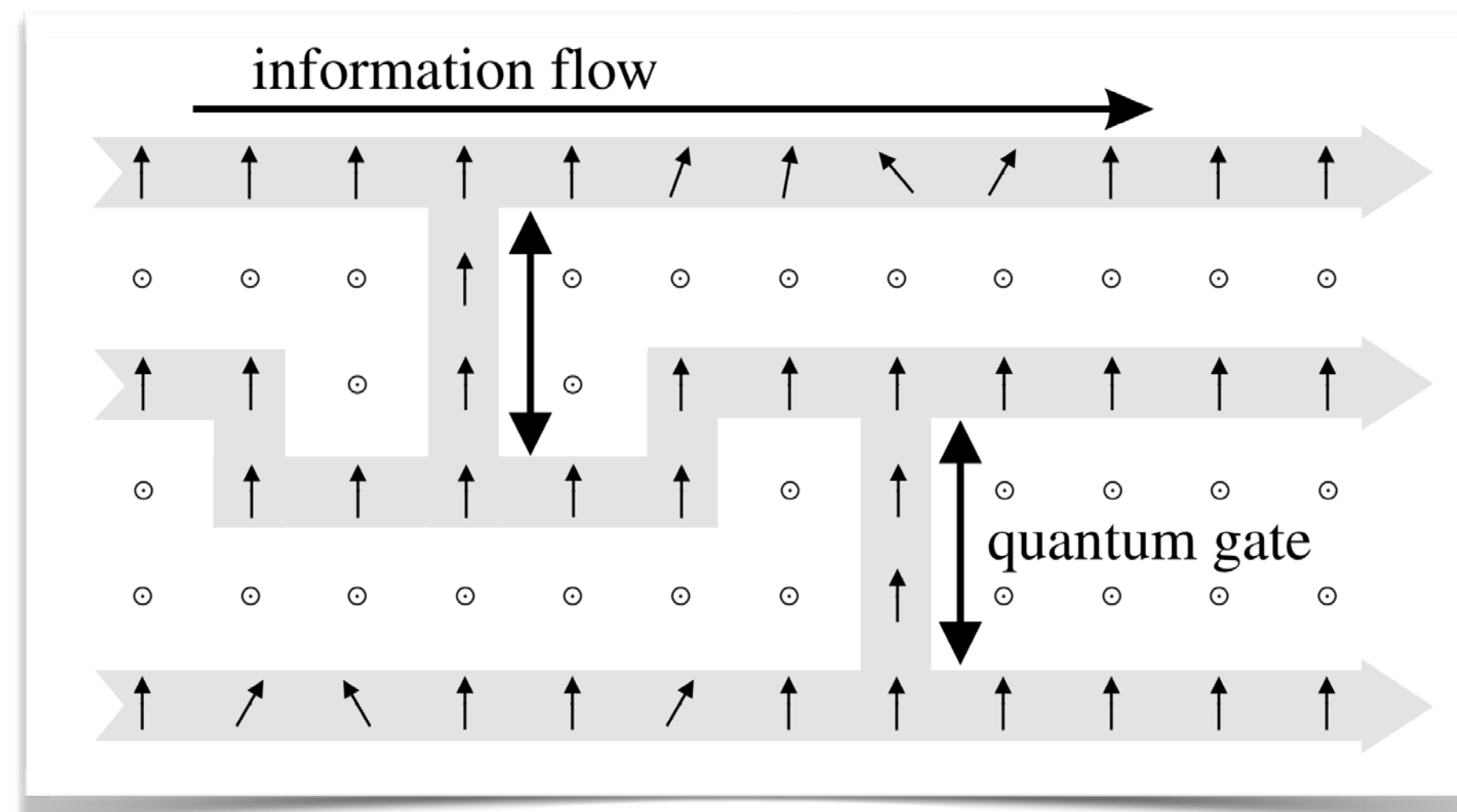
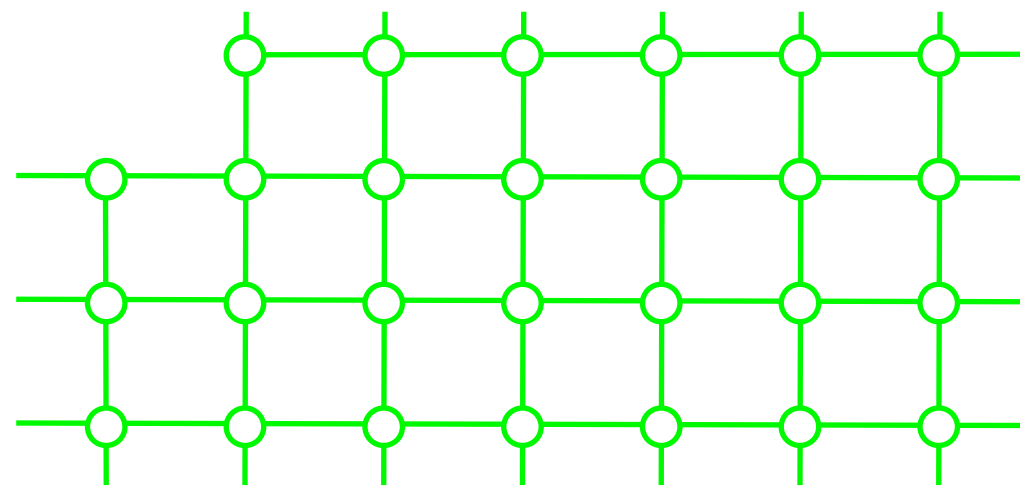
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# Teleportation as a qubit gate primitive

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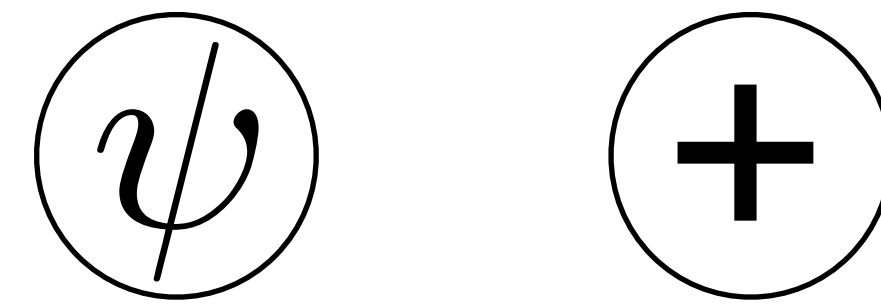
$$|\psi\rangle = \psi_0|0\rangle + \psi_1|1\rangle$$

$$\textcircled{\psi}$$



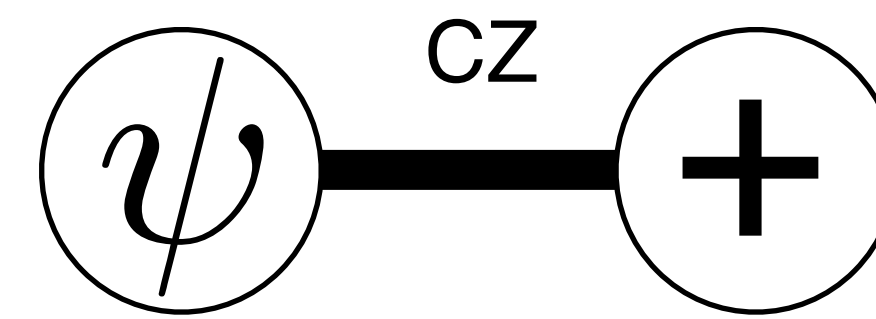
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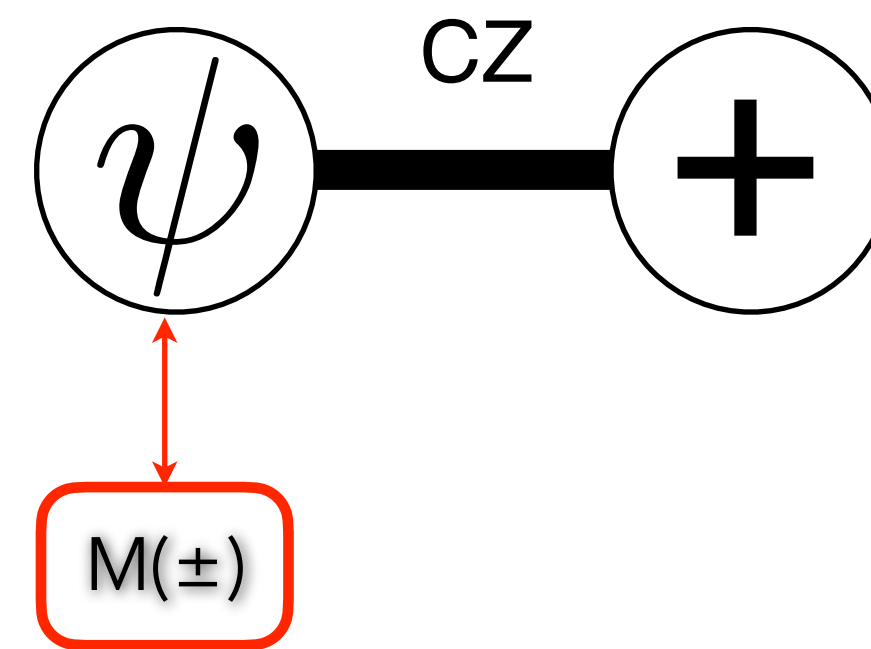
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$$\pm$$

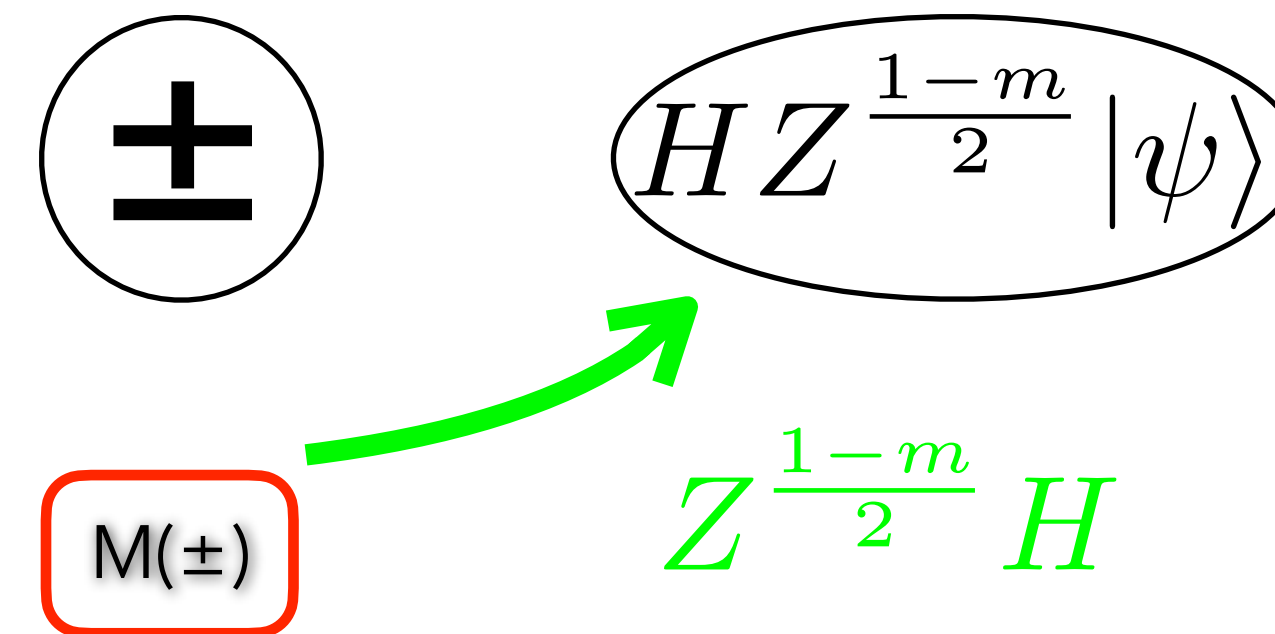
$$H Z^{\frac{1-m}{2}} |\psi\rangle$$

$$M(\pm)$$



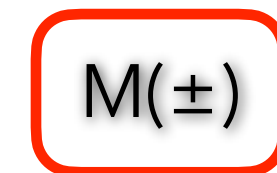
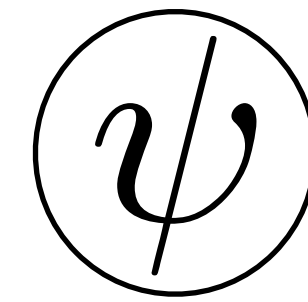
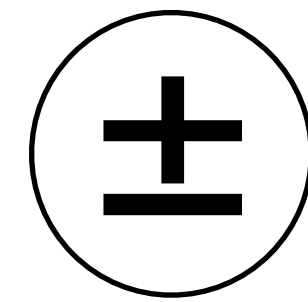
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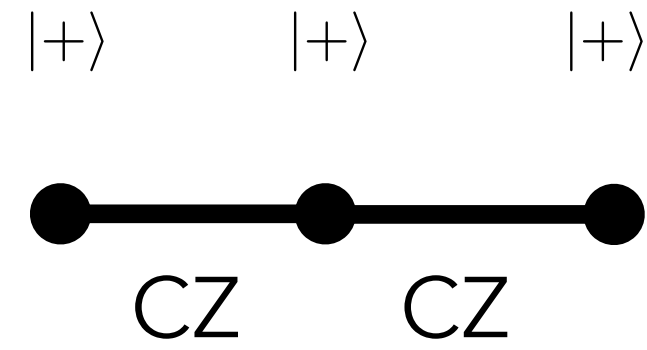


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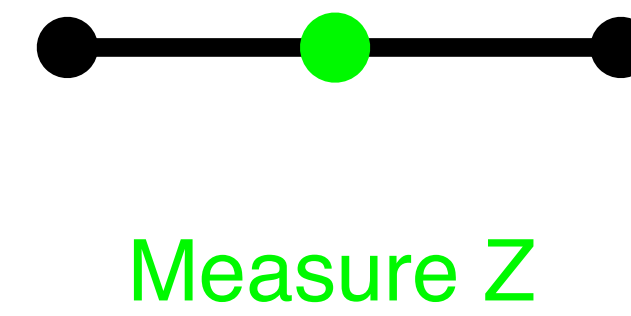
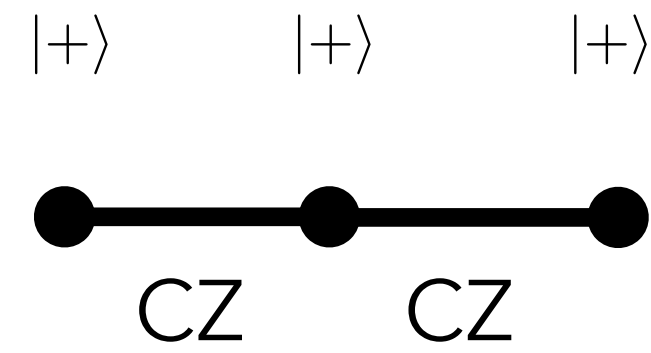


# Can also shape cluster states with measurements

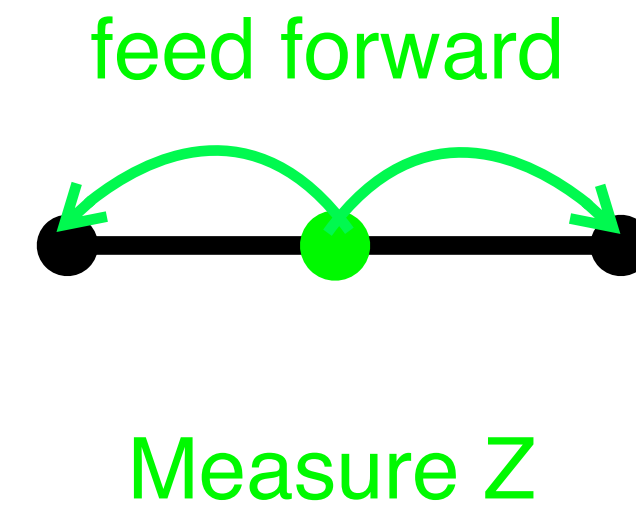
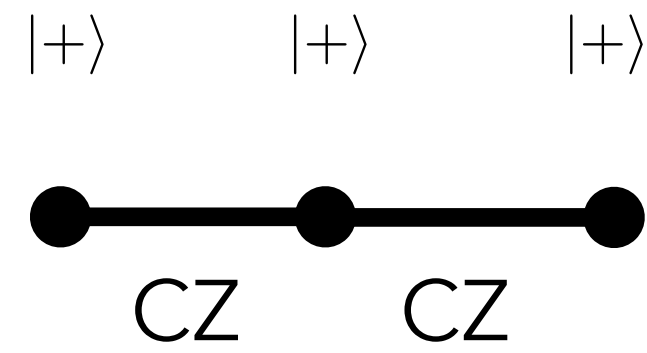




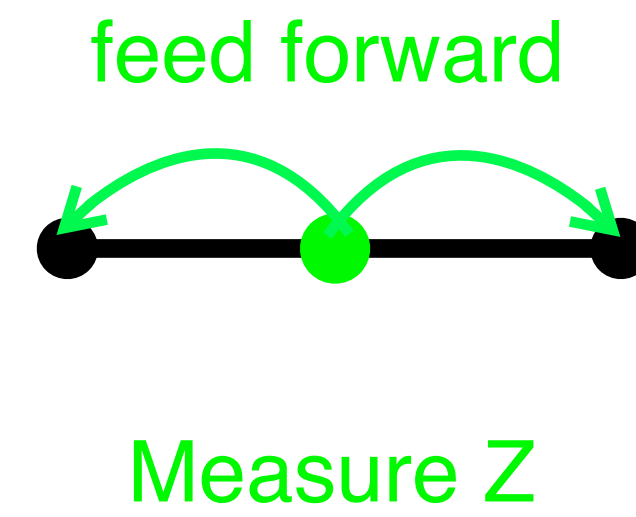
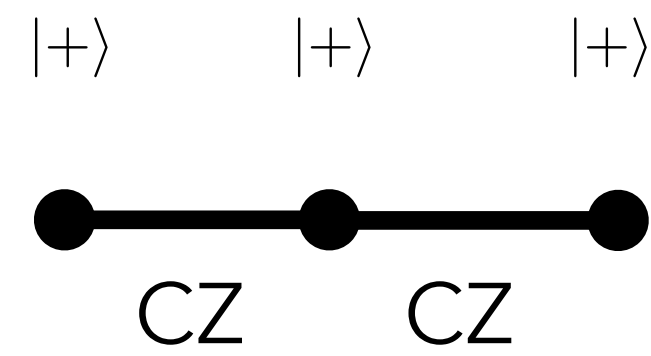
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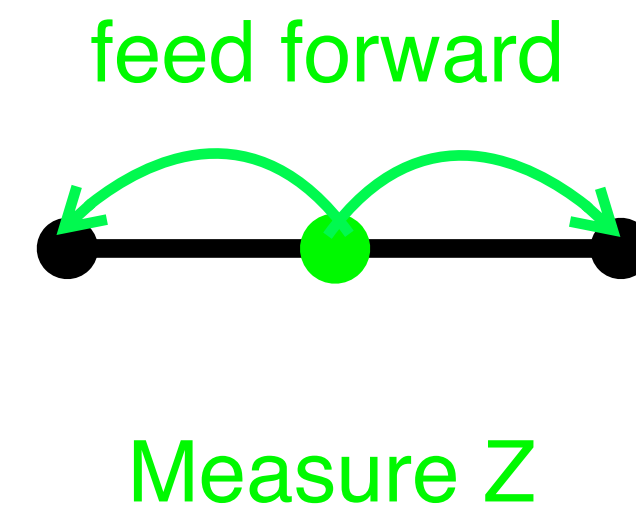
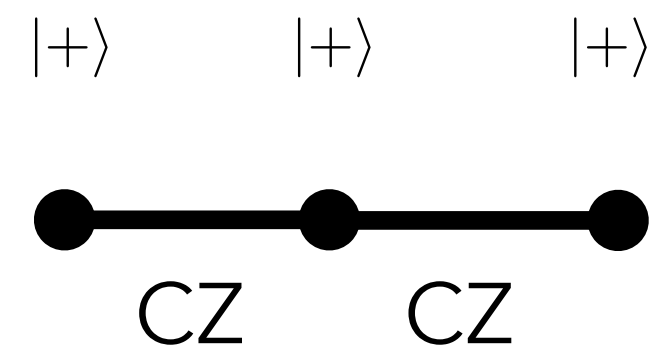


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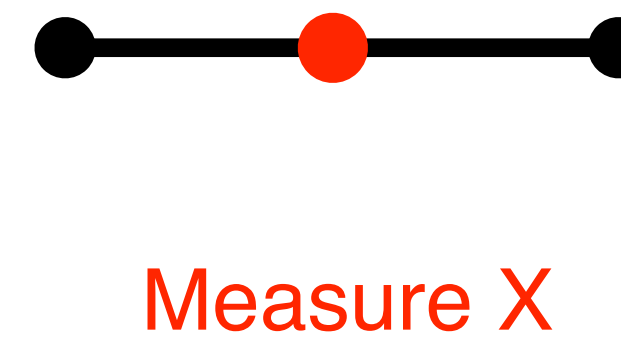
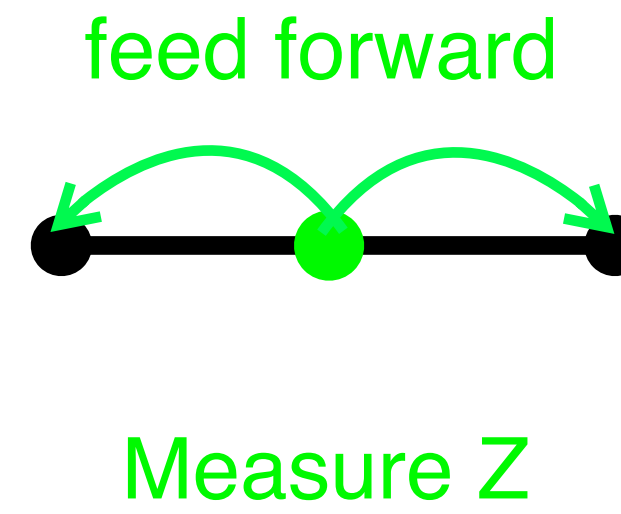
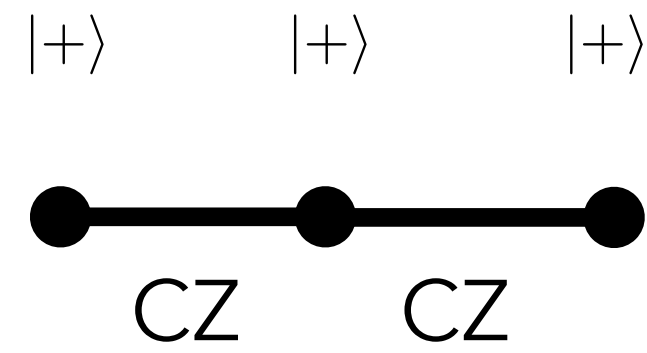




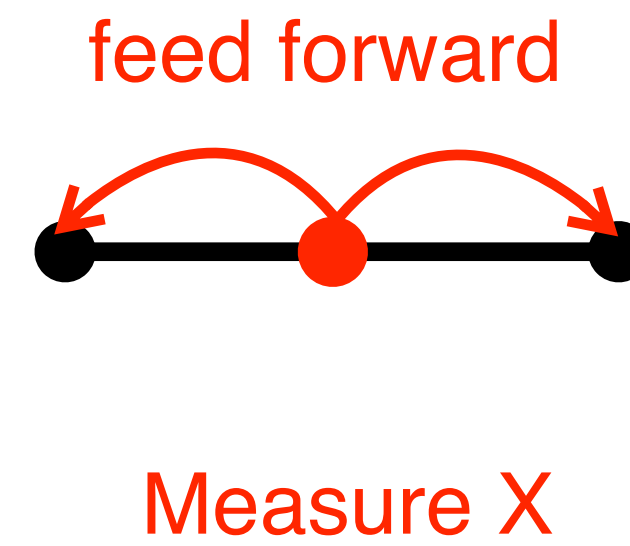
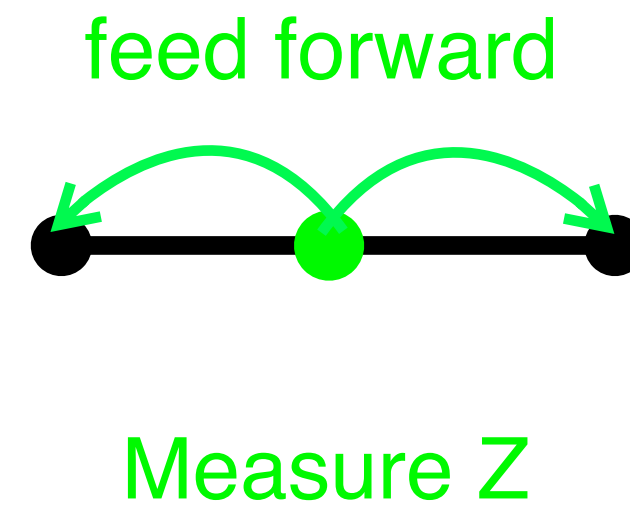
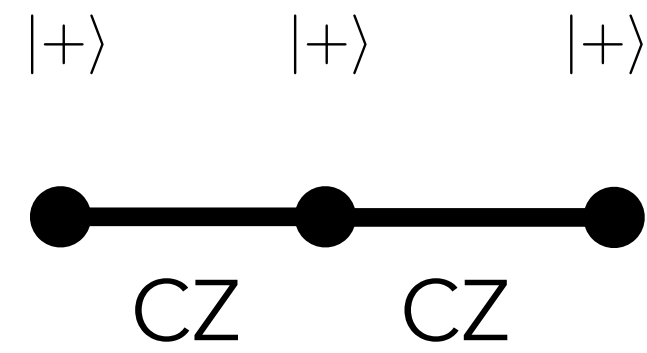
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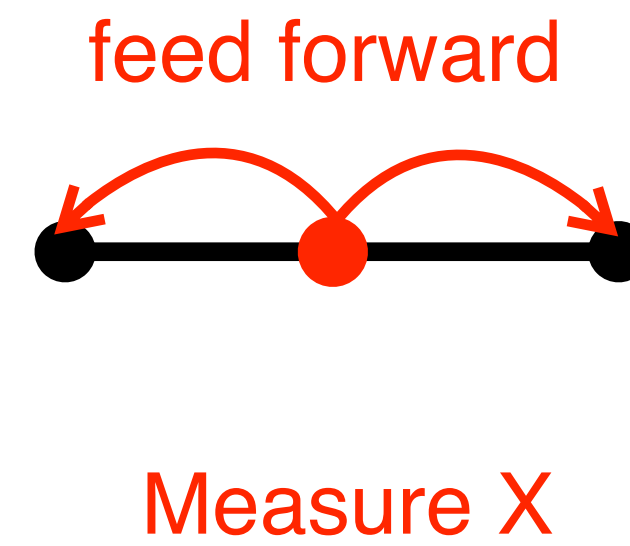
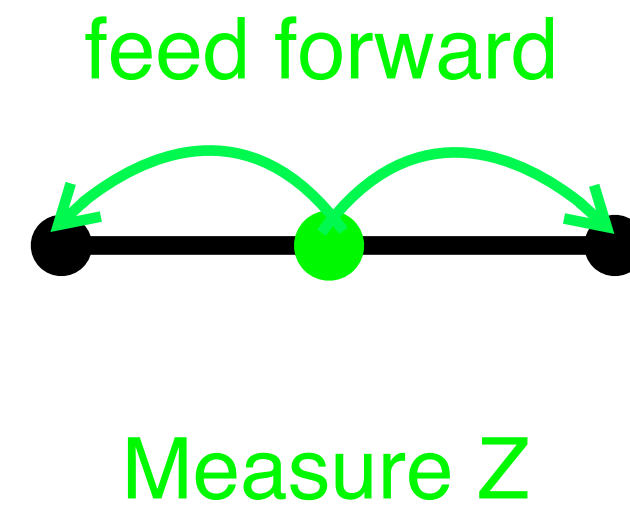
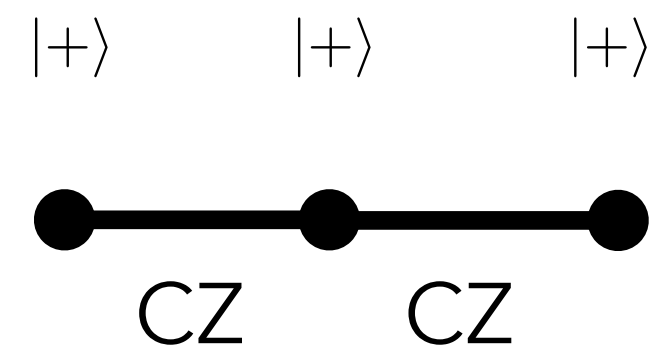


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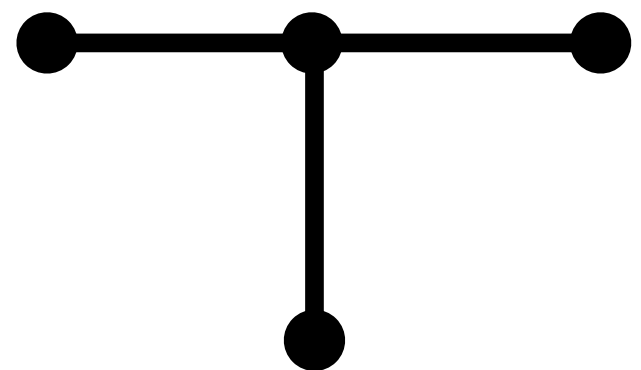
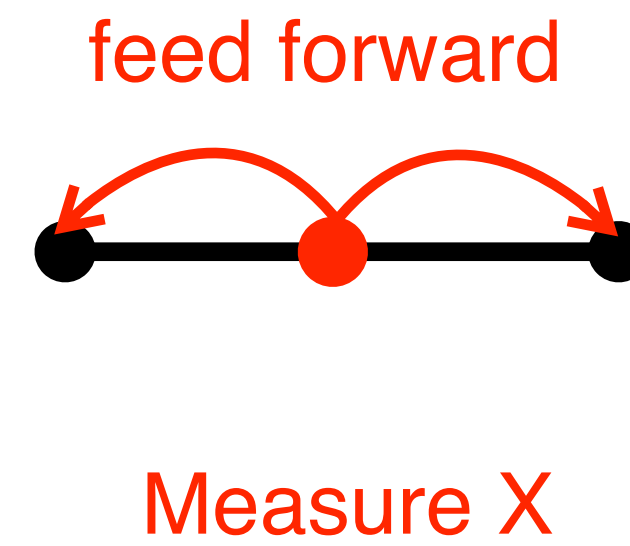
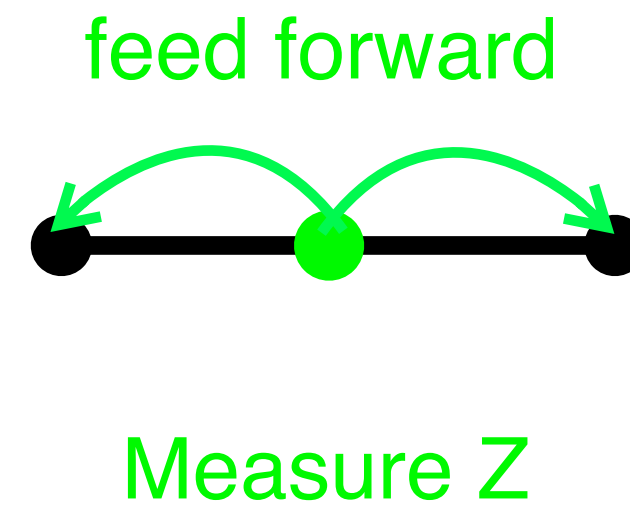
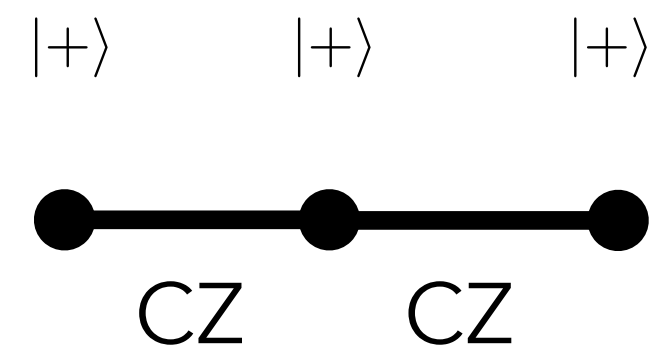




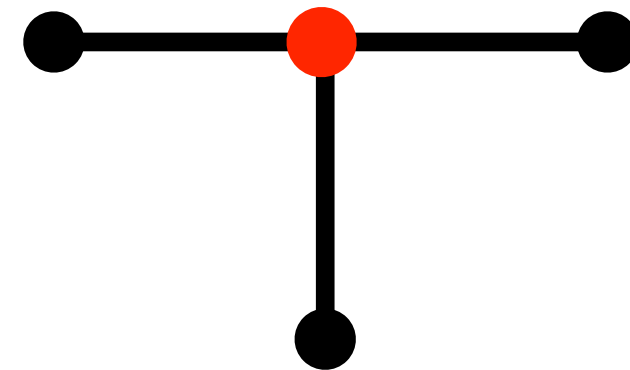
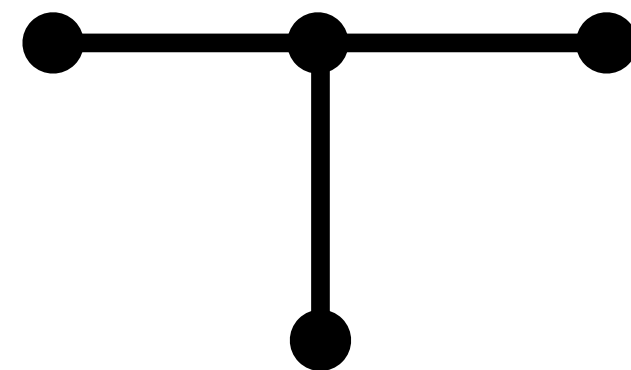
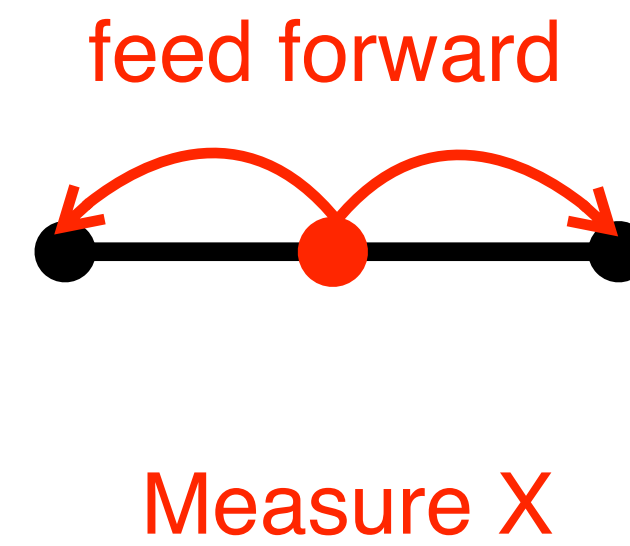
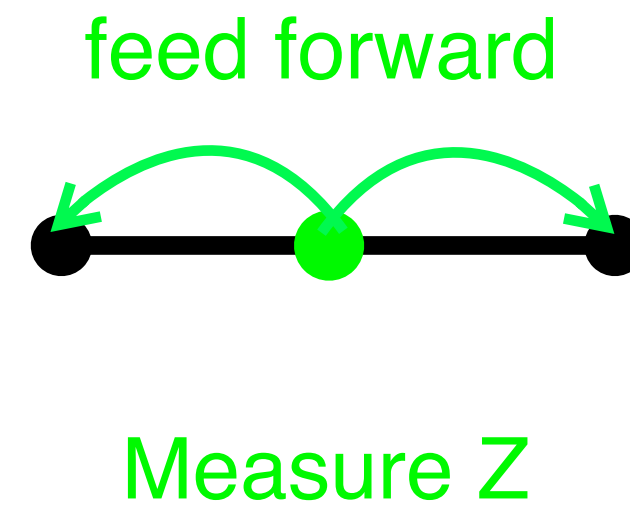
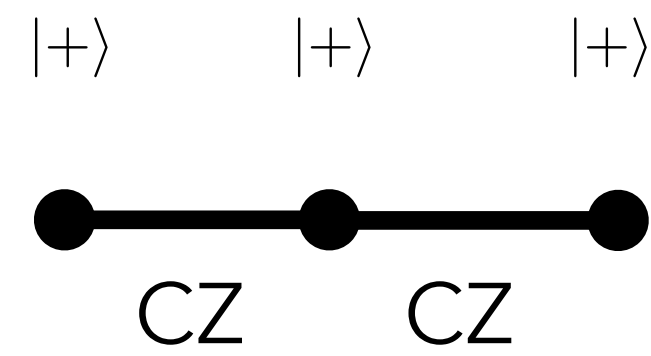
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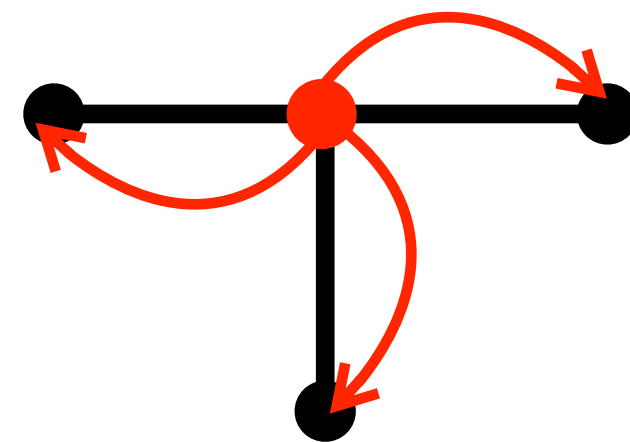
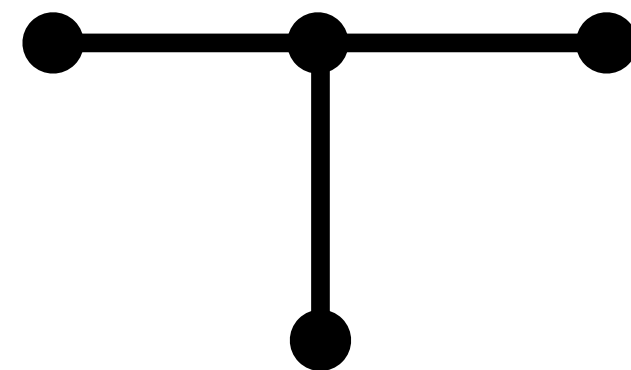
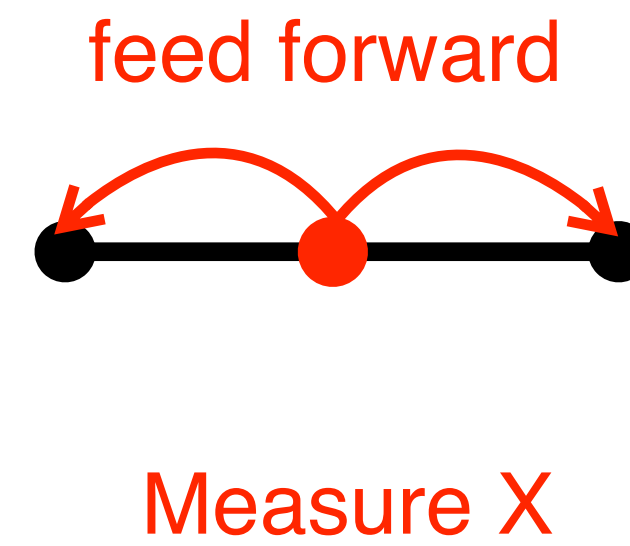
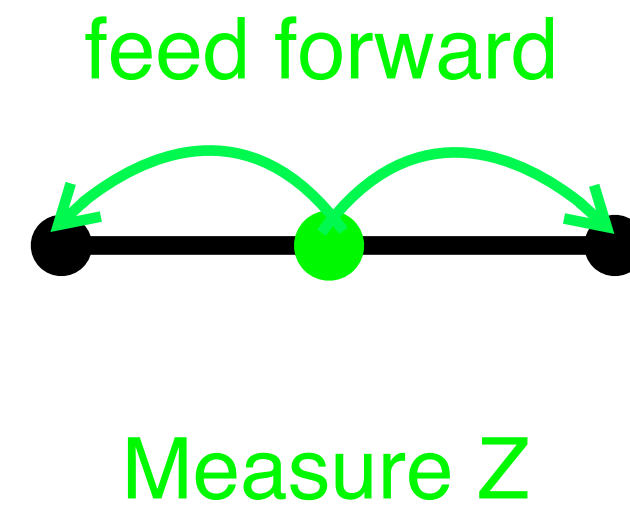
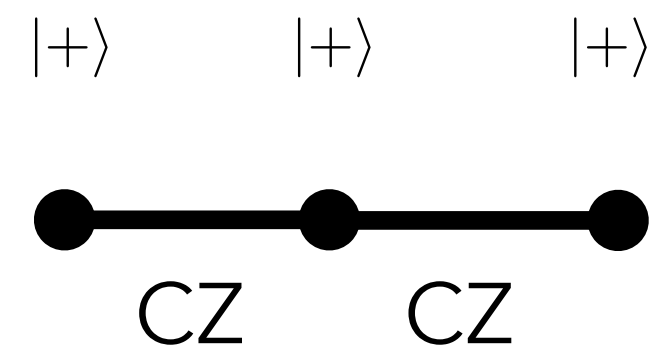


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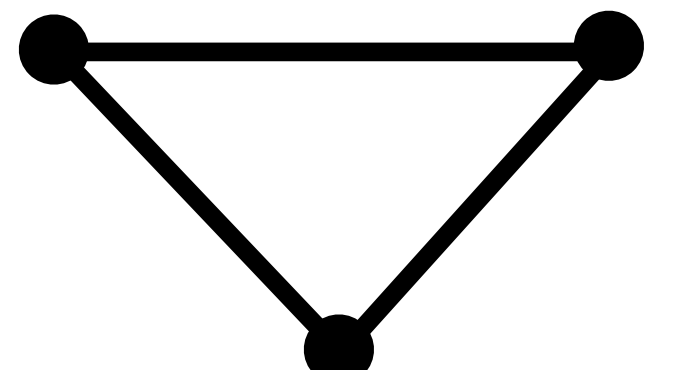
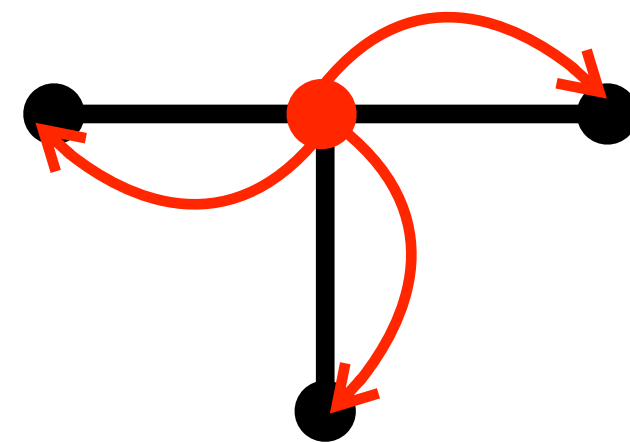
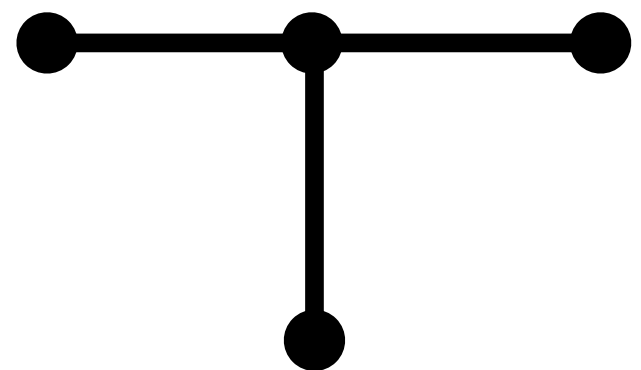
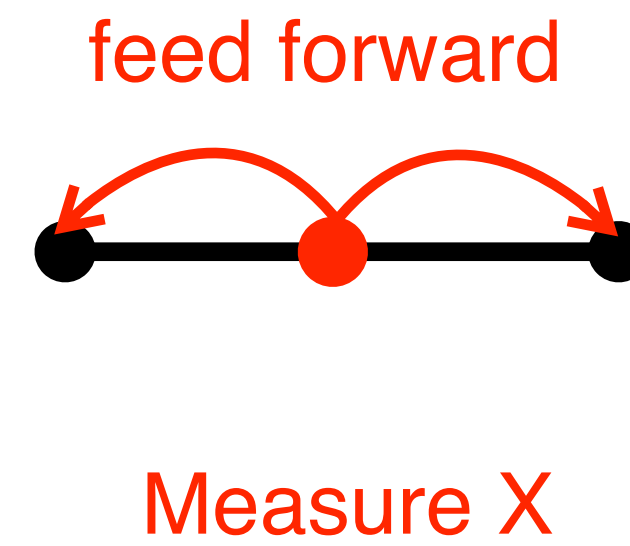
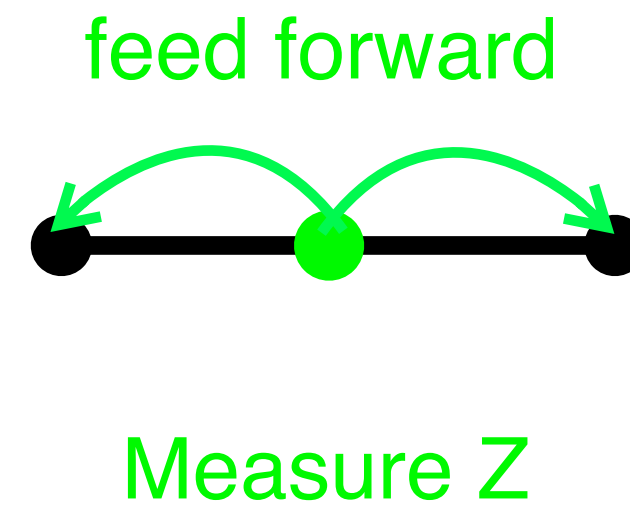
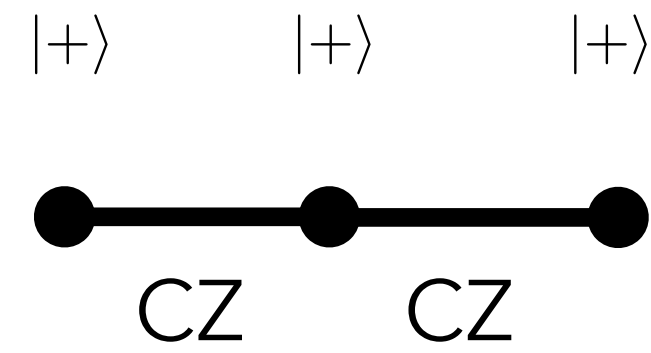




# Can also shape cluster states with measurements



# Can also shape cluster states with measurements



# Discrete vs. continuous variables, qubits vs. qumodes

## Qubits

$$\{|0\rangle, |1\rangle\}$$

e.g. eigenstates of  $S_z$

$$|\psi\rangle = \psi_0|0\rangle + \psi_1|1\rangle$$

↑  
Hadamard  
transform  
↓

$$\{|\pm\rangle\} = \{|0\rangle \pm |1\rangle\}$$

e.g. eigenstates of  $S_x$

$$|00\rangle + |11\rangle$$

Bell state

## Qumodes

$$\{|q\rangle\}_{q \in \mathbb{R}}$$

e.g. eigenstates of  $Q$ : position / amplitude quadr.

$$|\psi\rangle = \int \psi(q)|q\rangle dq$$

↑  
Fourier  
transform  
↓

$$\{|p\rangle\}_{p \in \mathbb{R}} = \left\{ \int e^{ipq} |q\rangle dq \right\}_{p \in \mathbb{R}}$$

e.g. eigenstates of  $P$ : momentum / phase quadr.

$$\int |q\rangle_1 |q\rangle_2 dq$$

EPR state

$$Q = \frac{1}{\sqrt{2}}(a + a^\dagger)$$

$$P = \frac{i}{\sqrt{2}}(a^\dagger - a)$$

$$[Q, P] = i$$

$$\Delta Q \Delta P = \frac{1}{2}$$

# Discrete vs. continuous variables,

## Qubits

$$\{|0\rangle, |1\rangle\}$$

e.g. eigenstates of  $S_z$

$$|\psi\rangle = \psi_0|0\rangle +$$

Hadamard  
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$$\{|\pm\rangle\} = \{|0\rangle \pm |1\rangle\}$$

e.g. eigenstates of  $S_x$

$$|00\rangle + |11\rangle$$

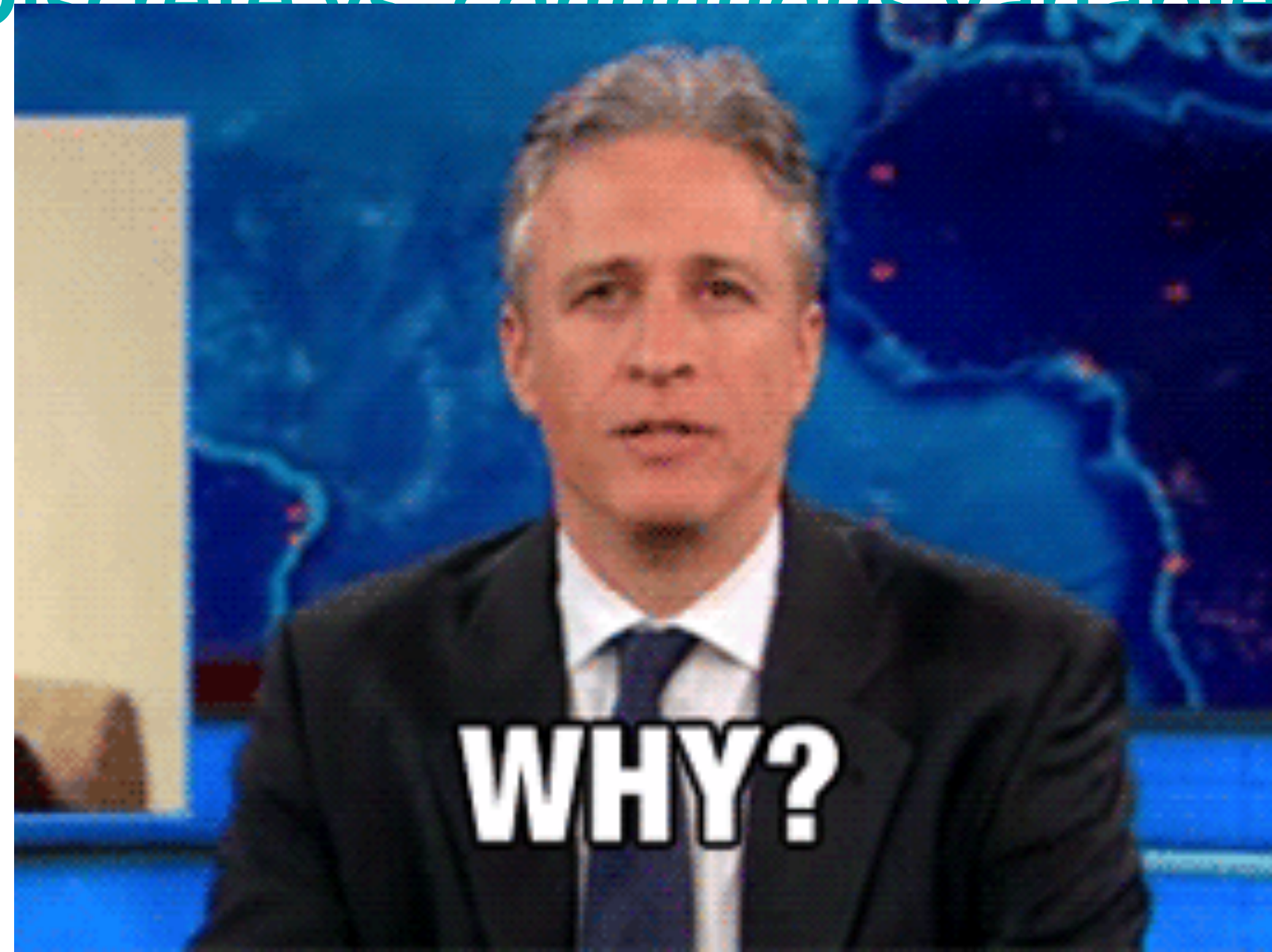
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EPR state



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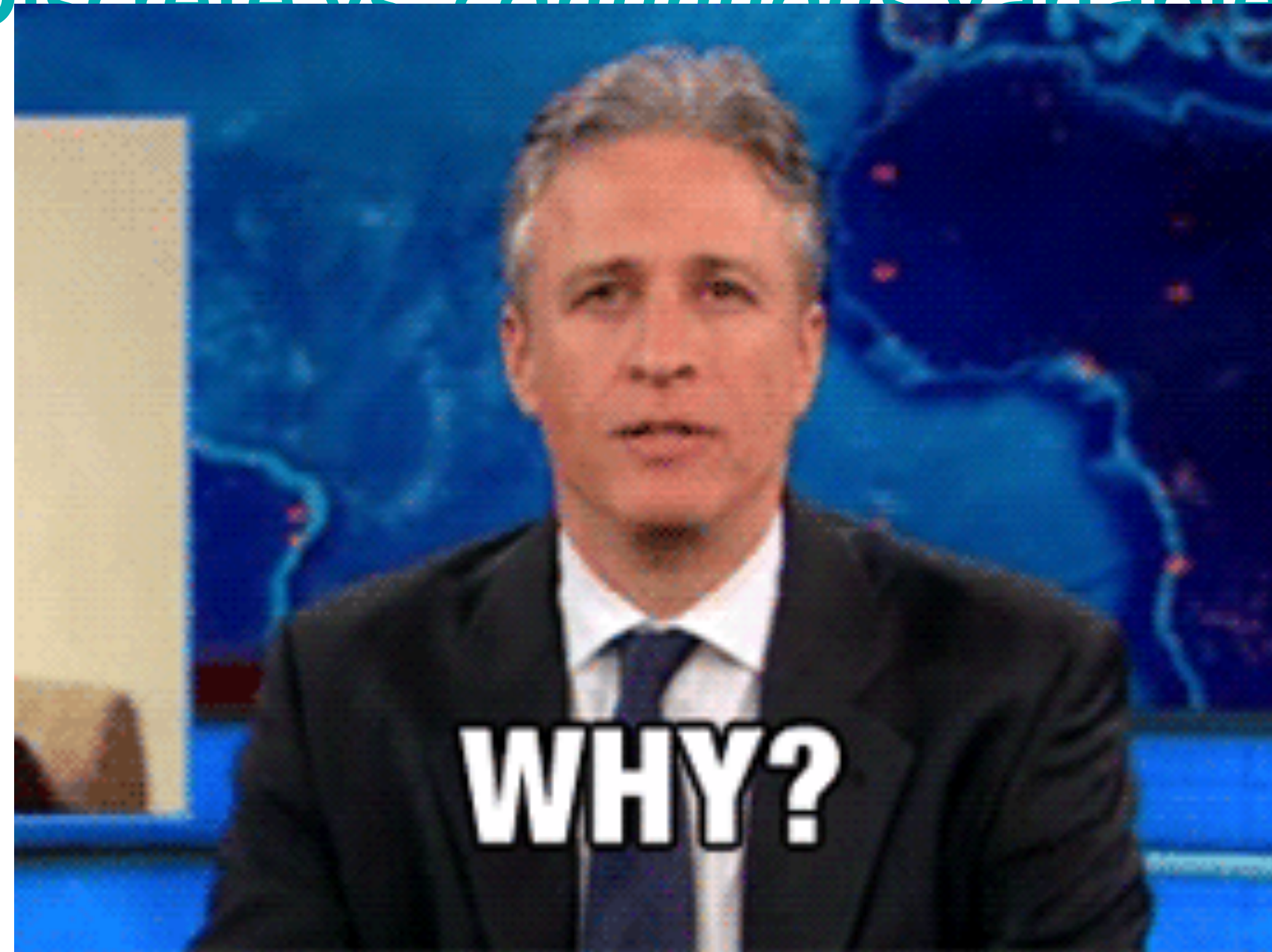
$$[Q, P] = i$$

$$\Delta Q \Delta P = \frac{1}{2}$$

de quadr.



# Discrete vs. continuous variables,



$$Q = \frac{1}{\sqrt{2}}(a + a^\dagger)$$

$$P = \frac{i}{\sqrt{2}}(a^\dagger - a)$$

$$[Q, P] = i$$

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Because CV SCALE

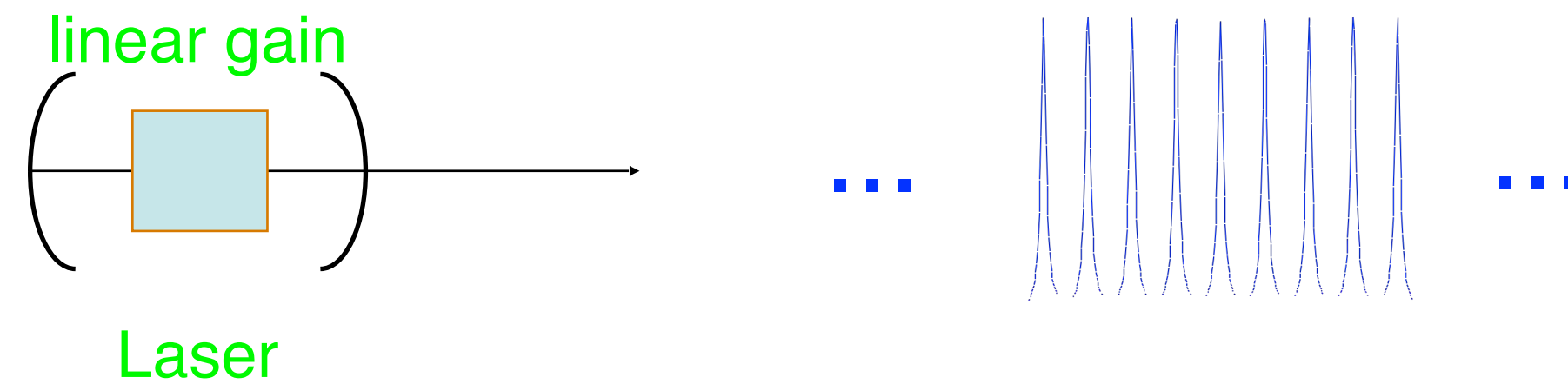


# Gaussian quantum optics (EM fields)



# A good starting point: the quantum optical frequency comb

The eigenmodes of a cavity form a **large** ensemble of **classically coherent** modes



Carrier-envelope-phase locked mode-locked laser = **optical frequency comb (OFC)**  
( $10^6$  modes oscillating in phase)

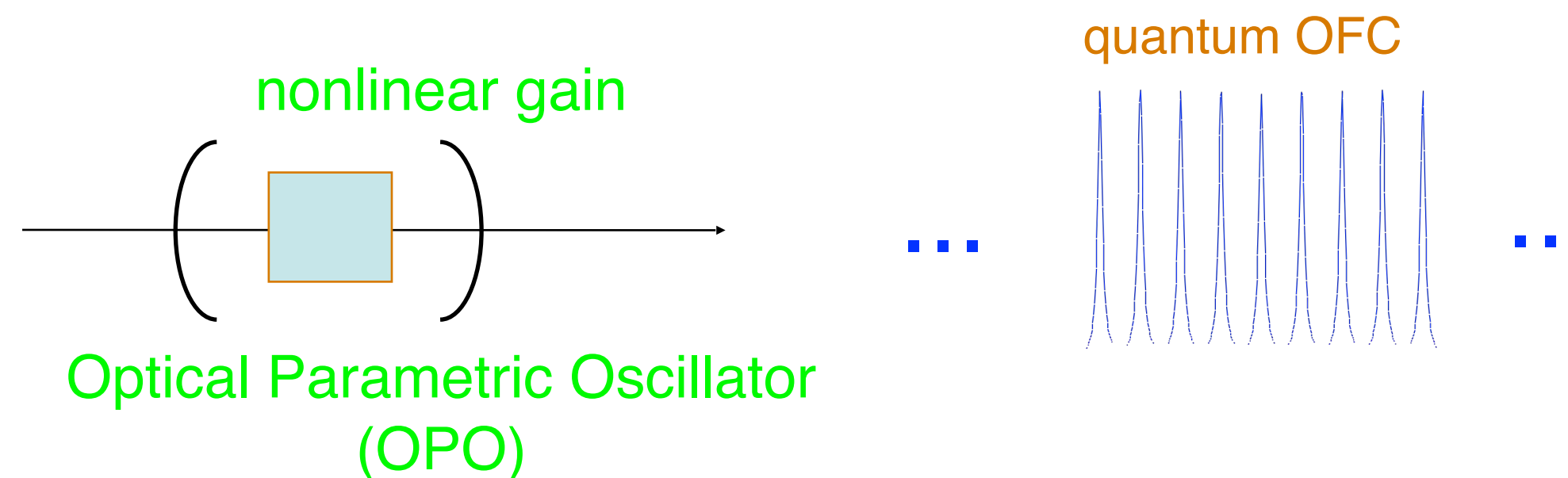


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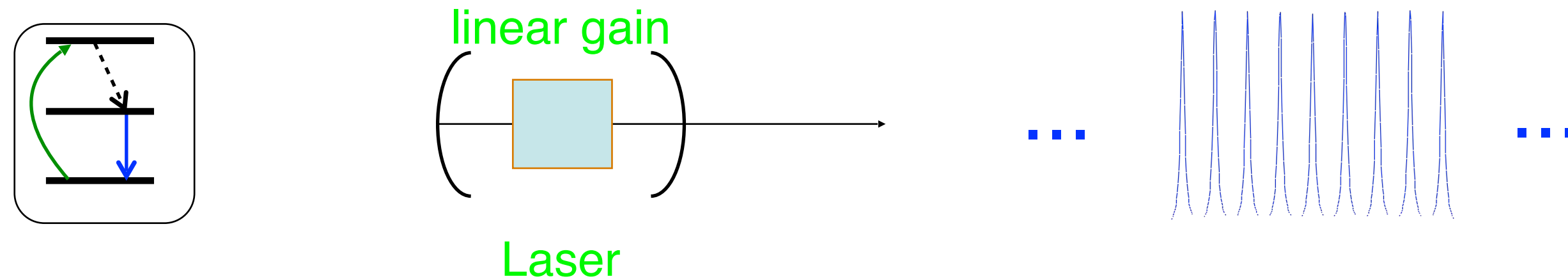
Theodor W. Hänsch

## Why not turn the QOFC into a quantum computer?



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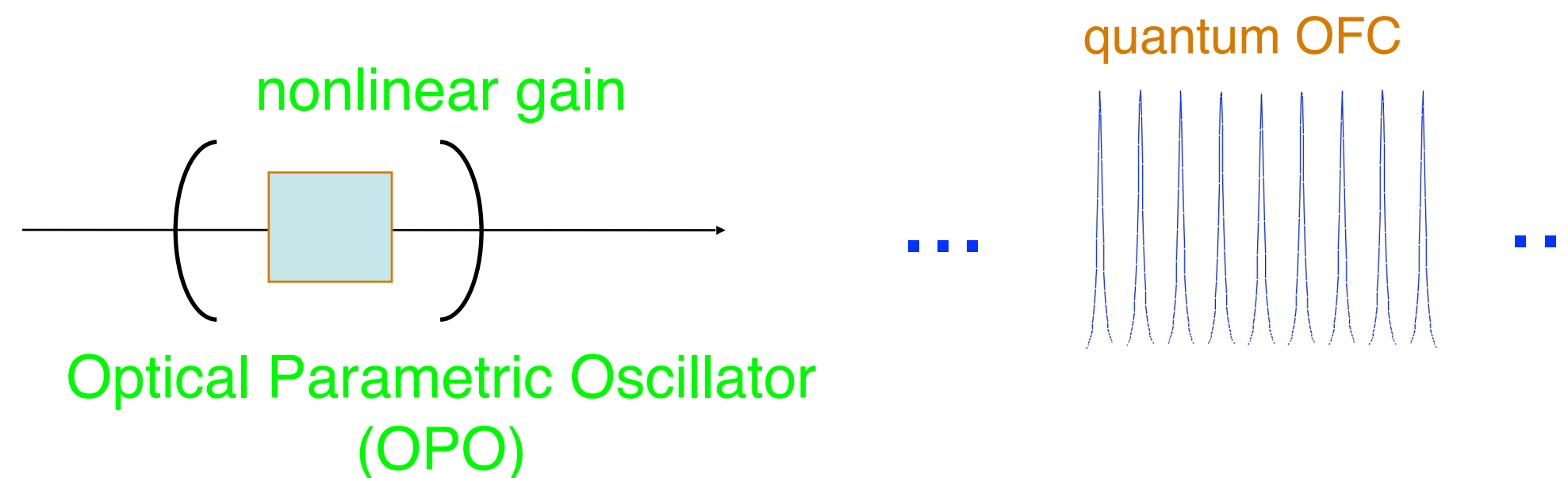


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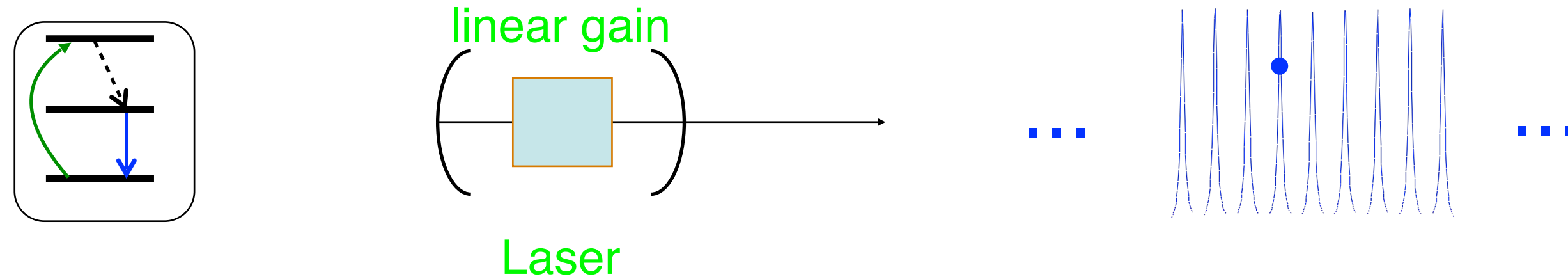
## Why not turn the QOFC into a quantum computer?





# A good starting point: the quantum optical frequency comb

The eigenmodes of a cavity form a **large** ensemble of **classically coherent** modes



Carrier-envelope-phase locked mode-locked laser = **optical frequency comb (OFC)**  
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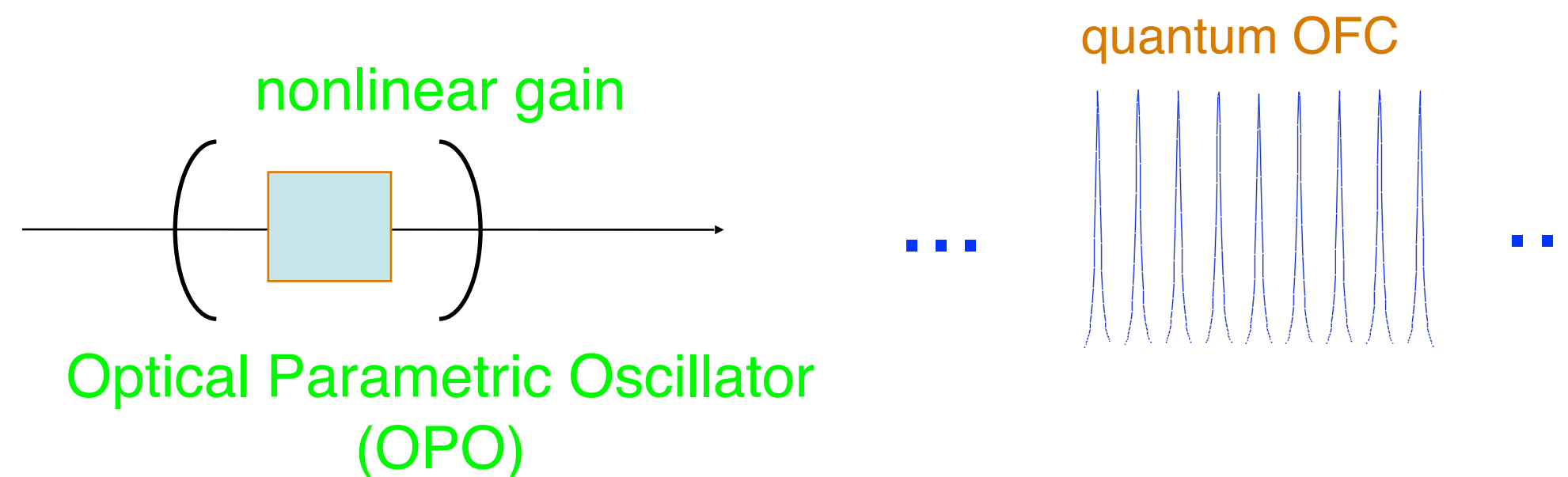


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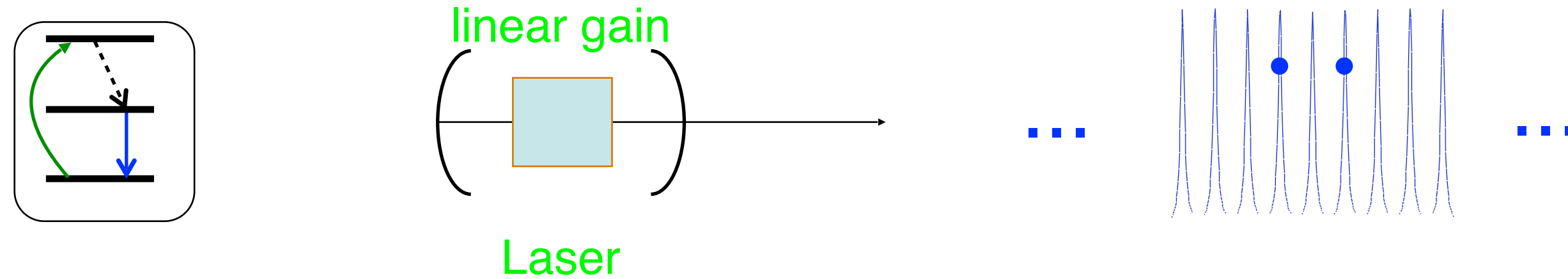
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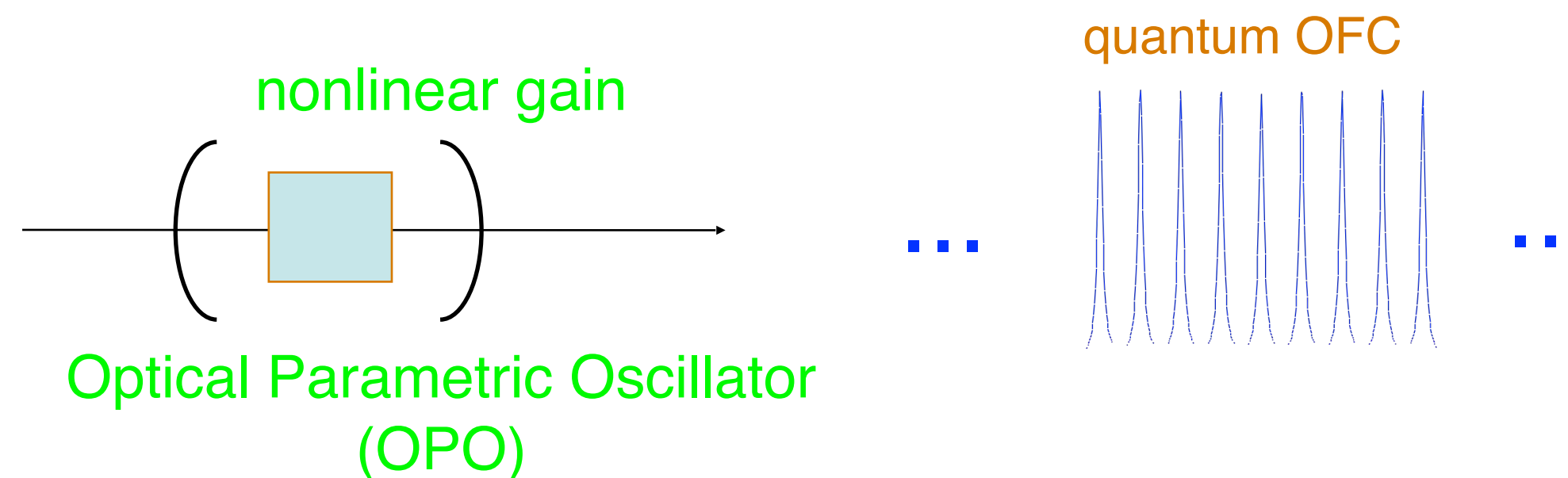


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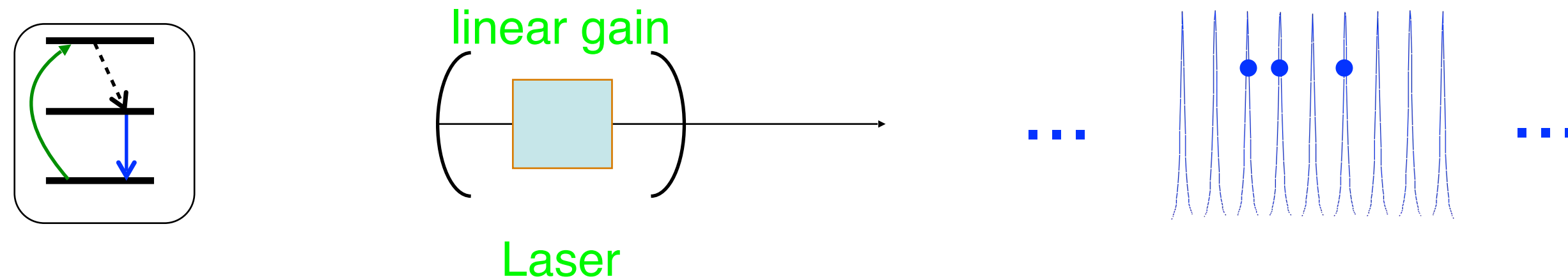
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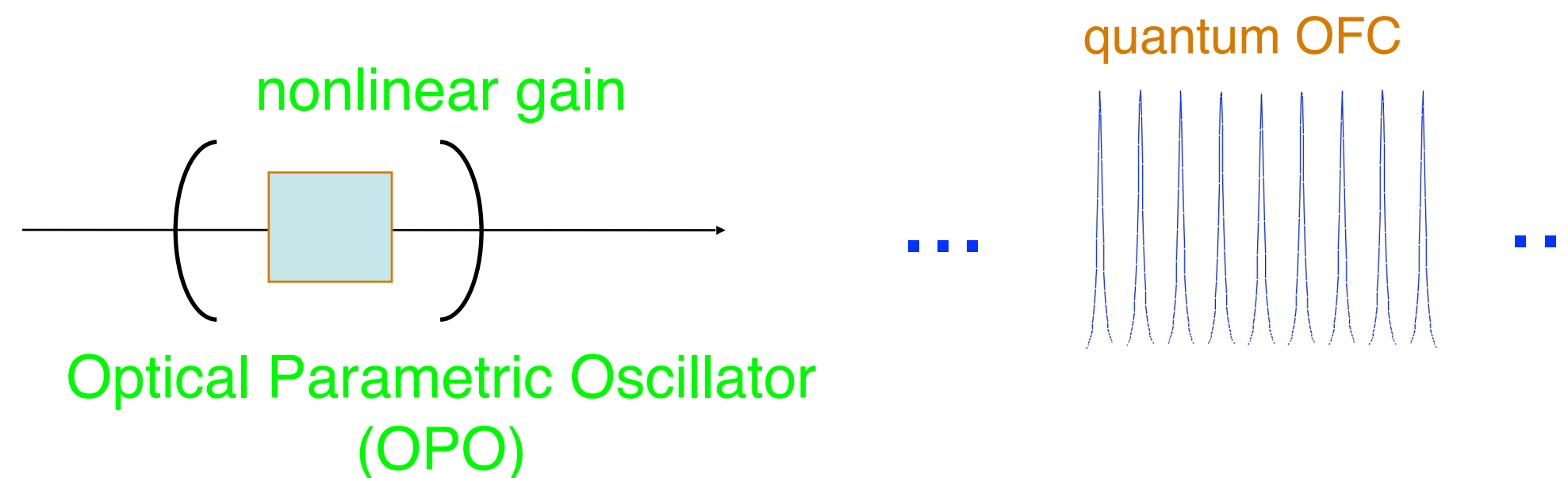


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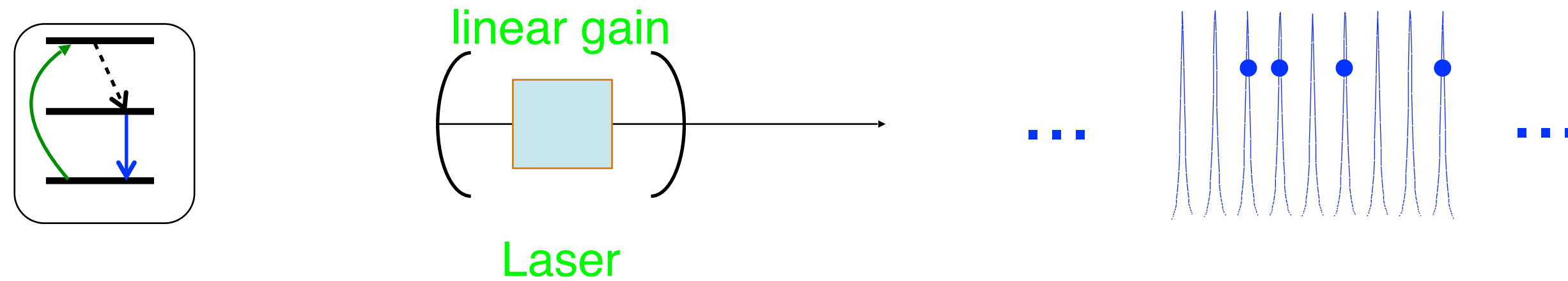
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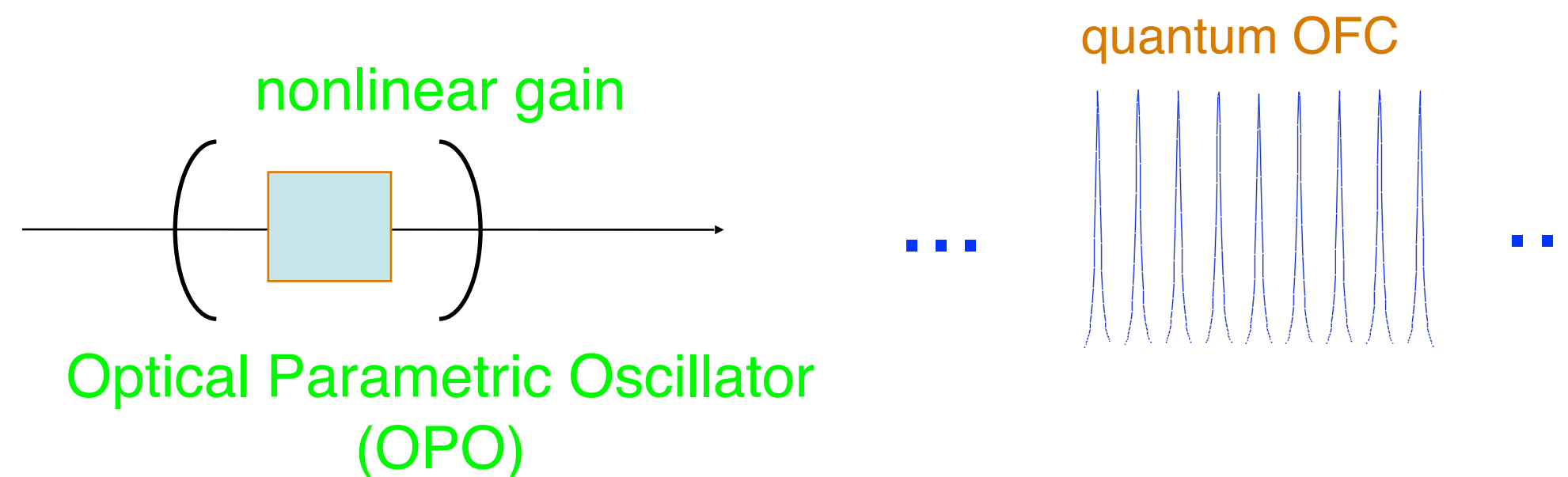


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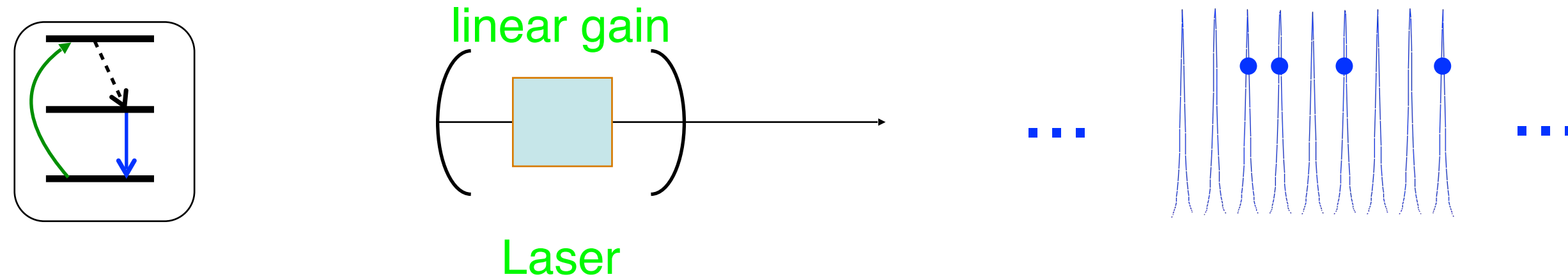
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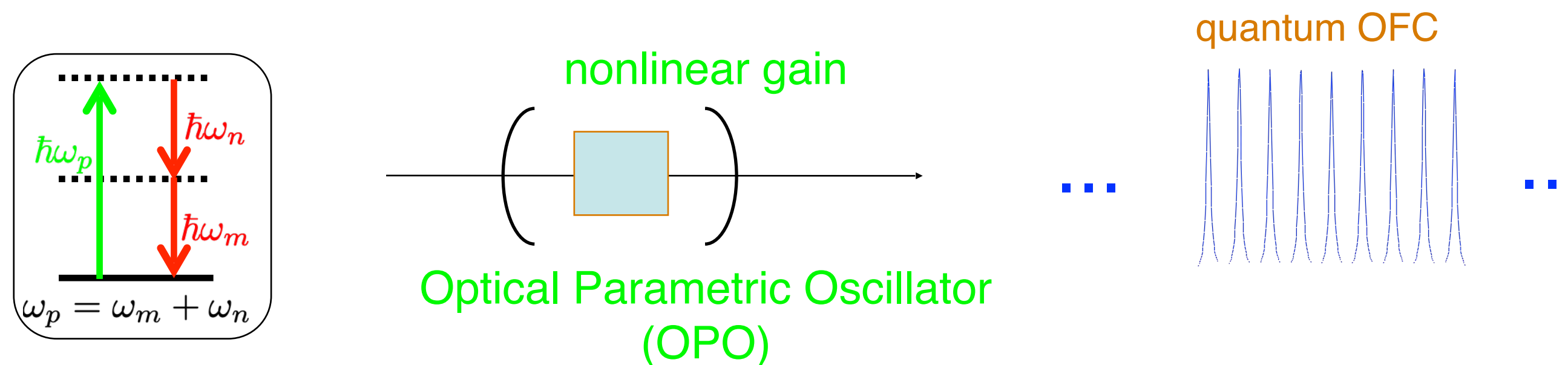


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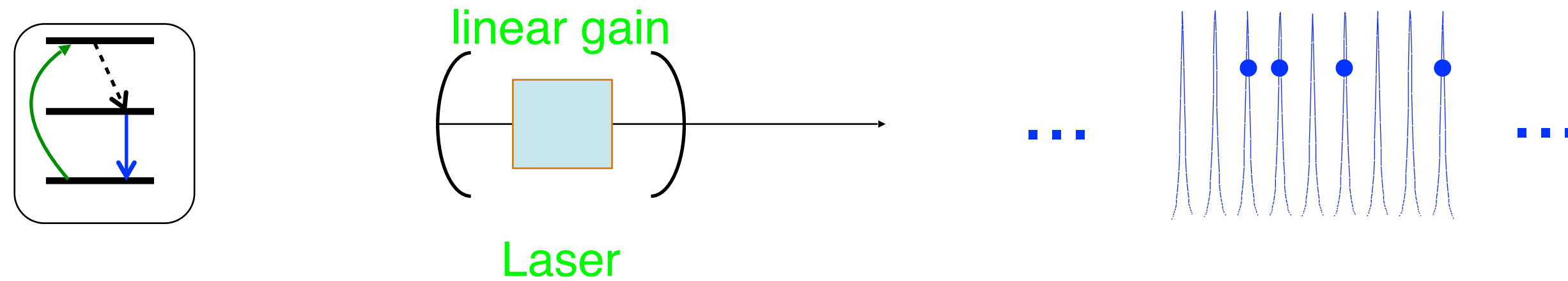
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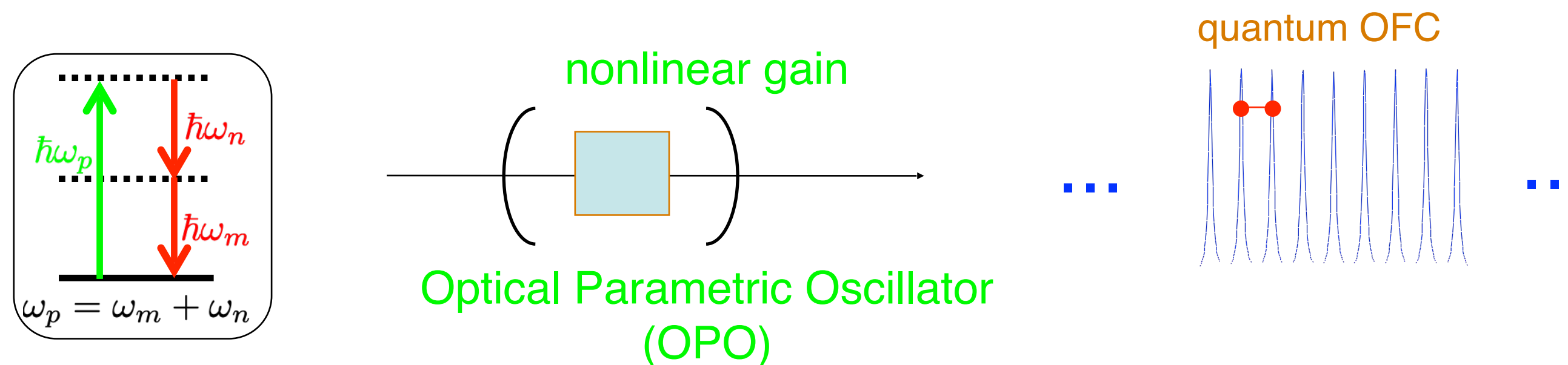


John L. Hall



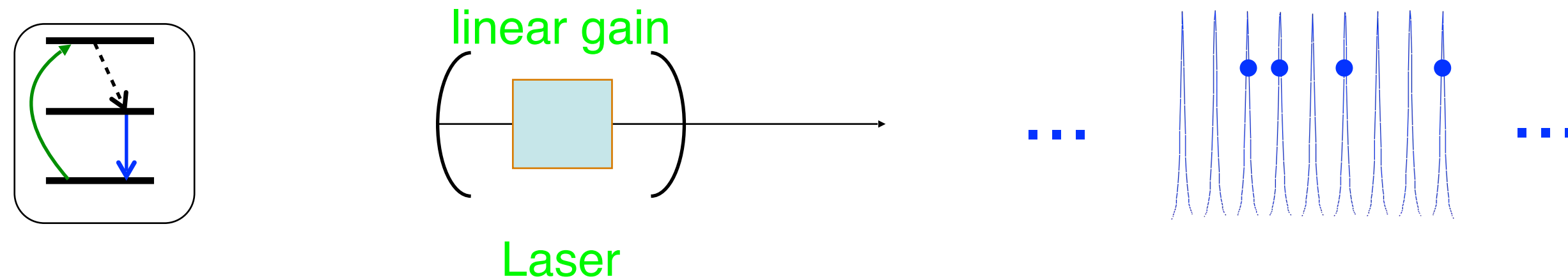
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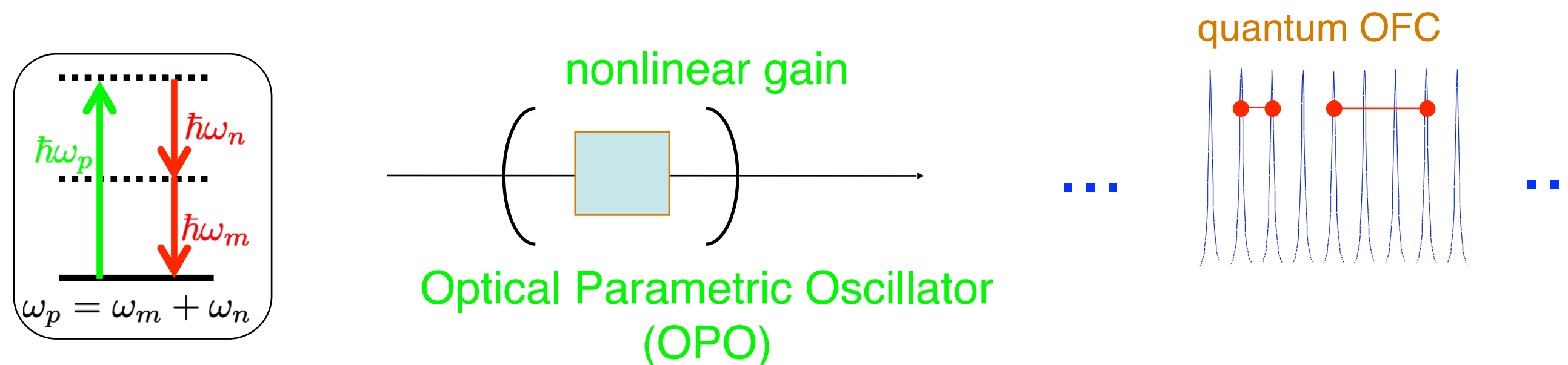


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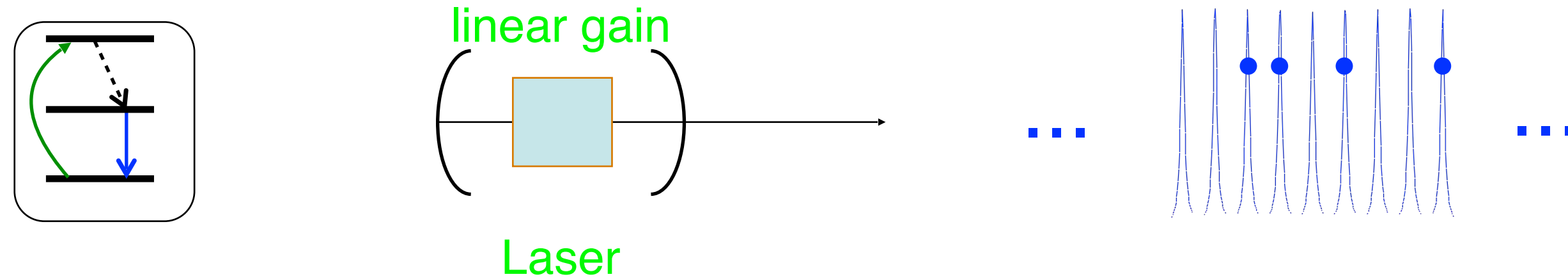
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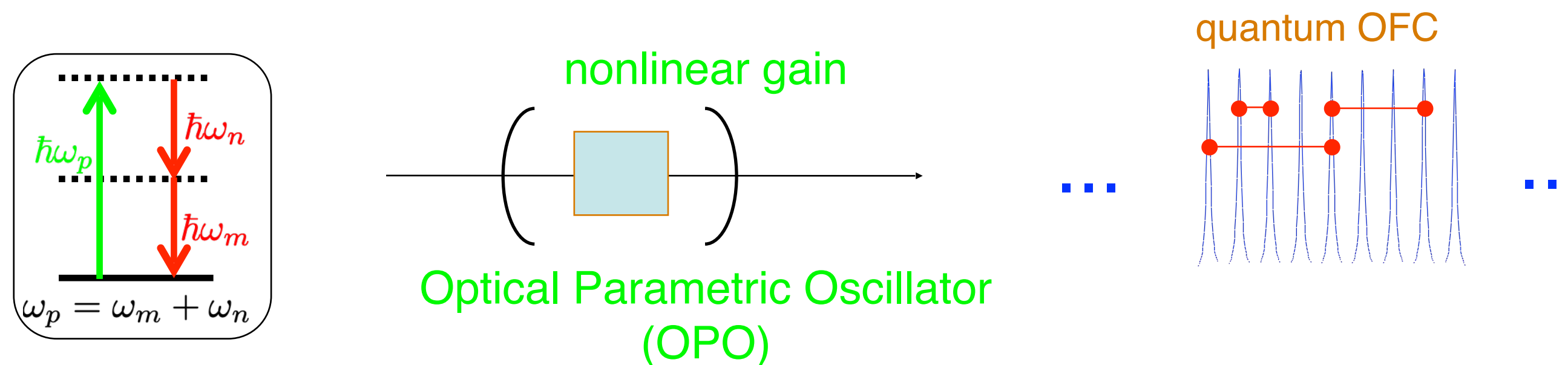


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Theodor W. Hänsch

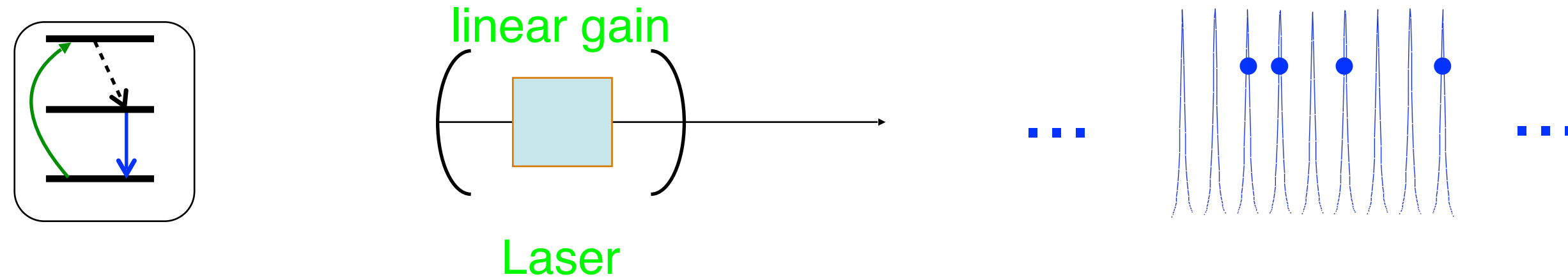
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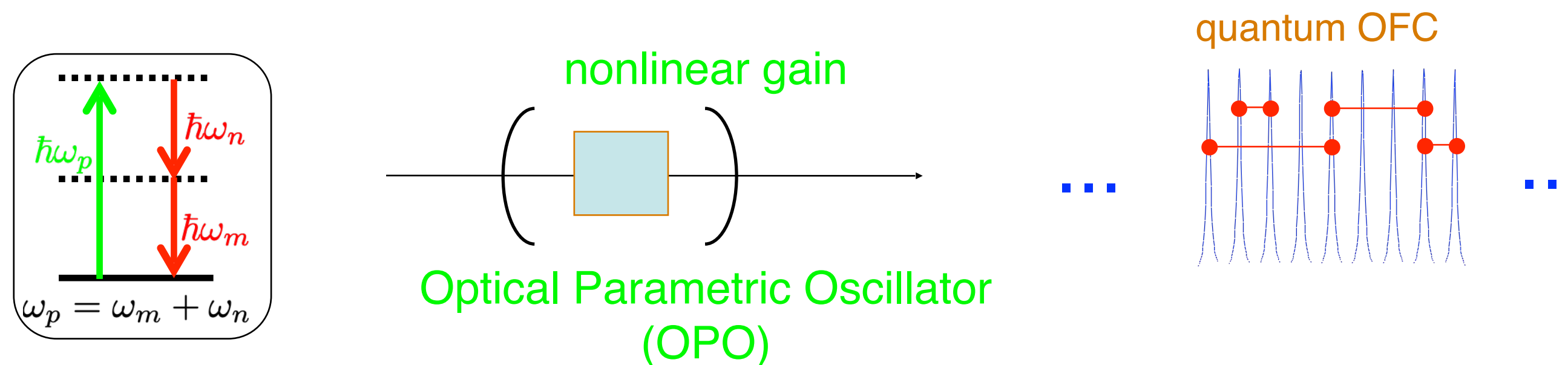


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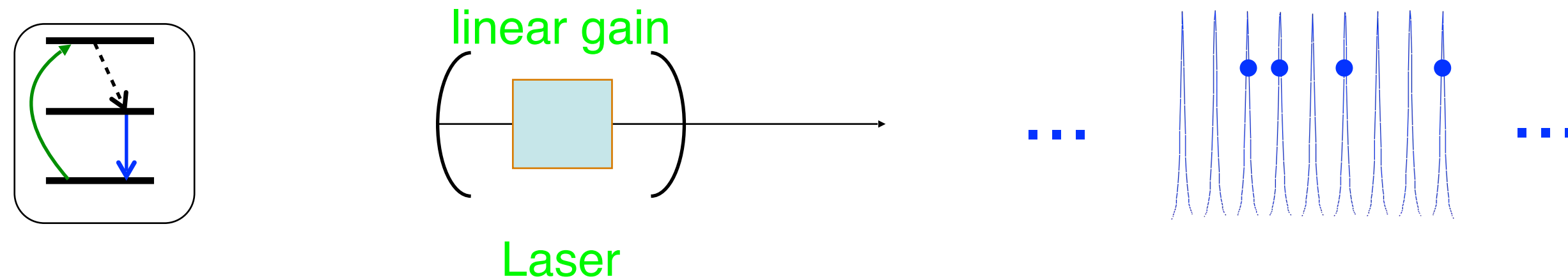
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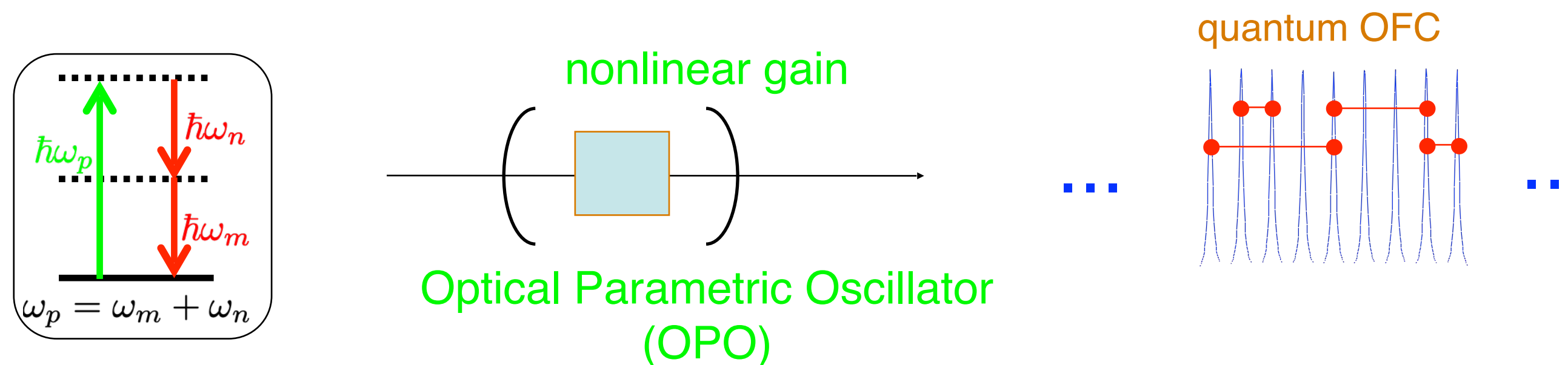


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## Why not turn the QOFC into a quantum computer?



**Multipartite cluster entanglement  
in one fell swoop:  
a top-down, large-scale quantum register**

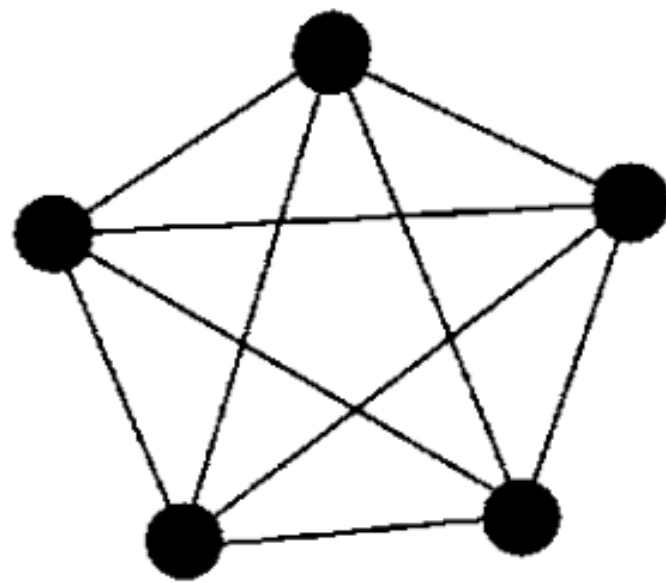


# Multipartite continuous-variable entanglement from concurrent nonlinearities

The figure consists of two schematic diagrams, (a) and (b), illustrating the energy levels and transitions in an optical parametric oscillator (OPO) based on a nonlinear crystal.

**(a) Singly resonant OPO:** This diagram shows a nonlinear crystal with two input/output ports. The left port has two input beams: a pump beam at frequency  $2\omega_0$  (represented by a black dot) and a signal beam at frequency  $\omega_0 + \omega_1$  (represented by a white circle). The right port has a single output beam at frequency  $\omega_1$  (represented by a white circle). The energy levels are shown as vertical lines with transitions indicated by arrows. The pump beam transitions from the ground state to the  $\omega_0$  level, and the signal beam transitions from the  $\omega_0$  level to the  $\omega_0 + \omega_1$  level. The output beam transitions from the  $\omega_0 + \omega_1$  level to the  $\omega_1$  level.

**(b) Doubly resonant OPO:** This diagram shows a nonlinear crystal with four input/output ports. The left port has two input beams: a pump beam at frequency  $\omega_0 + \omega_1$  (represented by a black dot) and a signal beam at frequency  $\omega_1 + \omega_2$  (represented by a white circle). The right port has two output beams: a signal beam at frequency  $2\omega_1$  (represented by a white circle) and a pump beam at frequency  $\omega_2 + \omega_3$  (represented by a black dot). The energy levels are shown as vertical lines with transitions indicated by arrows. The pump beam transitions from the ground state to the  $\omega_0$  level, and the signal beam transitions from the  $\omega_0$  level to the  $\omega_0 + \omega_1$  level. The output beam transitions from the  $\omega_0 + \omega_1$  level to the  $\omega_1$  level. The pump beam transitions from the  $\omega_1$  level to the  $\omega_1 + \omega_2$  level, and the signal beam transitions from the  $\omega_1 + \omega_2$  level to the  $\omega_2$  level.



## Bright tripartite entanglement in triply concurrent parametric oscillation

*ARC Centre of Excellence for Quantum-Atom Optics, School of Physical Sciences,  
University of Queensland, Brisbane, Queensland 4072, Australia*

QSS#13 07/16/2020





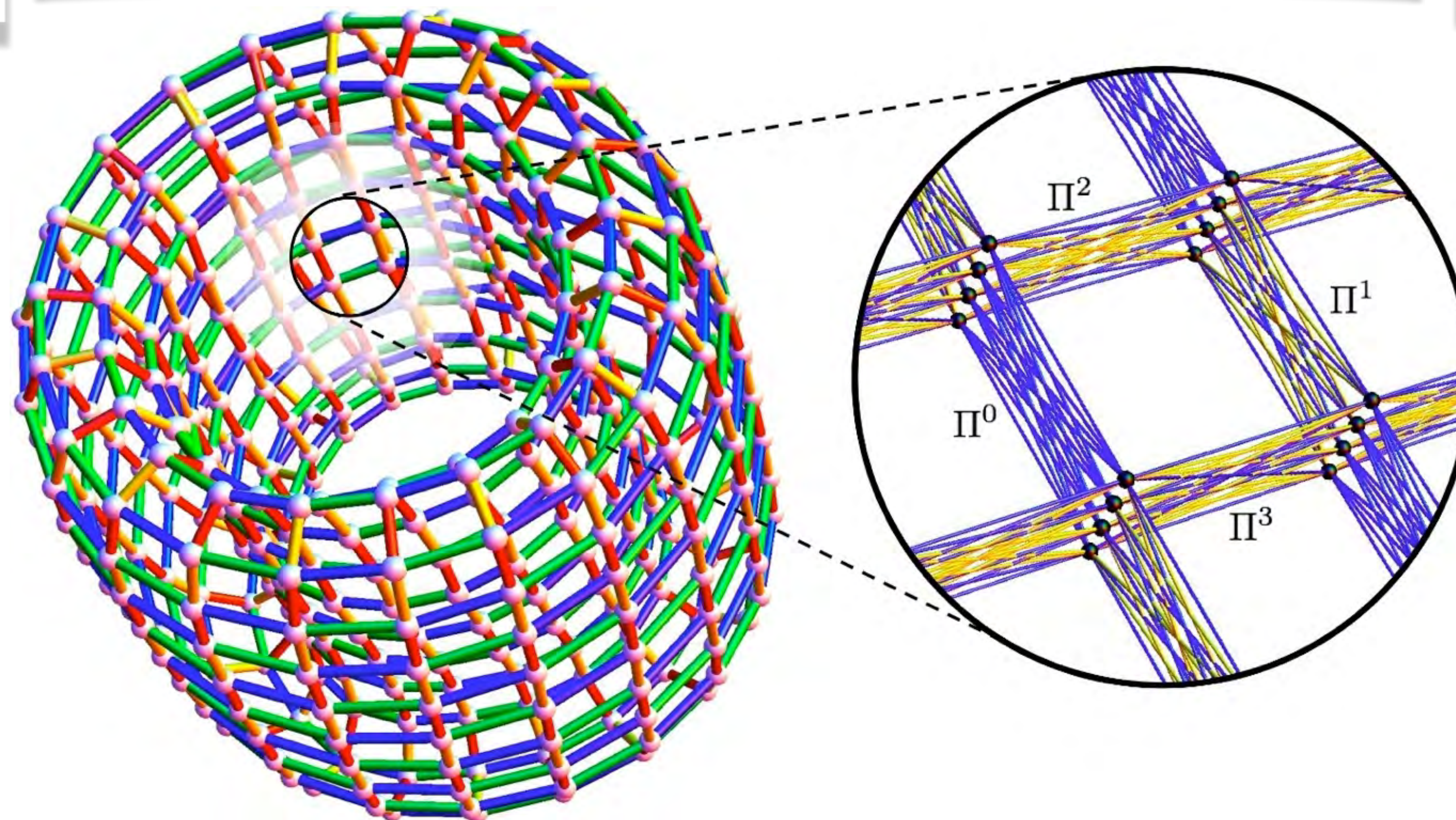
PRL 101, 130501 (2008)

Selected for a [Viewpoint](#) in *Physics*  
PHYSICAL REVIEW LETTERS

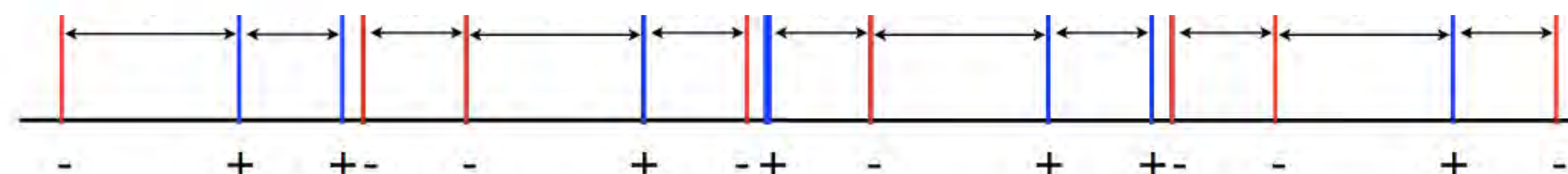
week ending  
26 SEPTEMBER 2008

## One-Way Quantum Computing in the Optical Frequency Comb

Nicolas C. Menicucci,<sup>1,2</sup> Steven T. Flammia,<sup>3</sup> and Olivier Pfister<sup>4</sup>



15 pump modes into YZY, ZZZ, ZYY







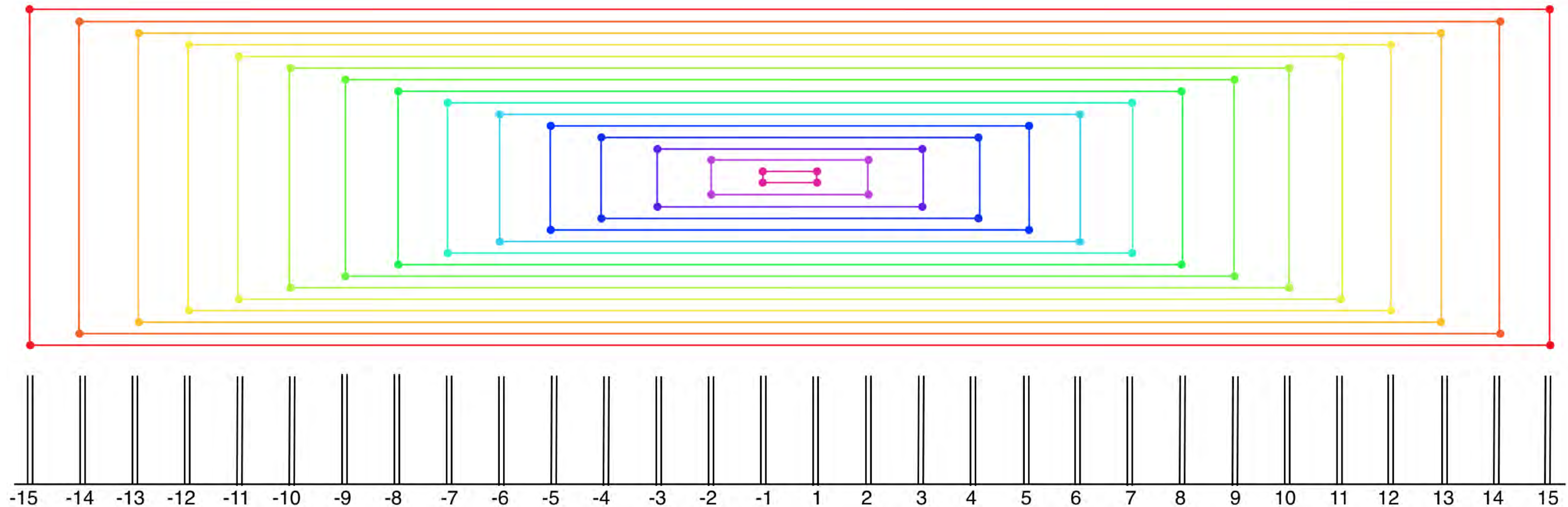
PRL 107, 030505 (2011)

PHYSICAL REVIEW LETTERS

week ending  
15 JULY 2011

## Parallel Generation of Quadripartite Cluster Entanglement in the Optical Frequency Comb

Matthew Pysher,<sup>1</sup> Yoshichika Miwa,<sup>2</sup> Reihaneh Shahrokhshahi,<sup>1</sup> Russell Bloomer,<sup>1</sup> and Olivier Pfister<sup>1,\*</sup>



### Entanglement gets scaled up in an optical frequency comb

Using a single nonlinear optical element, researchers have entangled dozens of the comb's optical modes.





PRL **112**, 120505 (2014)

PHYSICAL REVIEW LETTERS

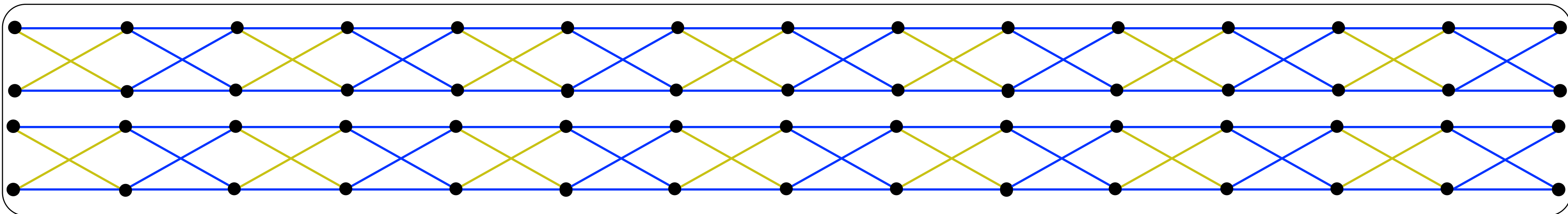
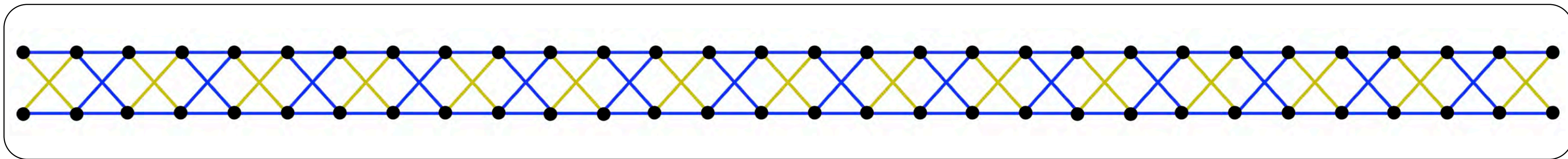
week ending  
28 MARCH 2014

## Experimental Realization of Multipartite Entanglement of 60 Modes of a Quantum Optical Frequency Comb

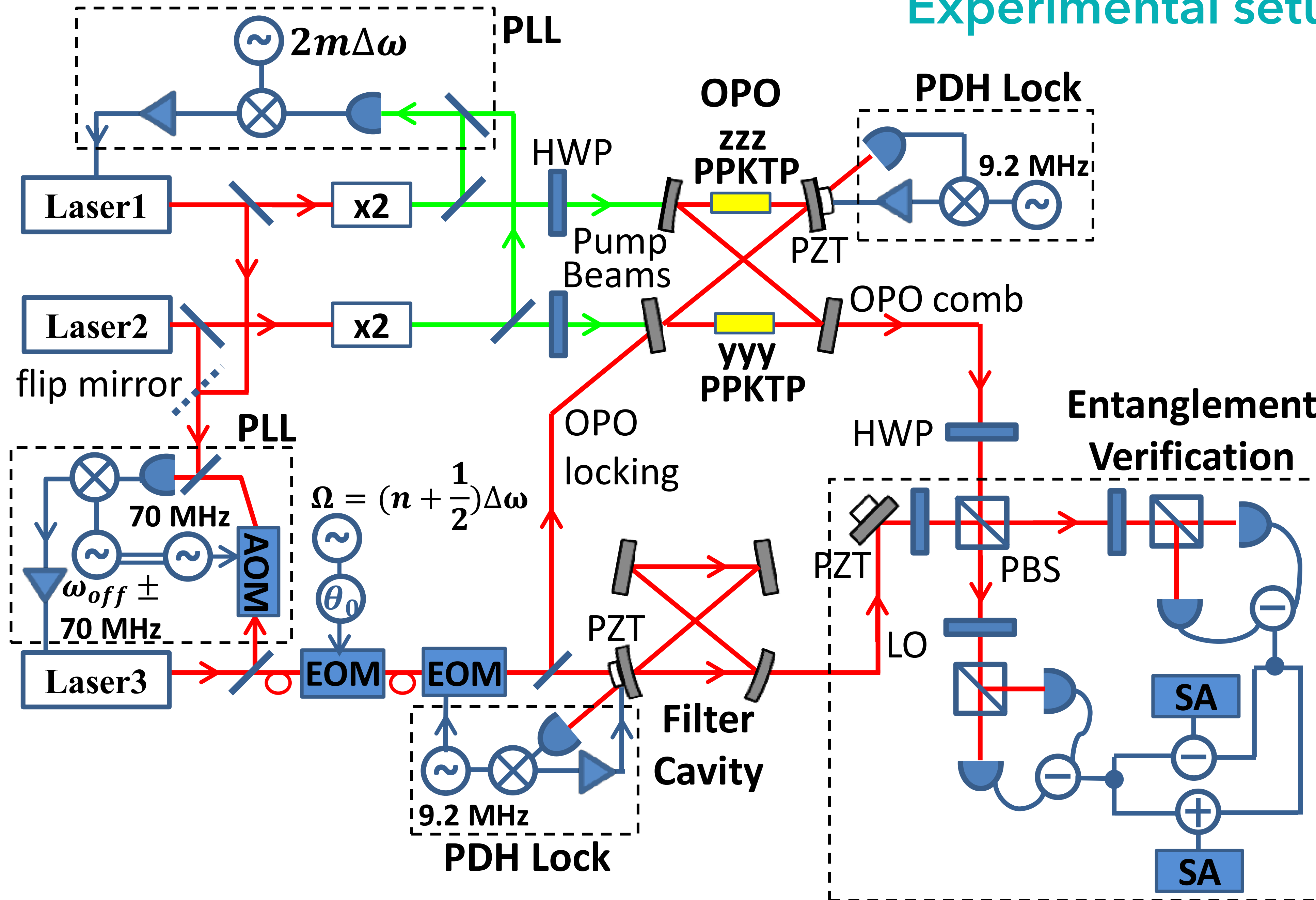
Moran Chen,<sup>1</sup> Nicolas C. Menicucci,<sup>2,\*</sup> and Olivier Pfister<sup>1,†</sup>

<sup>1</sup>*Department of Physics, University of Virginia, Charlottesville, Virginia 22903, USA*

<sup>2</sup>*School of Physics, The University of Sydney, Sydney, New South Wales 2006, Australia*



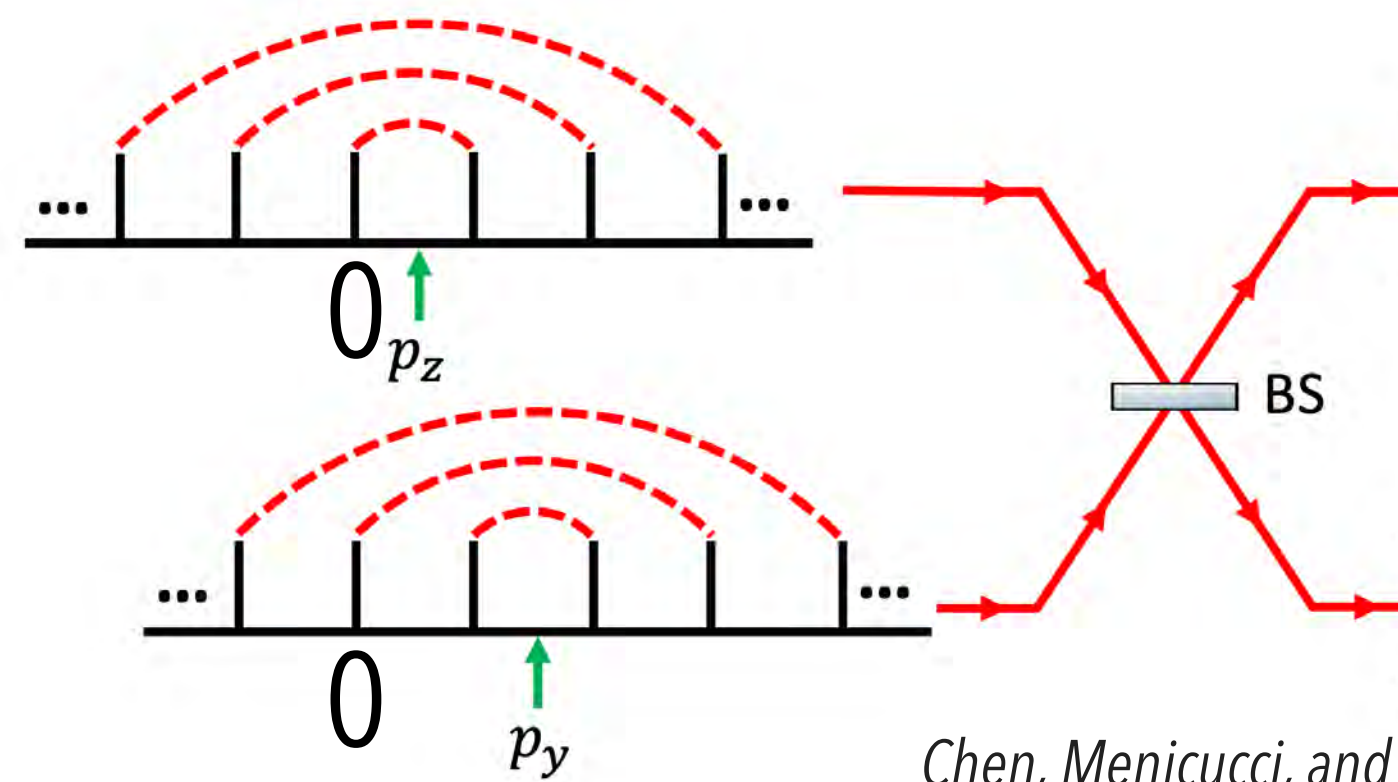
# Experimental setup





# Interfering quantum combs

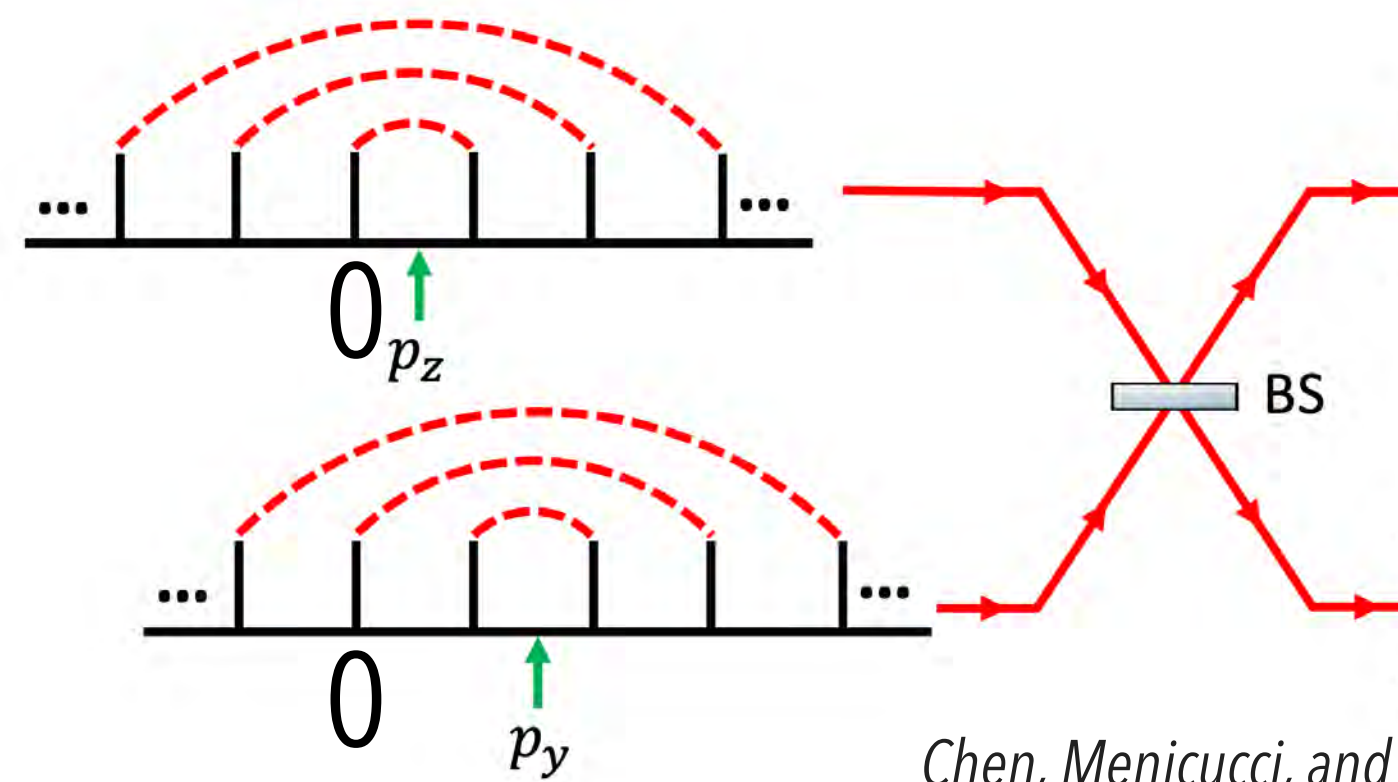
## Interference between identical frequencies of different combs



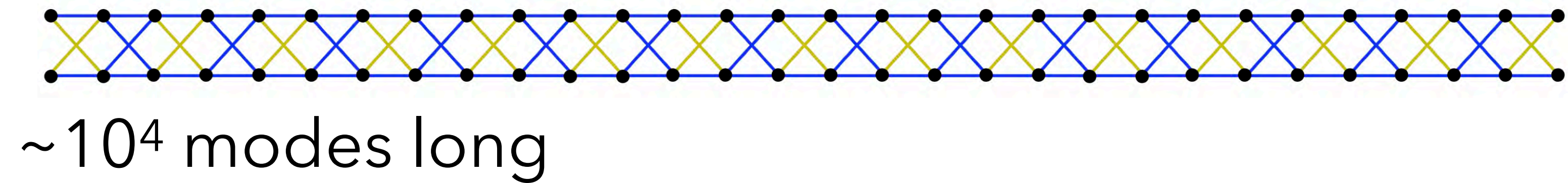
Chen, Menicucci, and Pfister, **Experimental realization of multipartite entanglement of 60 modes of a quantum optical frequency comb**, PRL 112, 120505 (2014)

# Interfering quantum combs

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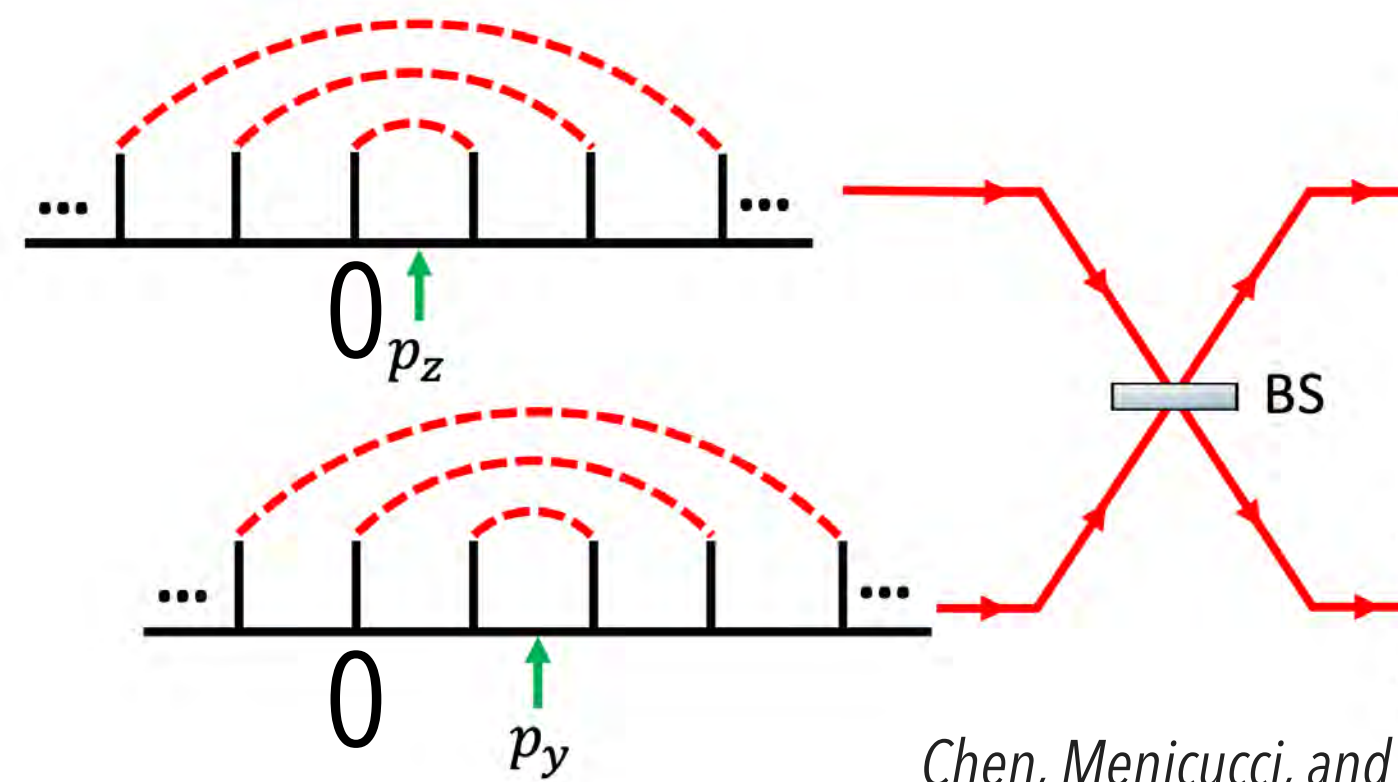


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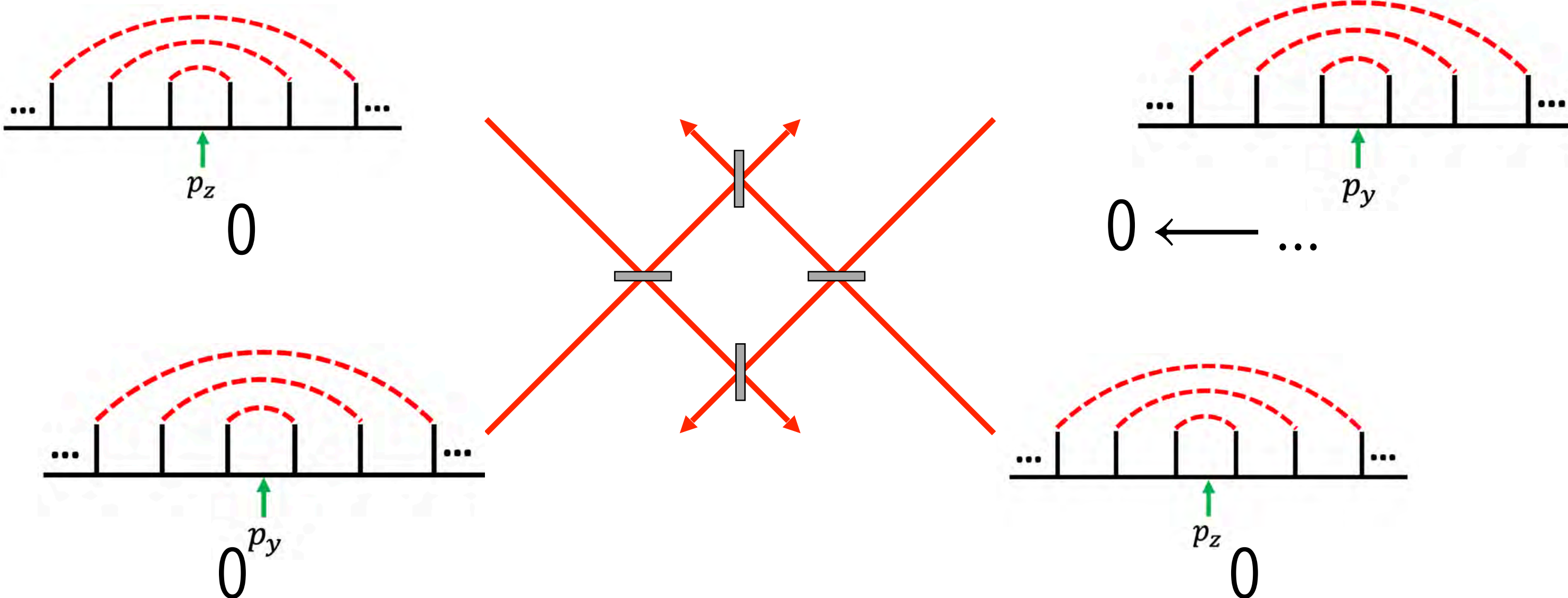
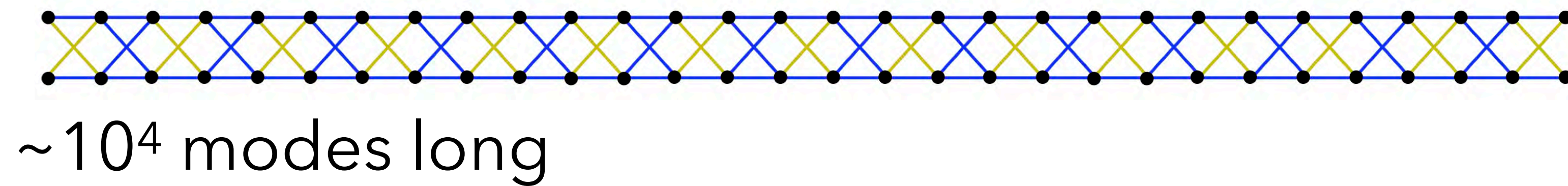


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Chen, Menicucci, and Pfister, **Experimental realization of multipartite entanglement of 60 modes of a quantum optical frequency comb**, PRL 112, 120505 (2014)

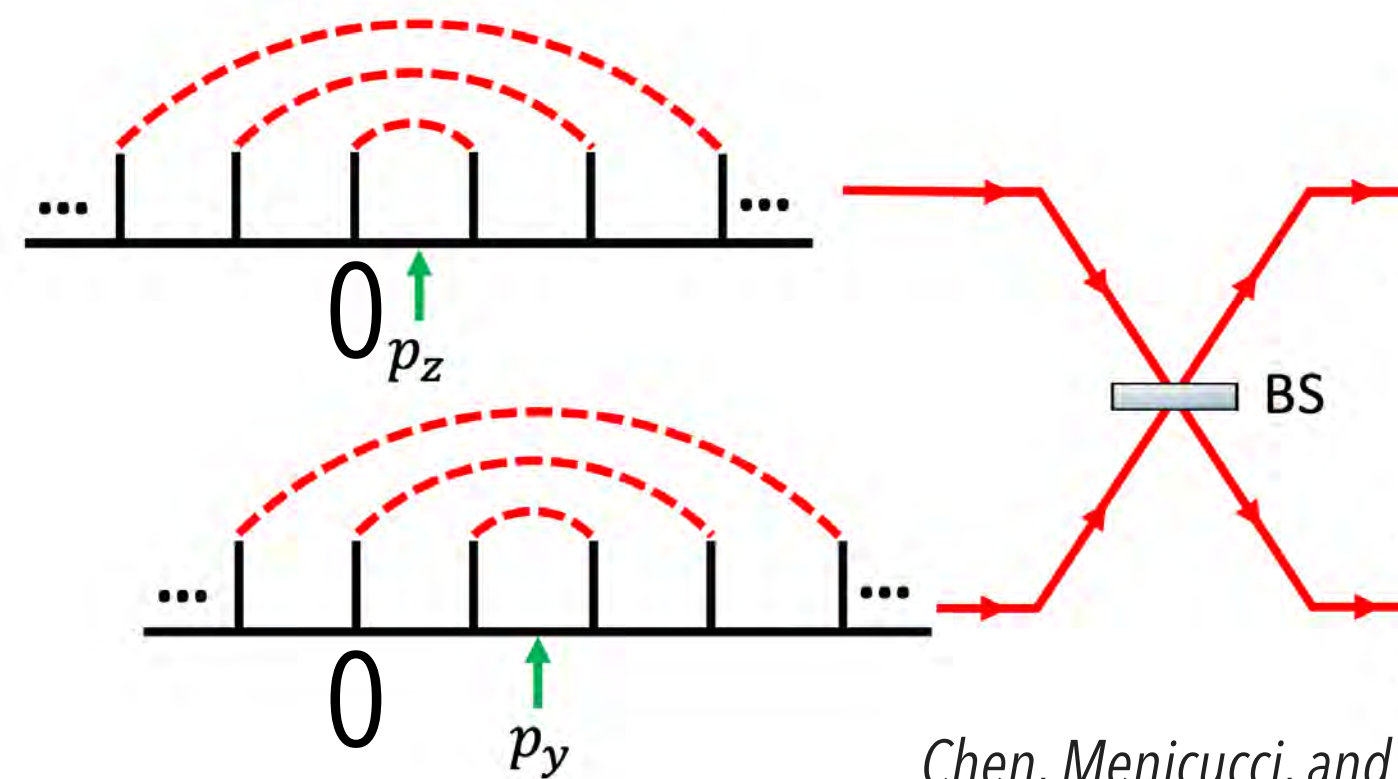


Wang, Chen, Menicucci, and Pfister, **Weaving quantum optical frequency combs into hypercubic cluster states**, PRA 90, 032325 (2014)

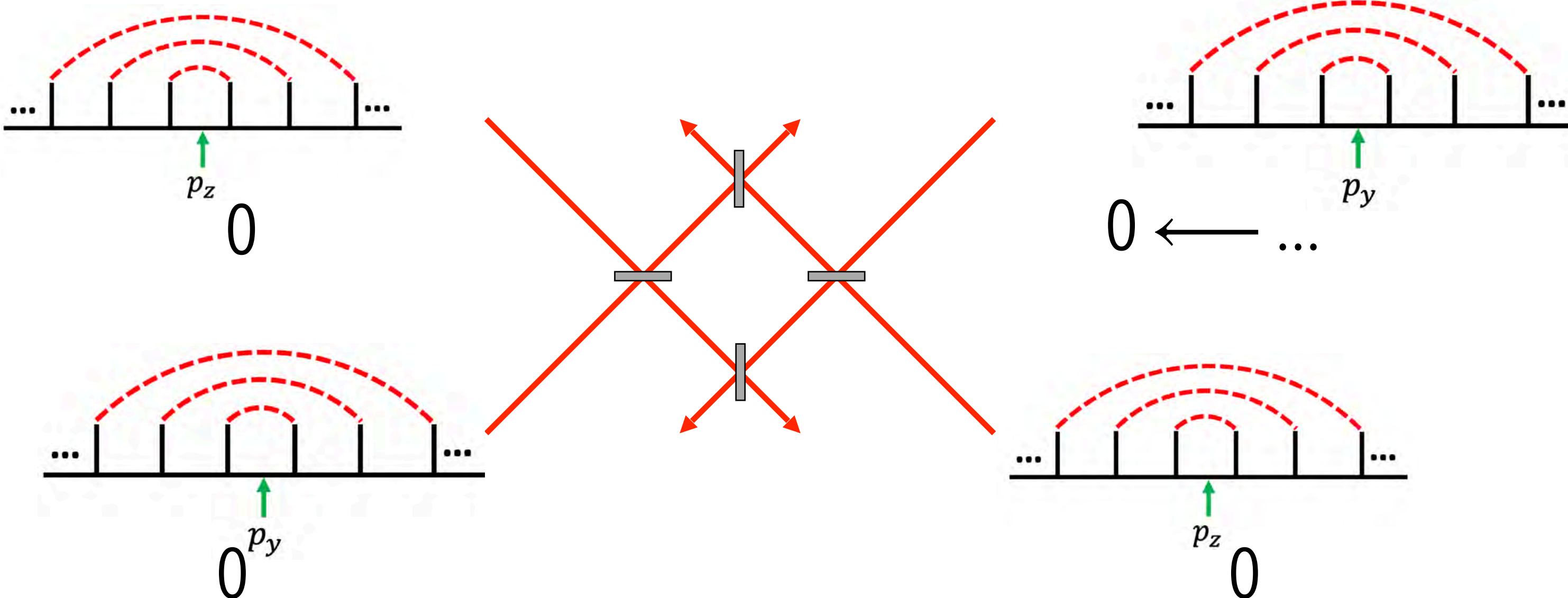
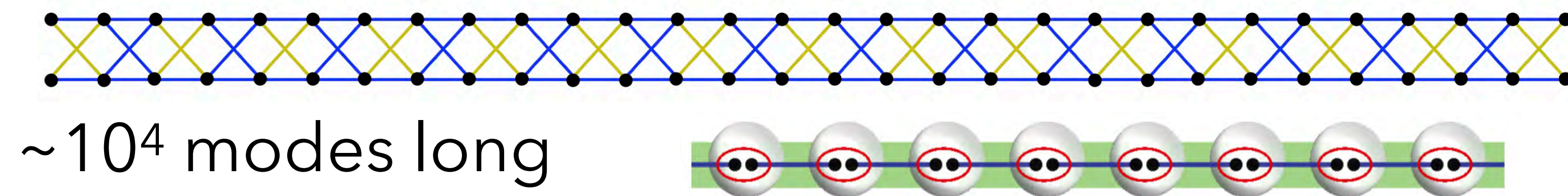


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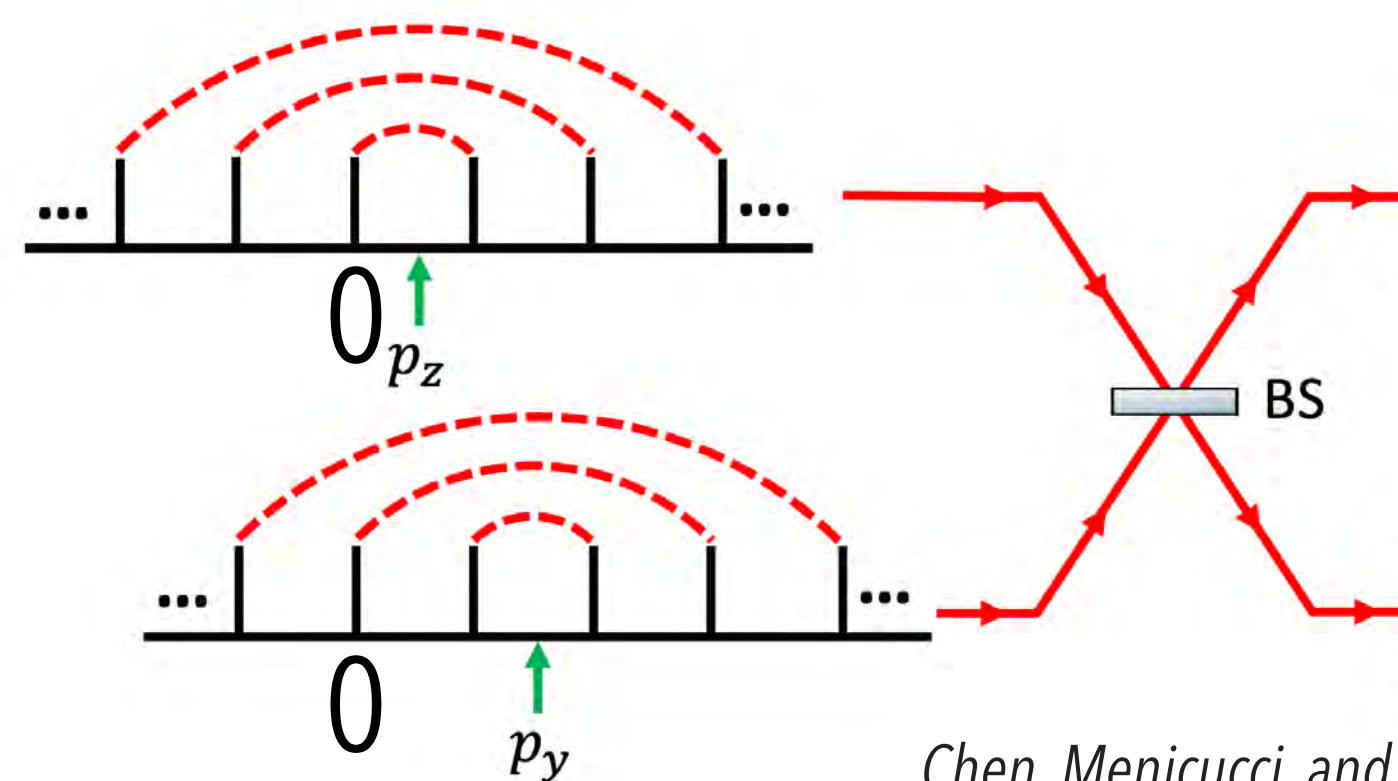


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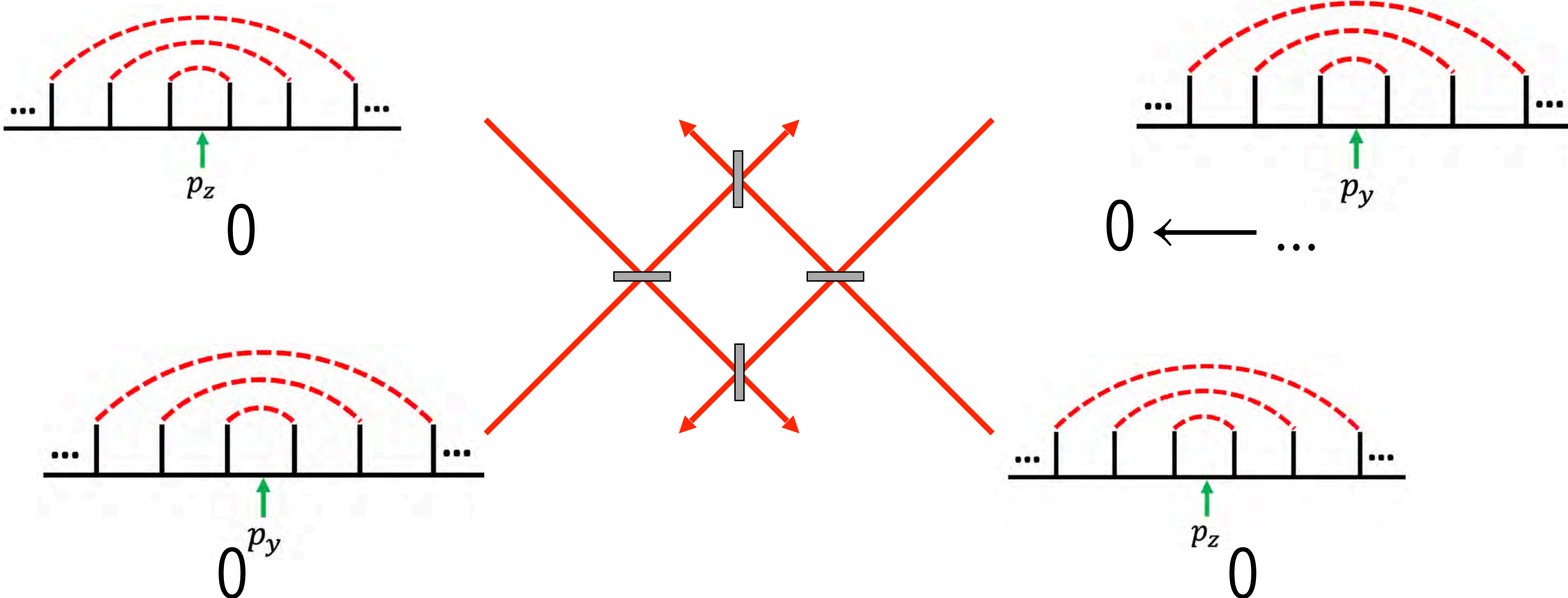
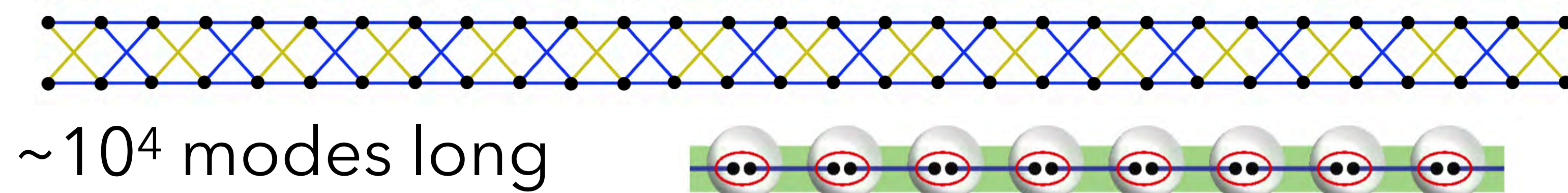


# Interfering quantum combs

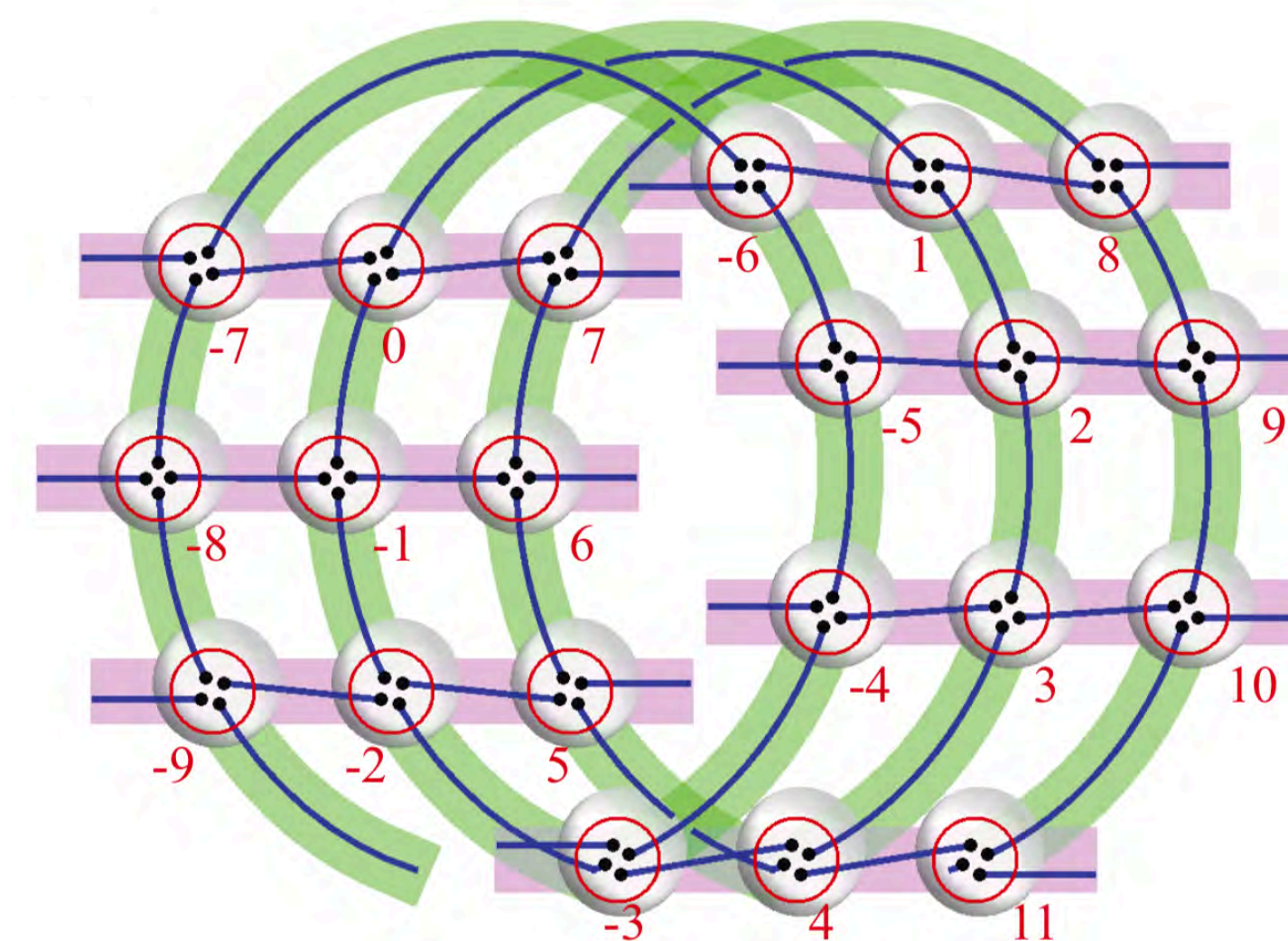
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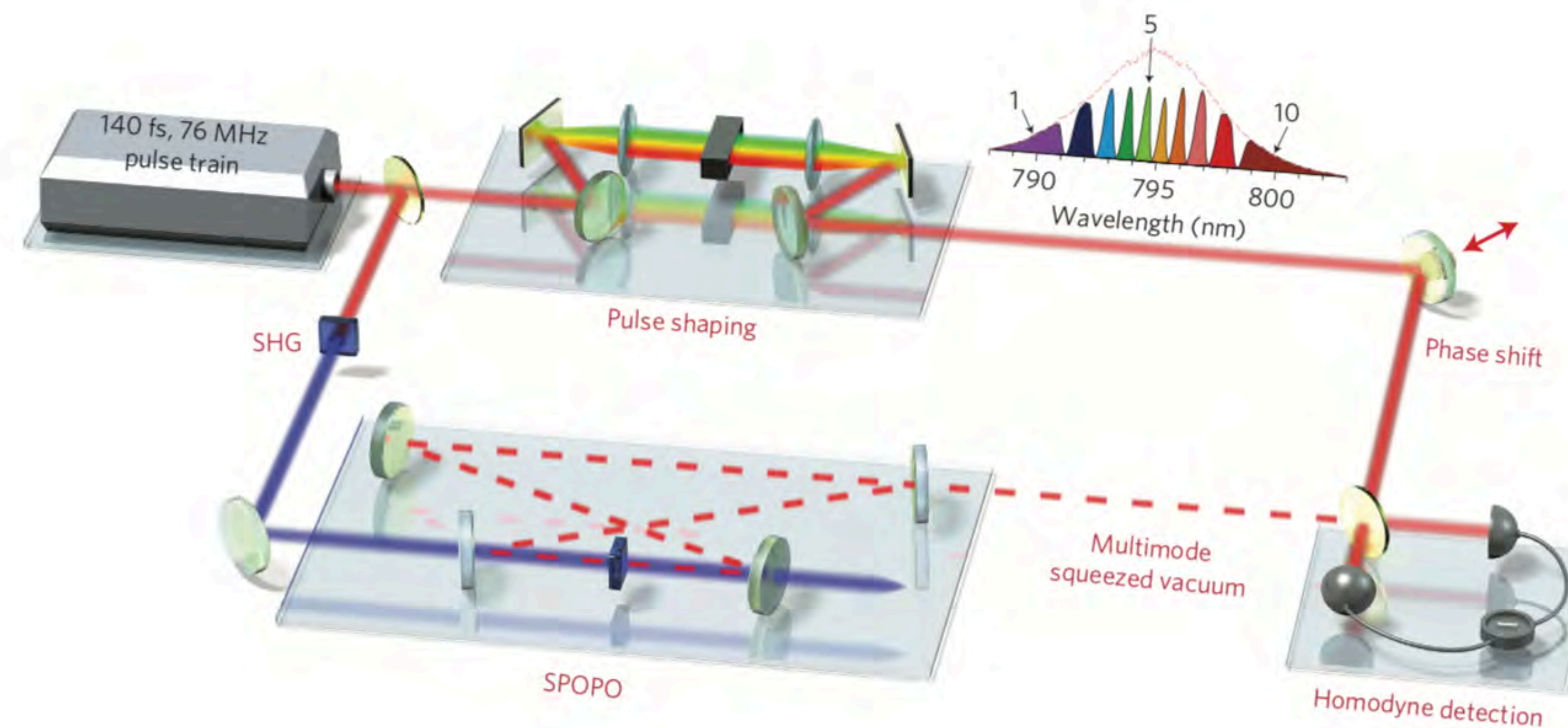
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# Wavelength-multiplexed quantum networks with ultrafast frequency combs

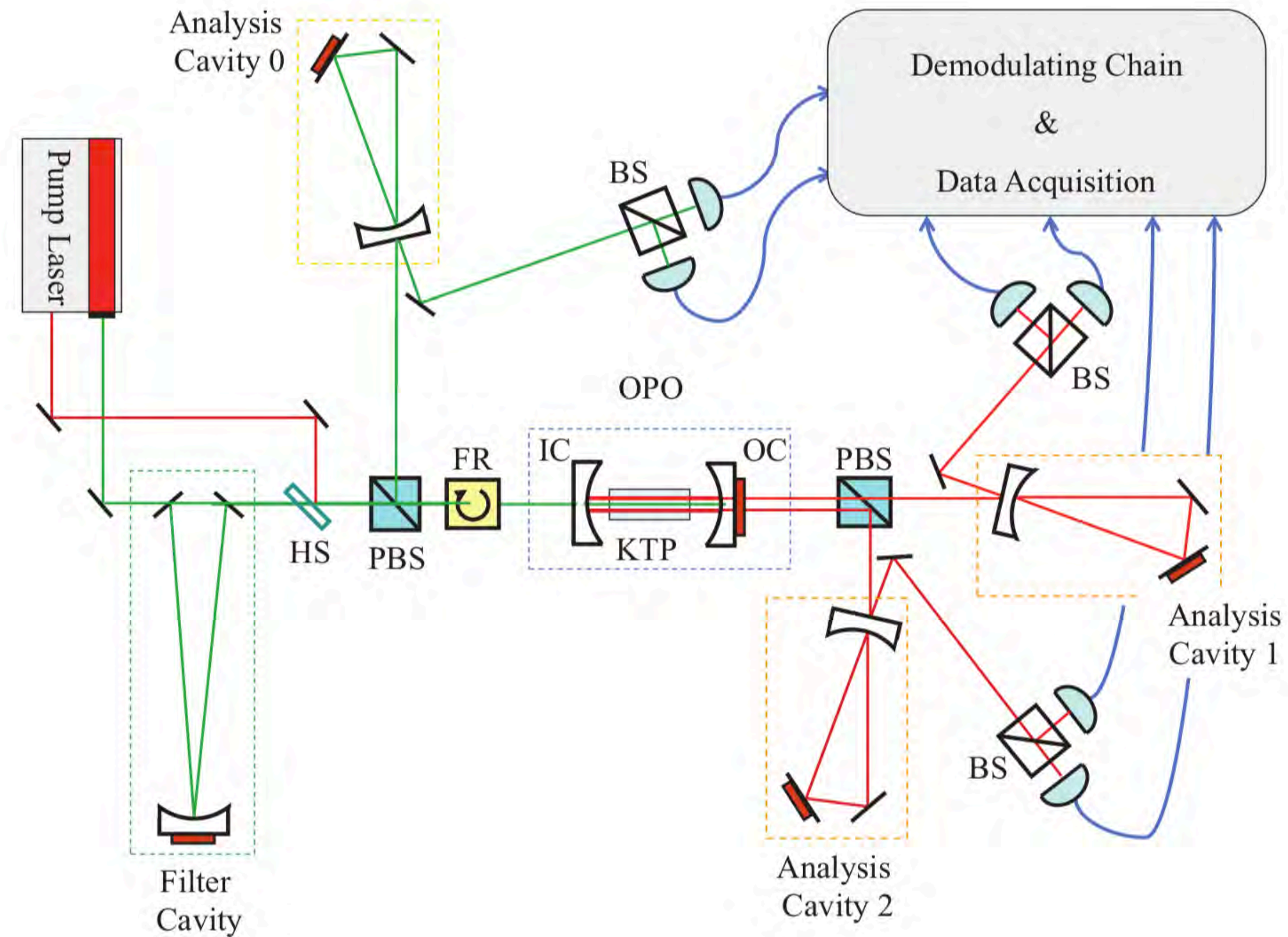
Jonathan Roslund, Renné Medeiros de Araújo, Shifeng Jiang, Claude Fabre and Nicolas Treps\*





# Hexapartite Entanglement in an above-Threshold Optical Parametric Oscillator

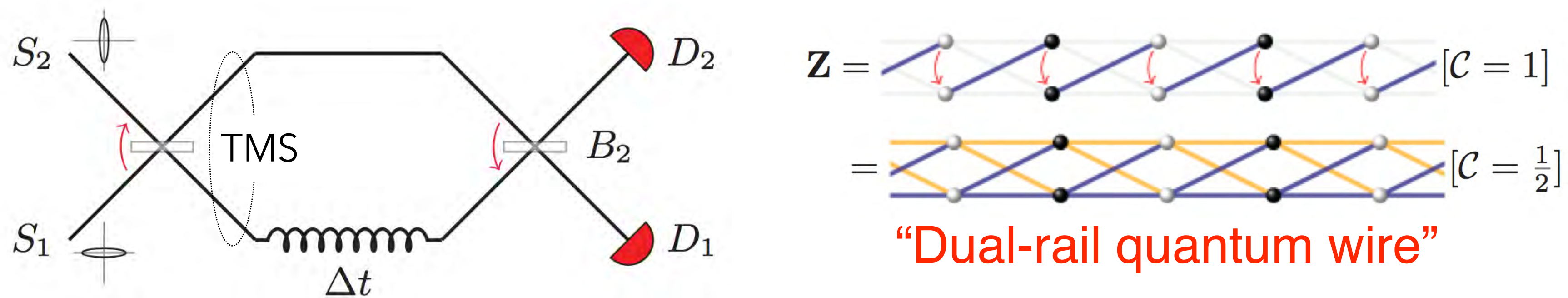
F. A. S. Barbosa,<sup>1</sup> A. S. Coelho,<sup>2,3</sup> L. F. Muñoz-Martínez,<sup>4</sup> L. Ortiz-Gutiérrez,<sup>5</sup>  
A. S. Villar,<sup>6</sup> P. Nussenzveig,<sup>7</sup> and M. Martinelli<sup>7,\*</sup>





## Temporal-mode continuous-variable cluster states using linear optics

Nicolas C. Menicucci



- Akira Furusawa's group: sequential entanglement of  $10^4$  qumodes (2 at a time)

### LETTERS

PUBLISHED ONLINE: 17 NOVEMBER 2013 | DOI: 10.1038/NPHOTON.2013.287

nature  
photonics

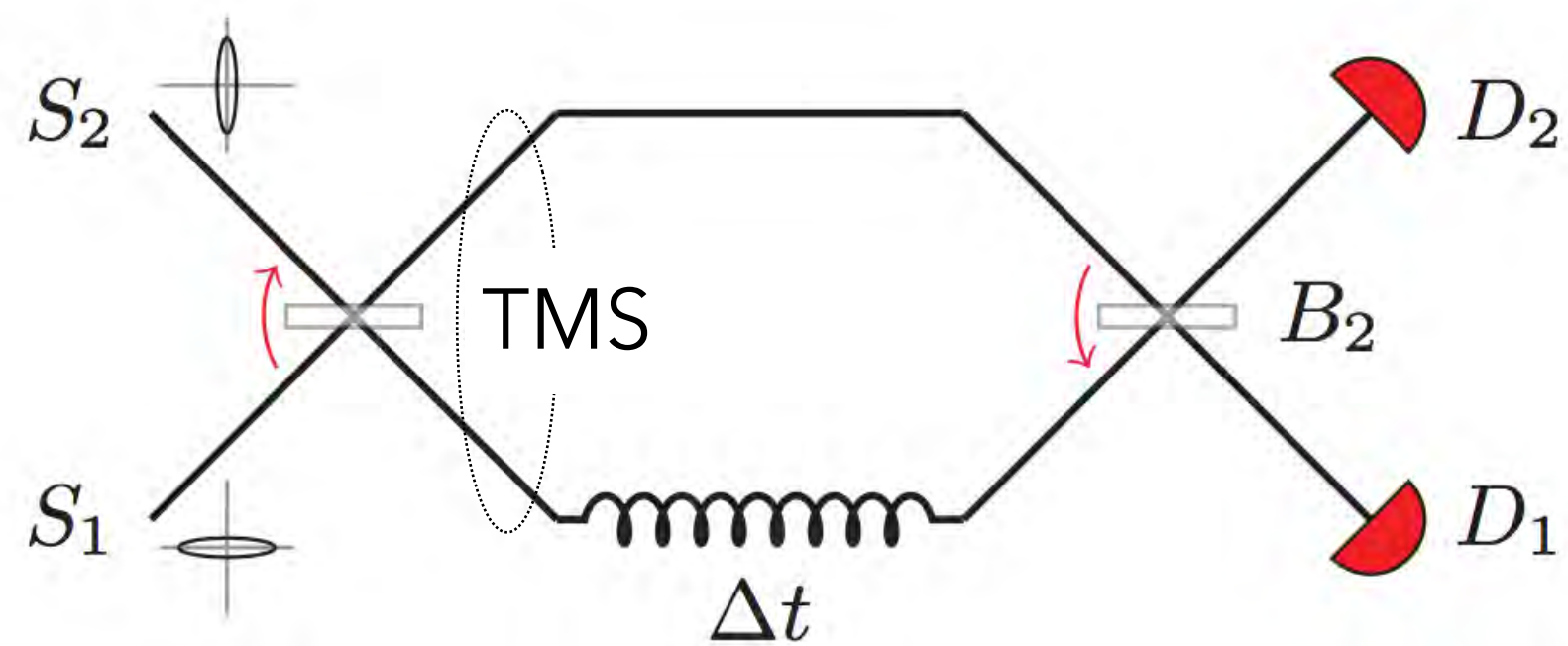
## Ultra-large-scale continuous-variable cluster states multiplexed in the time domain

Shota Yokoyama<sup>1</sup>, Ryuji Ukai<sup>1</sup>, Seiji C. Armstrong<sup>1,2</sup>, Chanond Sornphiphatphong<sup>1</sup>, Toshiyuki Kaji<sup>1</sup>, Shigenari Suzuki<sup>1</sup>, Jun-ichi Yoshikawa<sup>1</sup>, Hidehiro Yonezawa<sup>1</sup>, Nicolas C. Menicucci<sup>3</sup> and Akira Furusawa<sup>1\*</sup>



## Temporal-mode continuous-variable cluster states using linear optics

Nicolas C. Menicucci



$$\mathbf{Z} = \begin{array}{c} \text{[Diagram of a single-rail cluster state with } C=1 \text{]} \\ \text{[Diagram of a dual-rail cluster state with } C=\frac{1}{2} \text{]} \end{array}$$

“Dual-rail quantum wire”

~~$10^6$~~

- Akira Furusawa's group: sequential entanglement of  $10^4$  qumodes (2 at a time)

## Invited Article: Generation of one-million-mode continuous-variable cluster state by unlimited time-domain multiplexing

Jun-ichi Yoshikawa,<sup>1</sup> Shota Yokoyama,<sup>1,2</sup> Toshiyuki Kaji,<sup>1</sup>  
Chanond Sornphiphatphong,<sup>1</sup> Yu Shiozawa,<sup>1</sup> Kenzo Makino,<sup>1</sup>  
and Akira Furusawa<sup>1,a</sup>



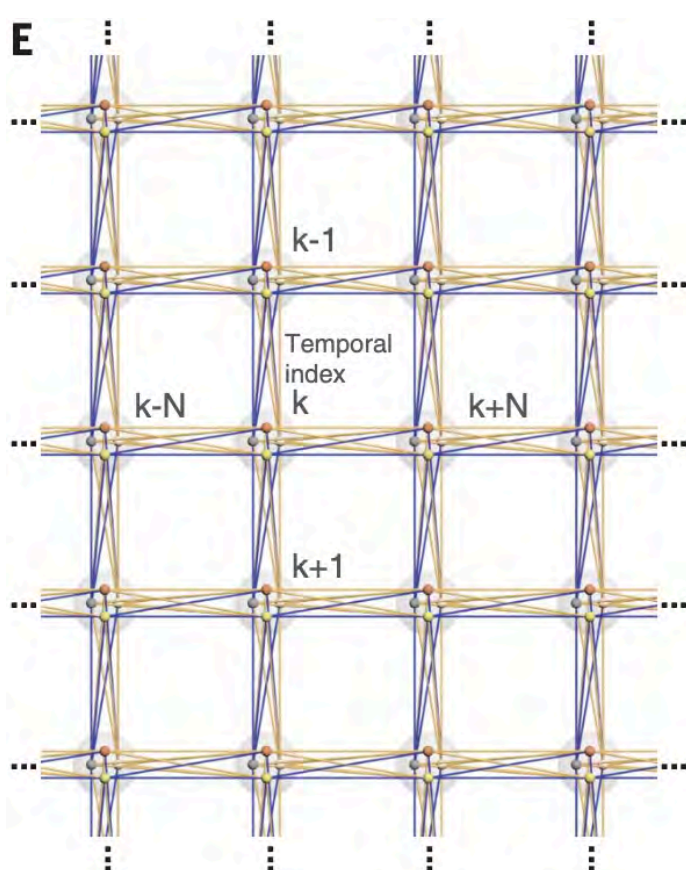
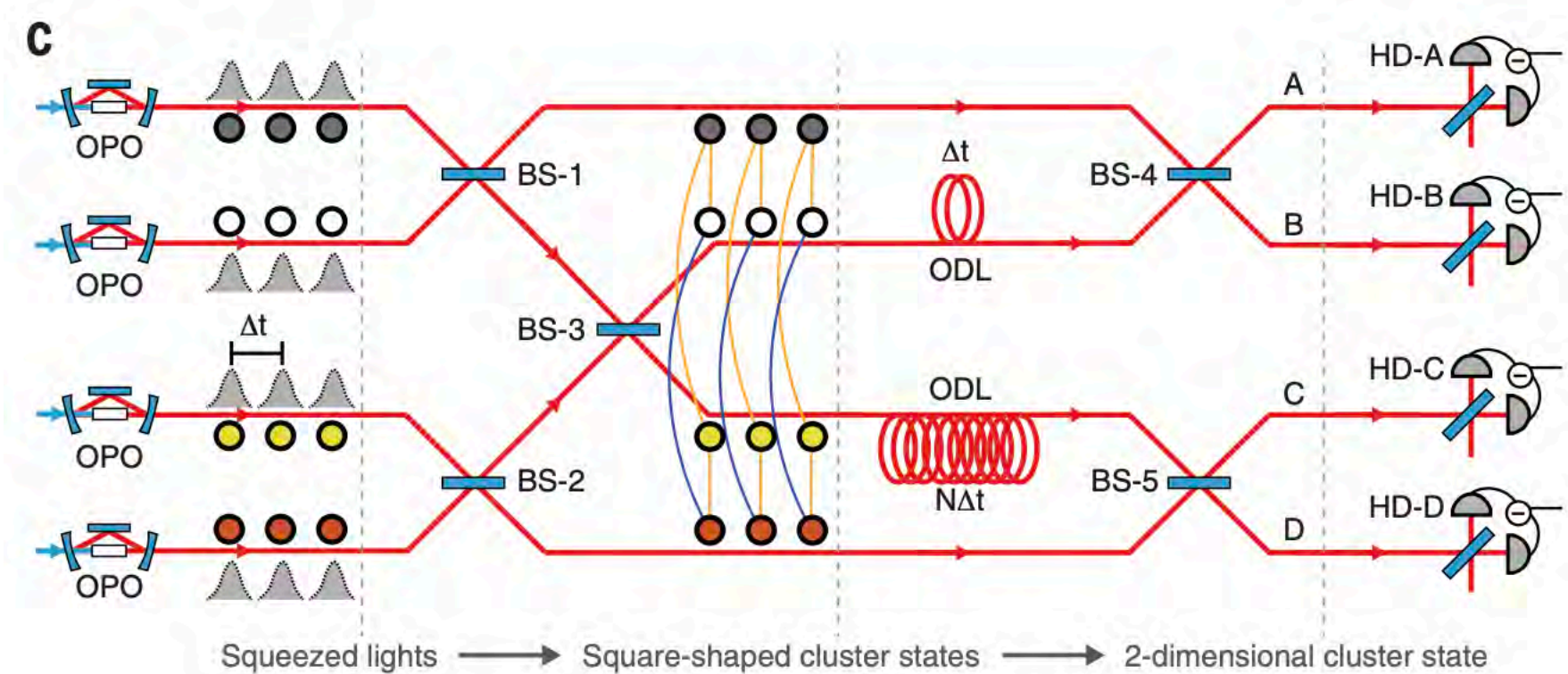
# 2D cluster states have arrived

## QUANTUM COMPUTING

### Generation of time-domain-multiplexed two-dimensional cluster state

Warit Asavanant<sup>1</sup>, Yu Shiozawa<sup>1</sup>, Shota Yokoyama<sup>2</sup>, Baramée Charoensombutamon<sup>1</sup>, Hiroki Emura<sup>1</sup>, Rafael N. Alexander<sup>3</sup>, Shuntaro Takeda<sup>1,4</sup>, Jun-ichi Yoshikawa<sup>1</sup>, Nicolas C. Menicucci<sup>5</sup>, Hidehiro Yonezawa<sup>2</sup>, Akira Furusawa<sup>1\*</sup>

Asavanant *et al.*, *Science* **366**, 373–376 (2019) 18 October 2019

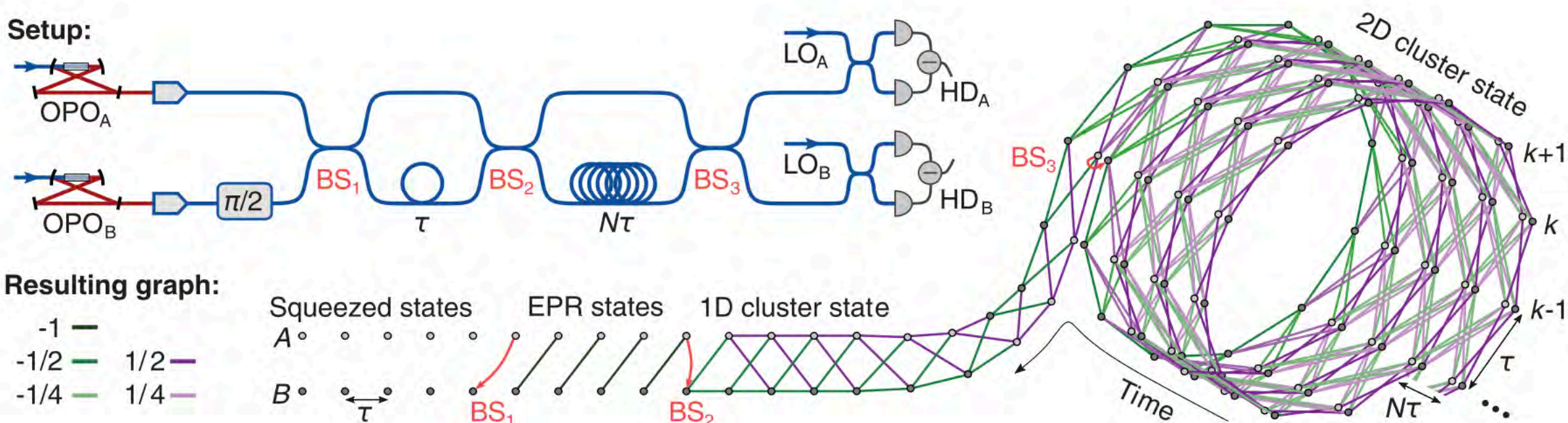


## QUANTUM COMPUTING

### Deterministic generation of a two-dimensional cluster state

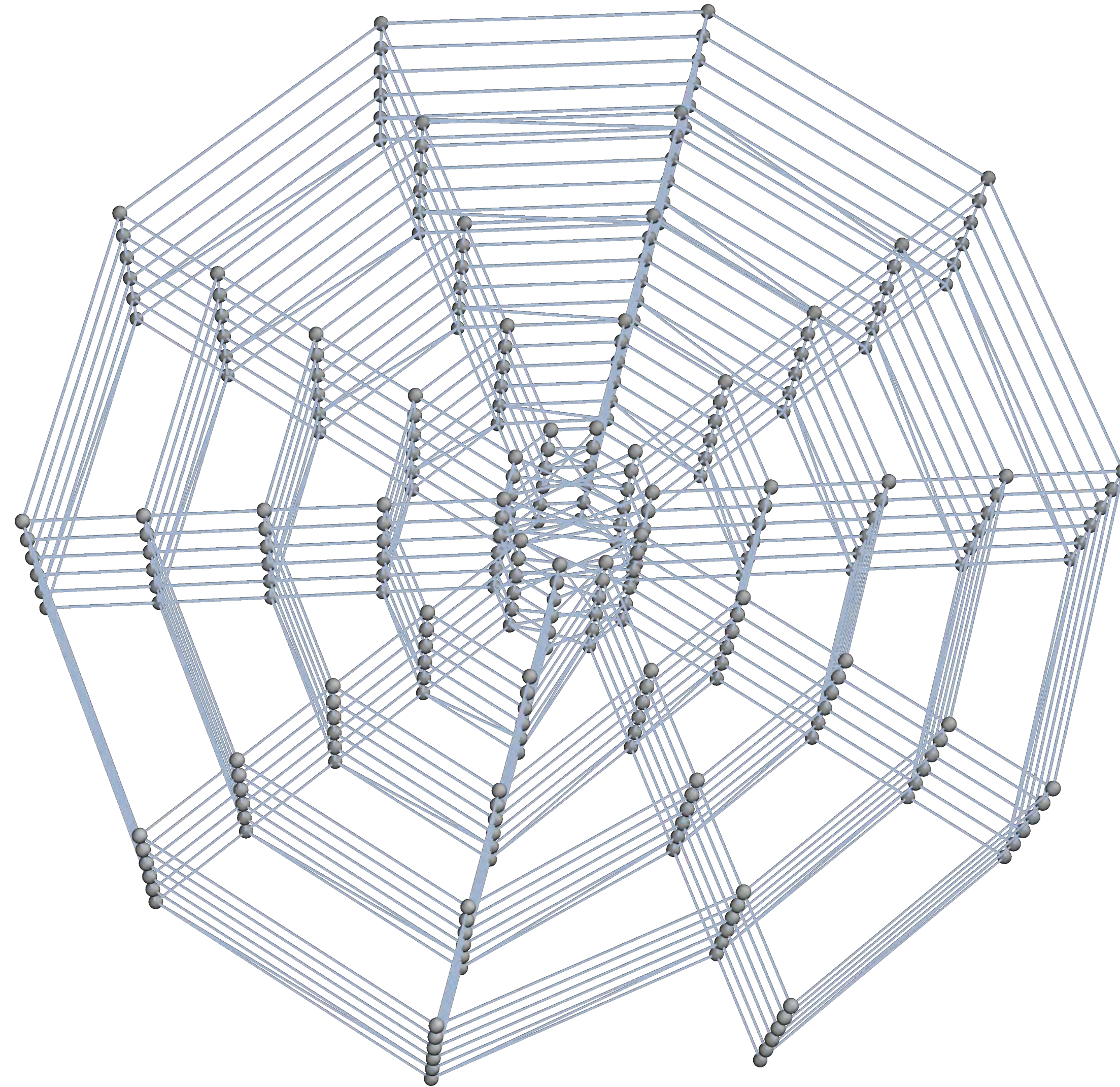
Mikkel V. Larsen\*, Xueshi Guo, Casper R. Breum, Jonas S. Neergaard-Nielsen, Ulrik L. Andersen\*

Larsen *et al.*, *Science* **366**, 369–372 (2019) 18 October 2019





# 3D states (and beyond) just around the corner are



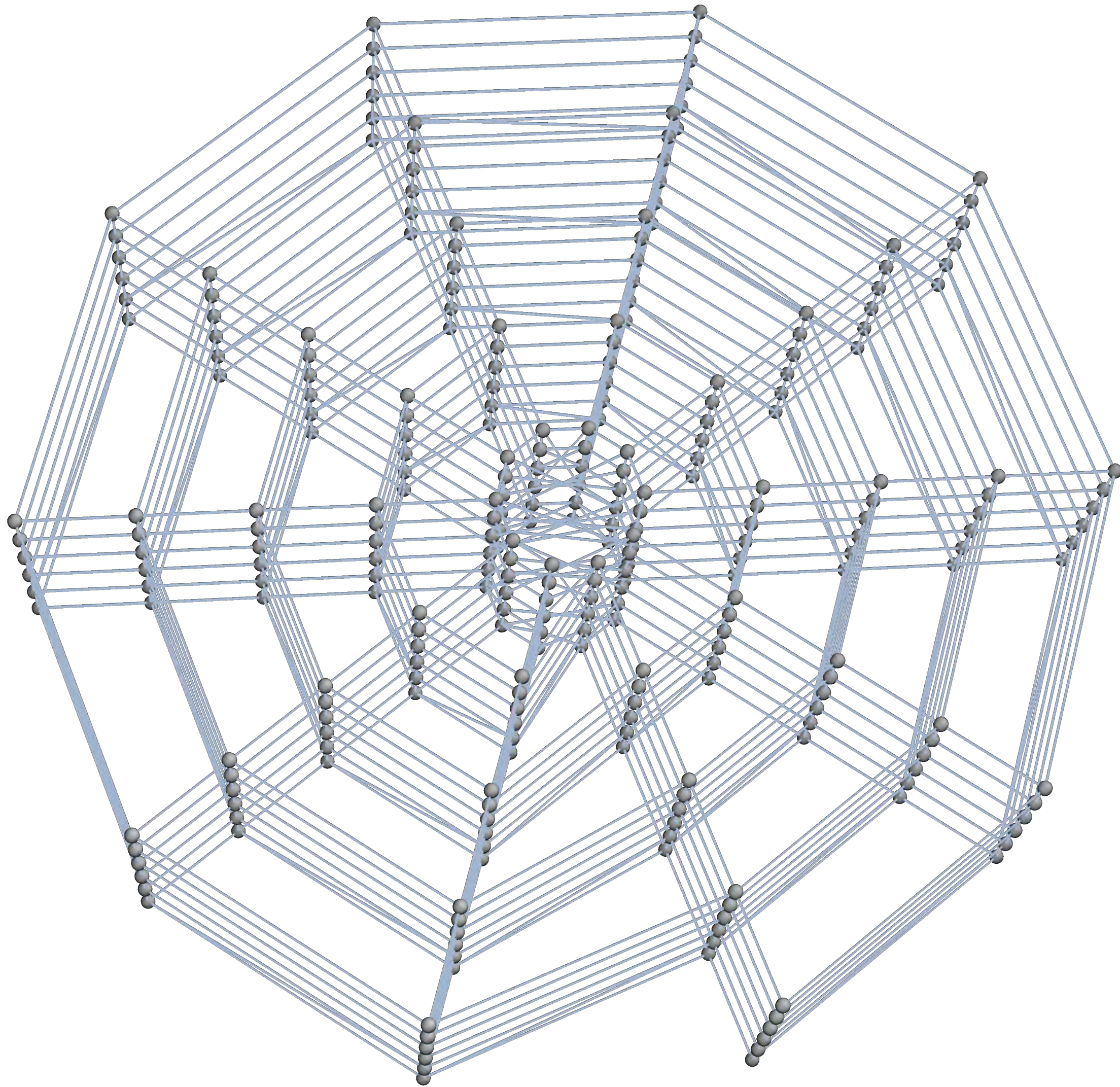


3D states (and beyond) just around the corner are



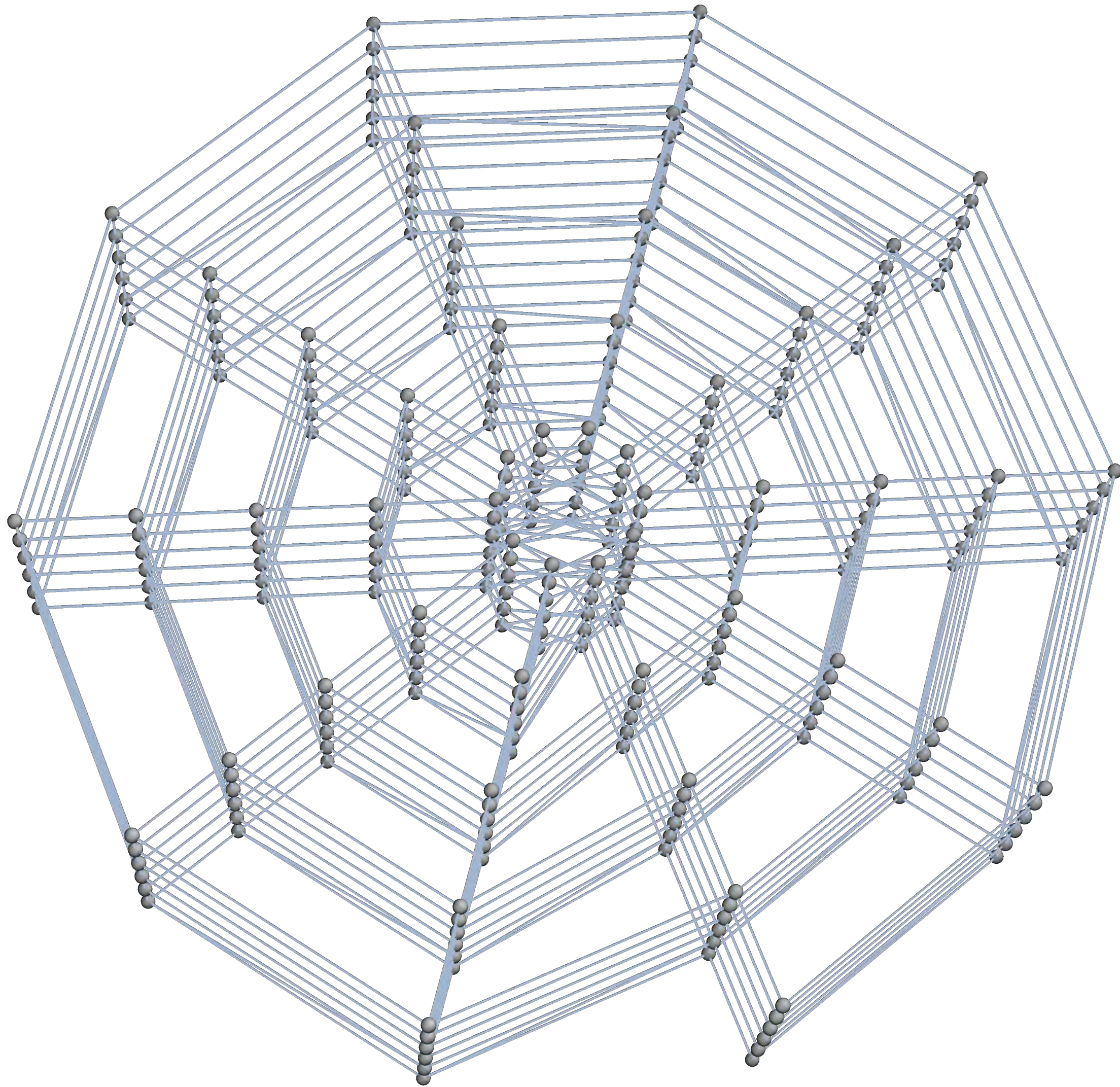


# 3D states (and beyond) just around the corner are





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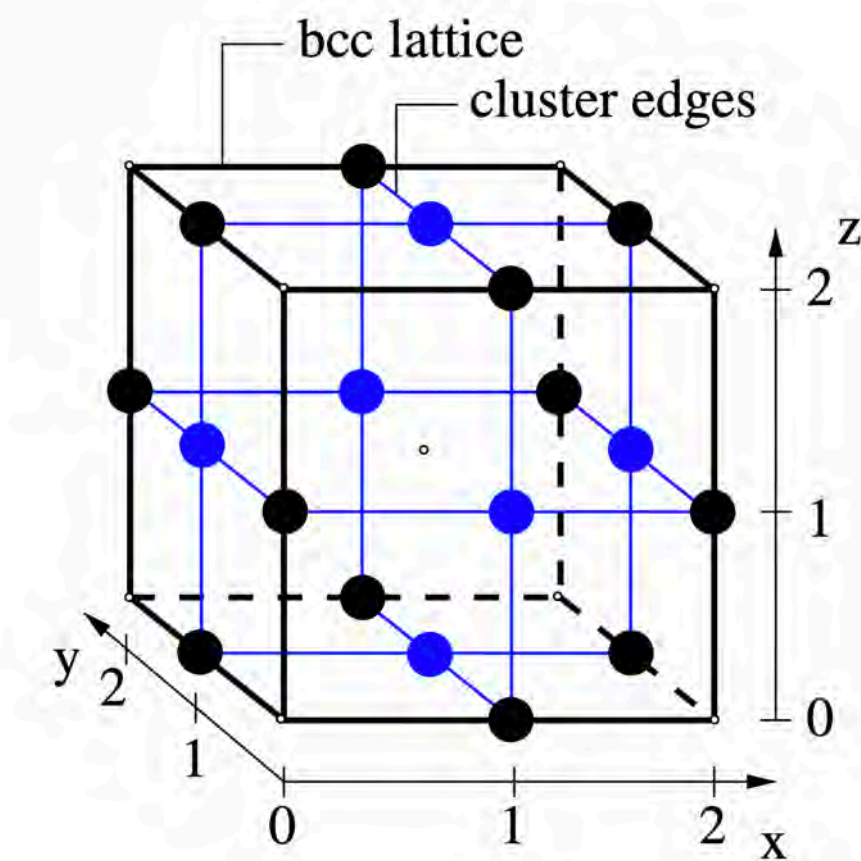
Annals of Physics 321 (2006) 2242–2270

ANNALS  
of  
PHYSICS

[www.elsevier.com/locate/aop](http://www.elsevier.com/locate/aop)

## A fault-tolerant one-way quantum computer

R. Raussendorf <sup>a,\*</sup>, J. Harrington <sup>b</sup>, K. Goyal <sup>a</sup>





# Take-home message

The generation of sophisticated, QC-universal cluster states on a very large scale (thousands of qumodes) is

EASY

over continuous variables and is also

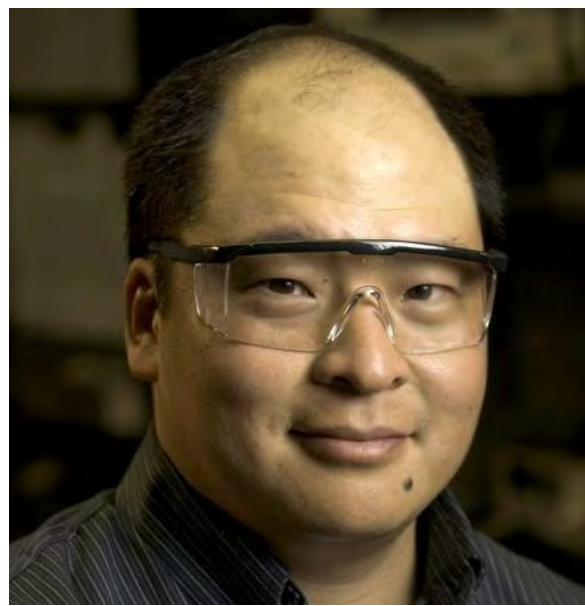
highly compatible with an integrated optics approach. (Stay tuned.)



# Non-Gaussian quantum optics (photons)



# An experimentally accessible non-Gaussian operation: photon-number detection

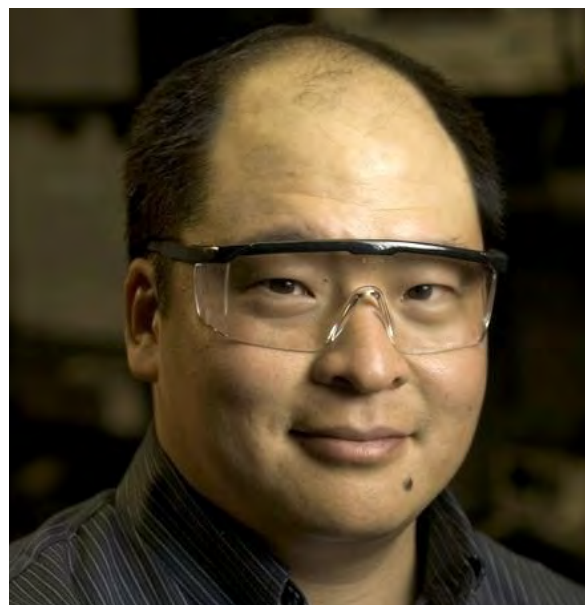


Sae Woo Nam  
(NIST)



Aaron Miller  
(Quantum Opus)





Sae Woo Nam  
(NIST)

# An experimentally accessible non-Gaussian operation: photon-number detection

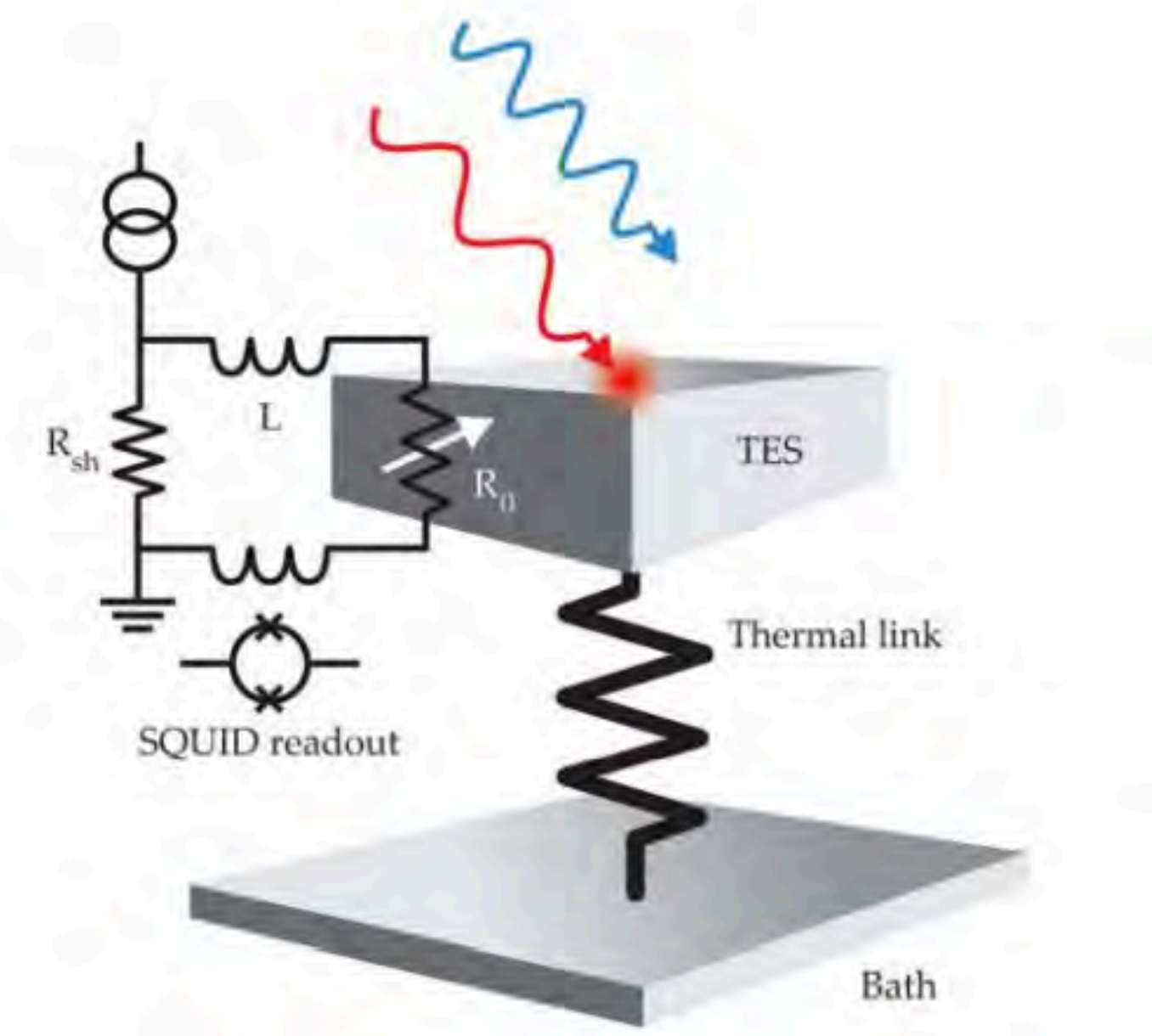
$$\text{POVM set} \equiv \{ |n\rangle \langle n| \}_{n=0, \dots, n_{\max}}$$

(ideally!)

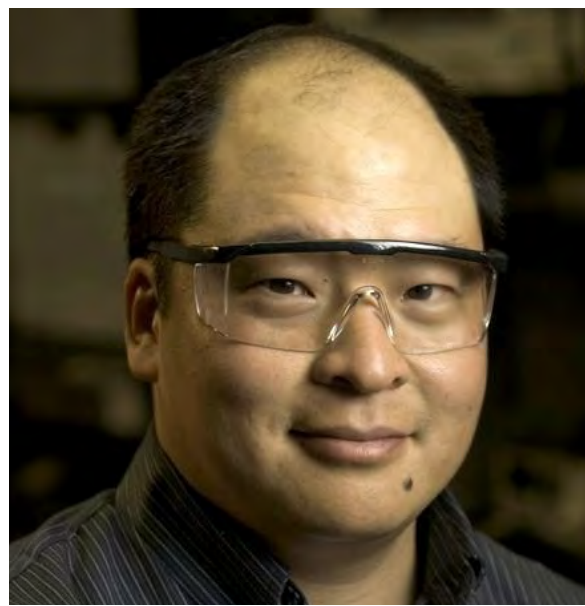


Aaron Miller  
(Quantum Opus)

Superconducting transition-edge sensor



Physics Today **71**, 8, 28 (2018)



Sae Woo Nam  
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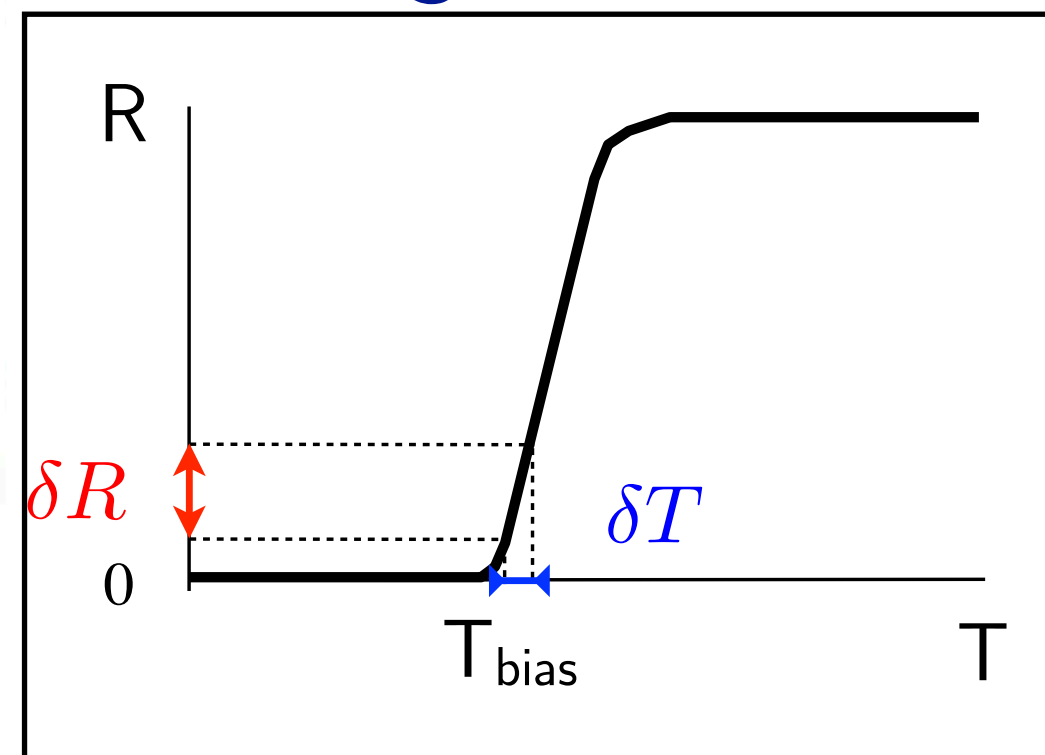
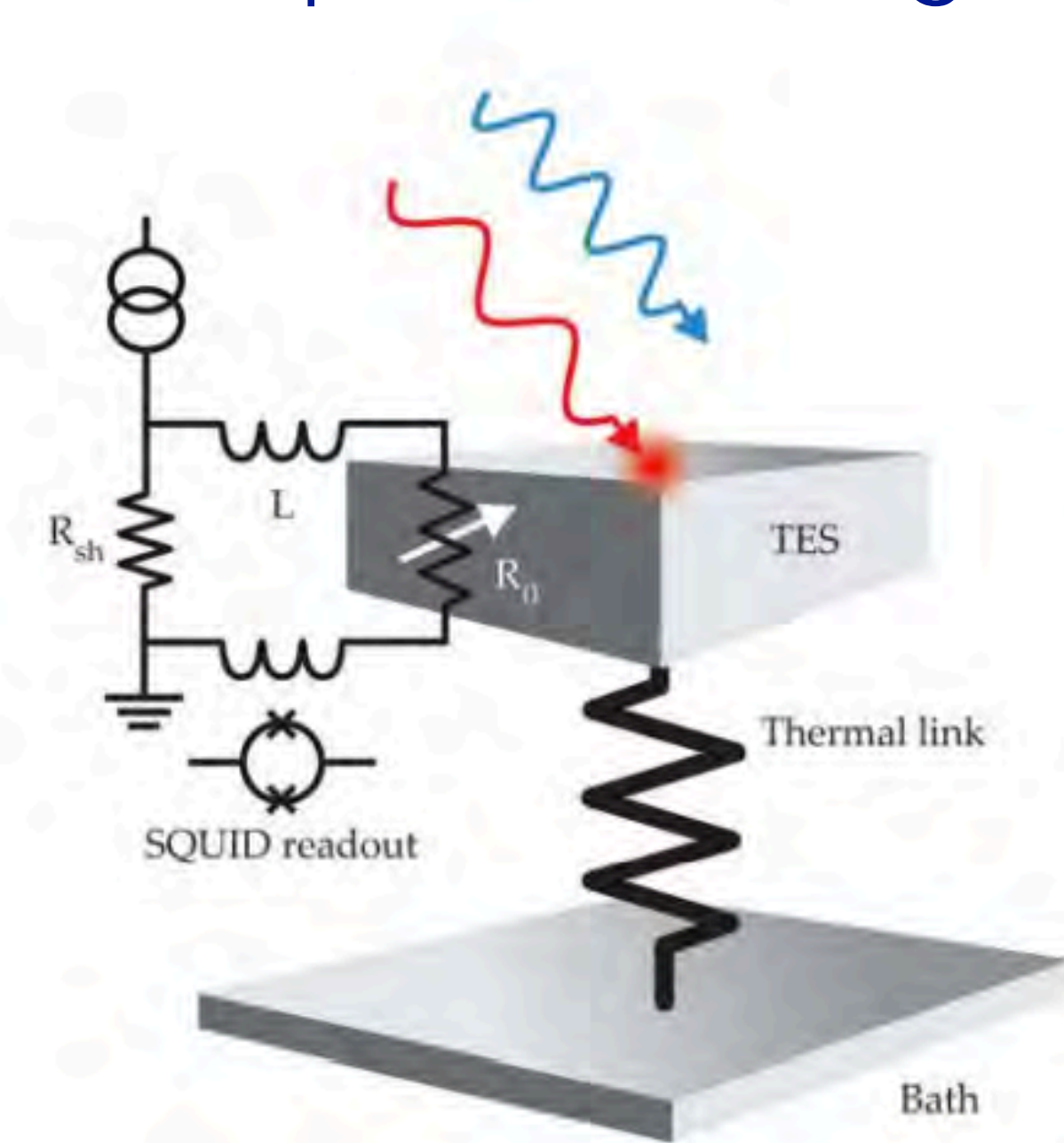
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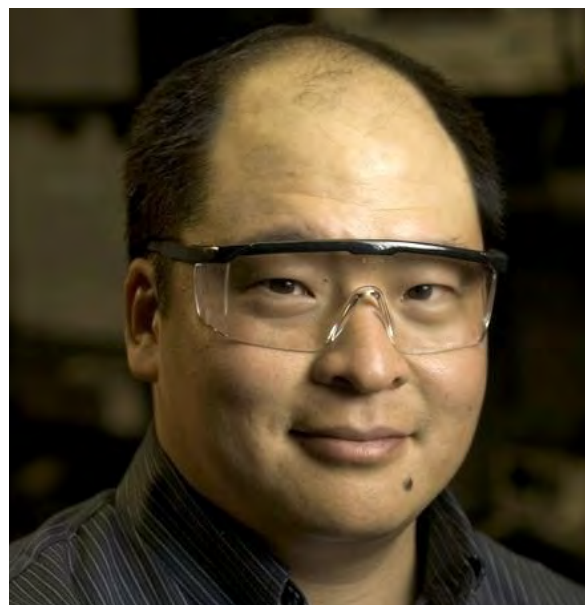
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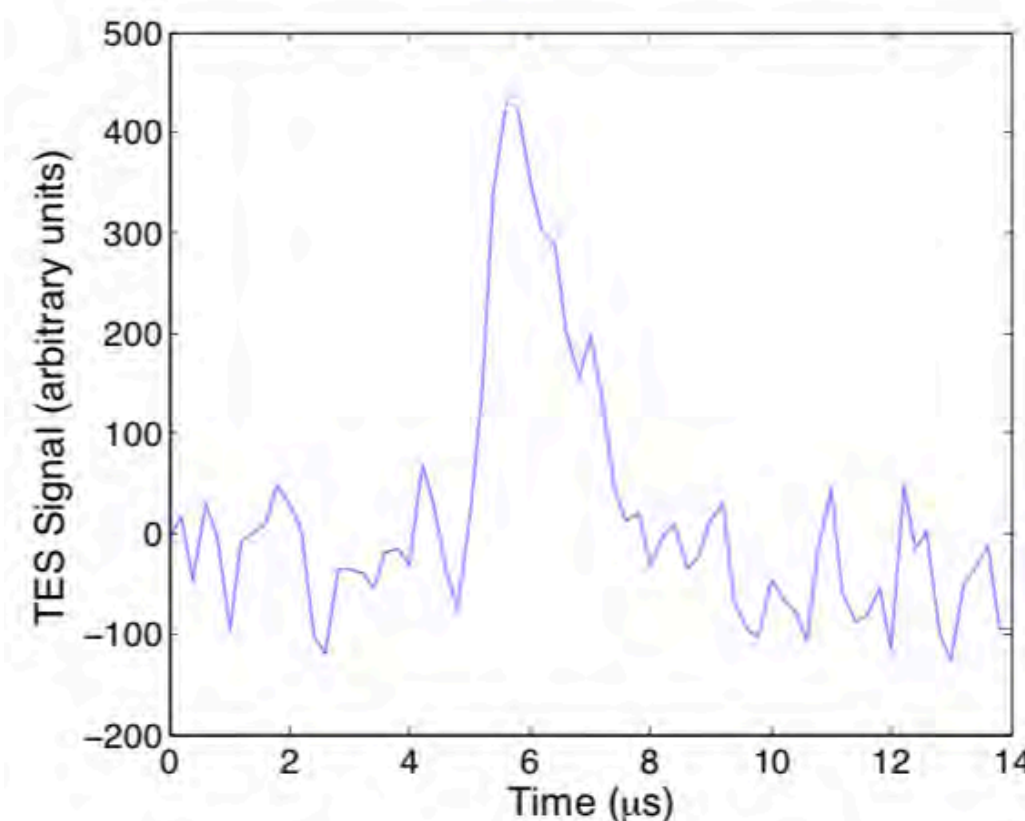
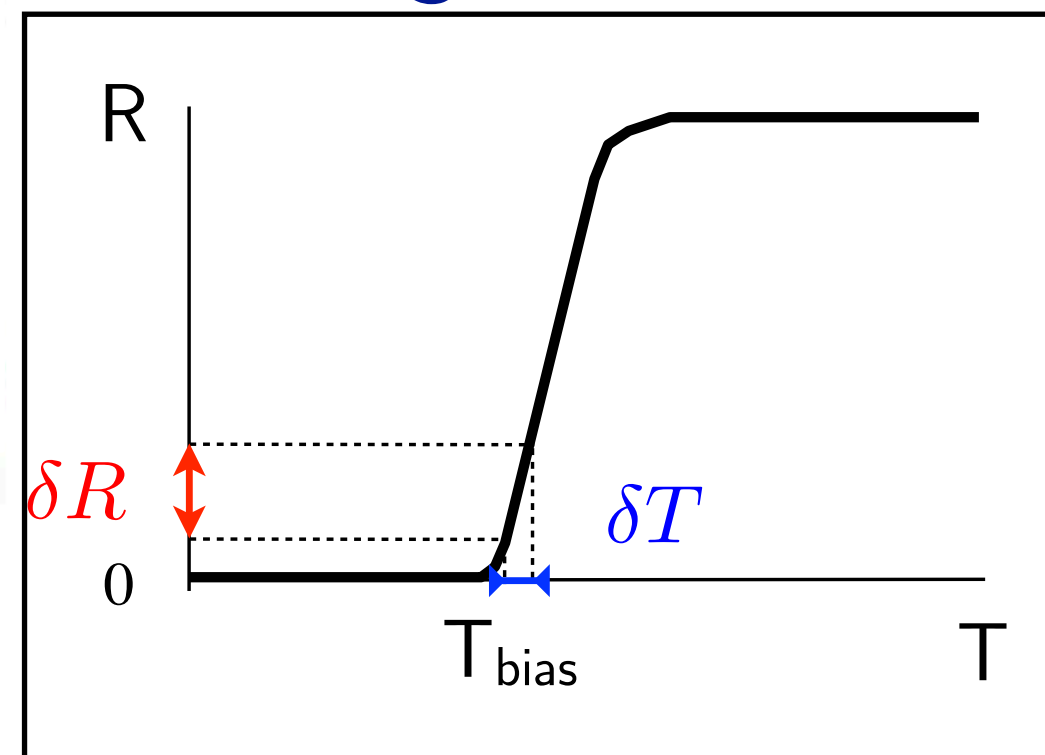
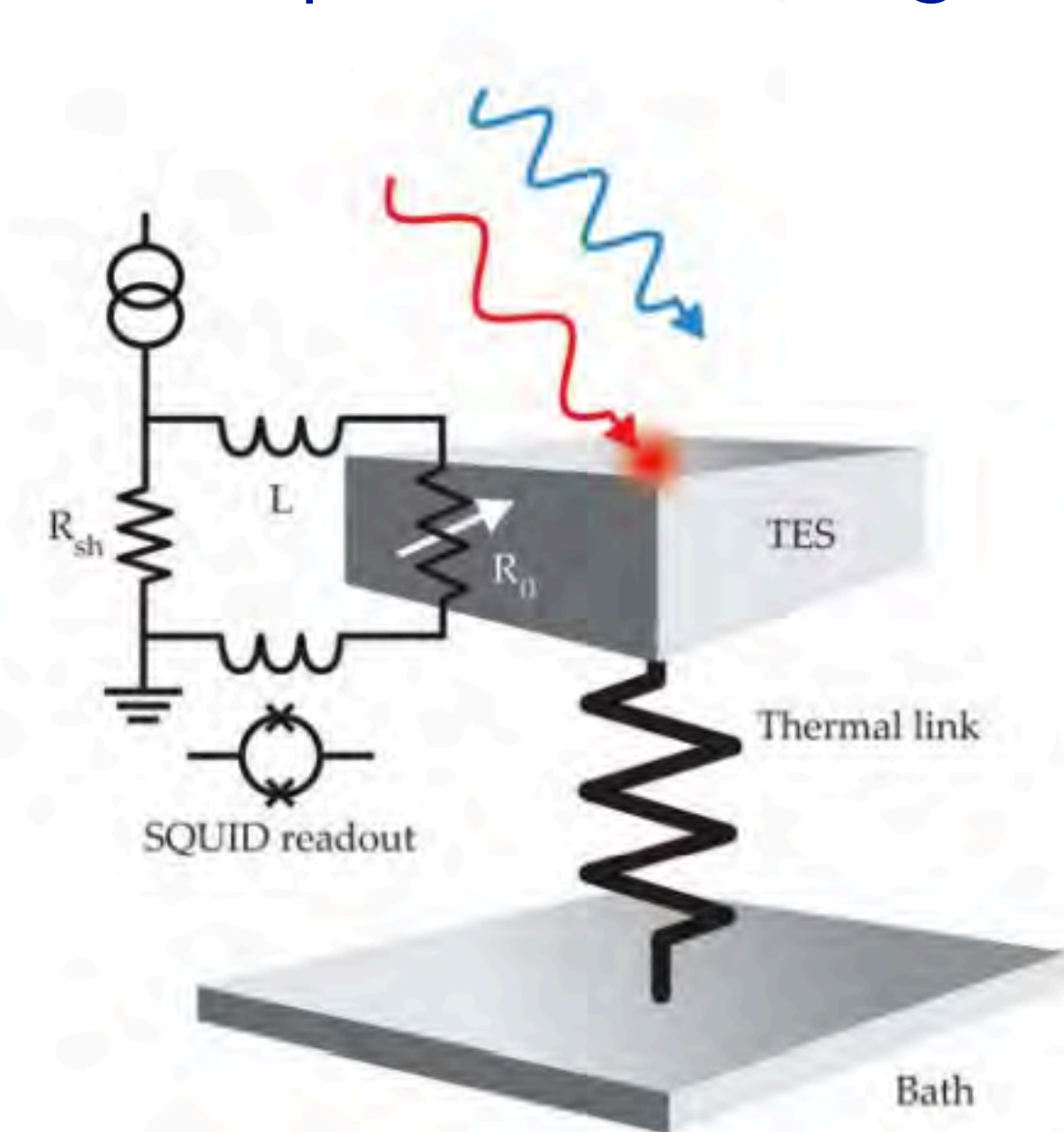
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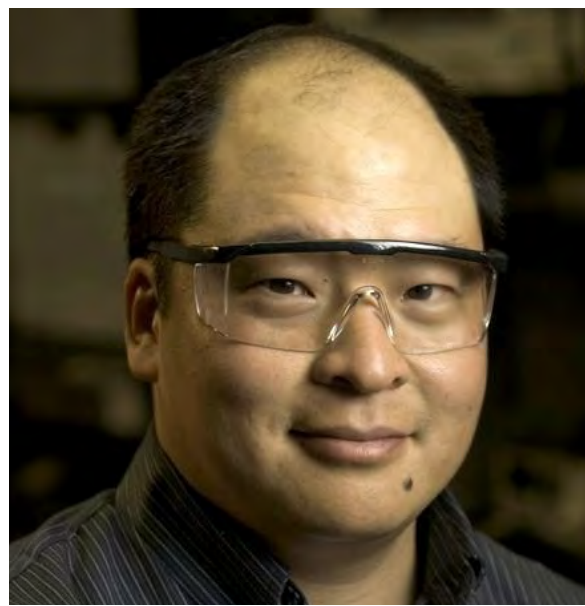


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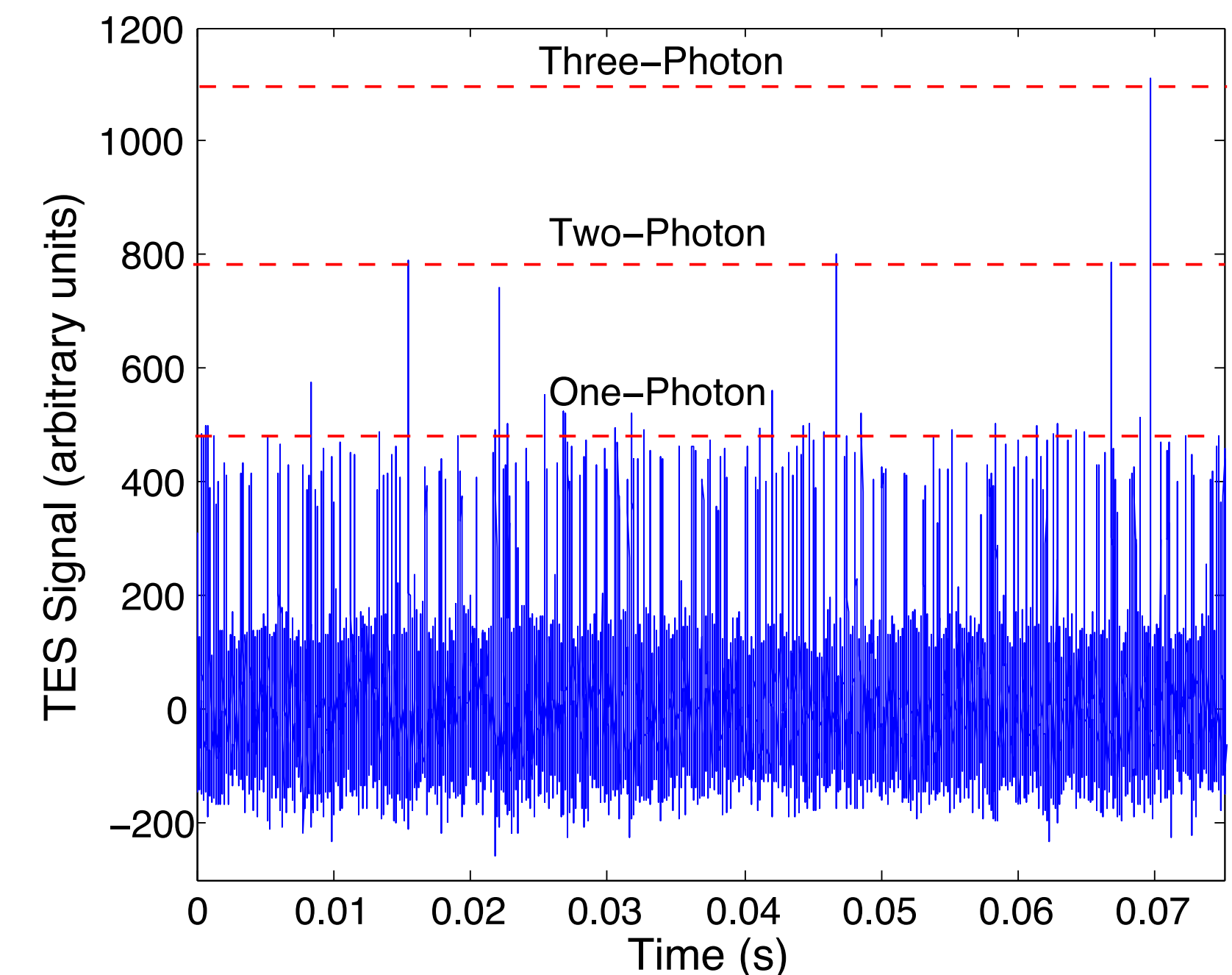
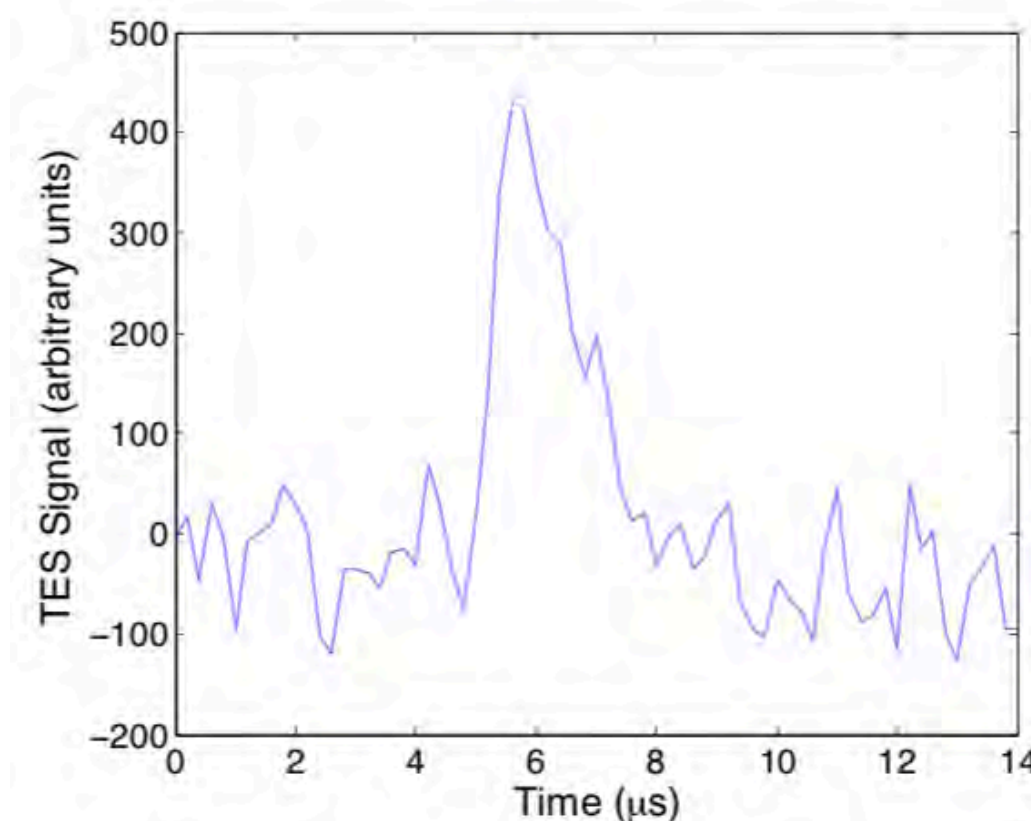
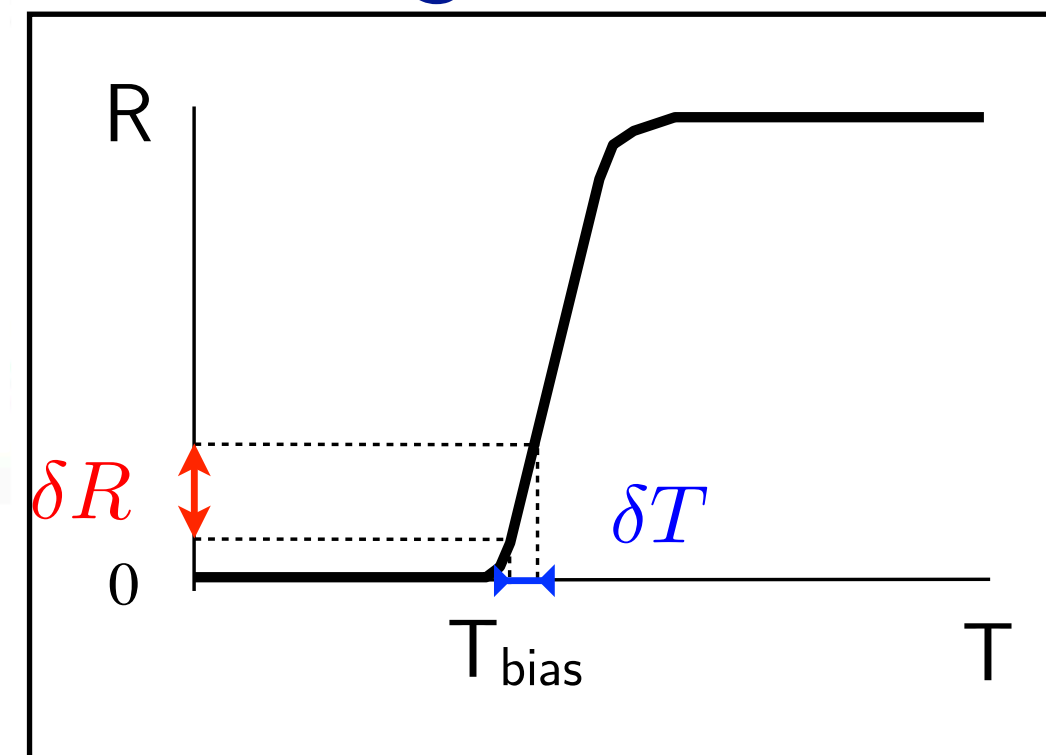
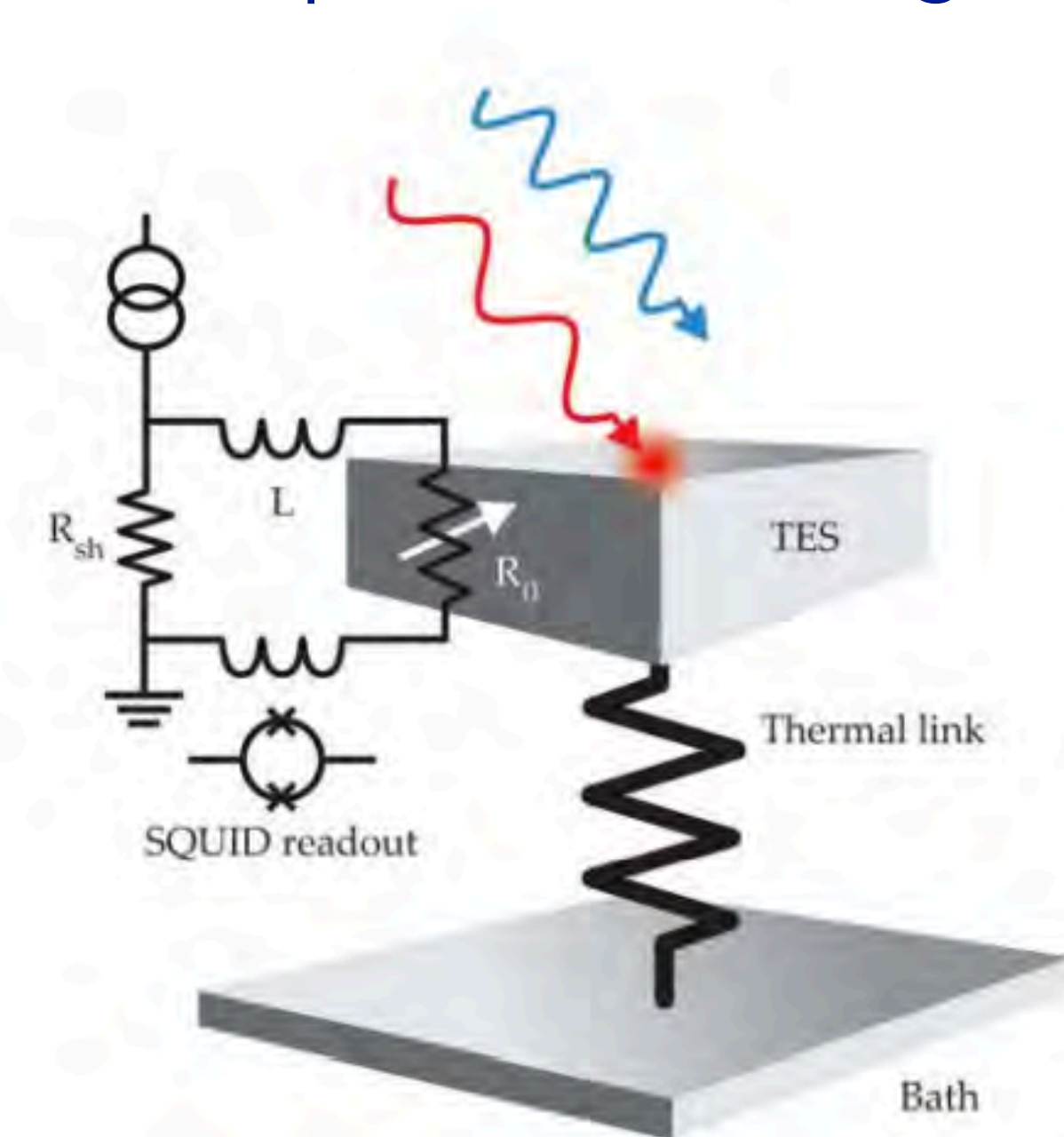
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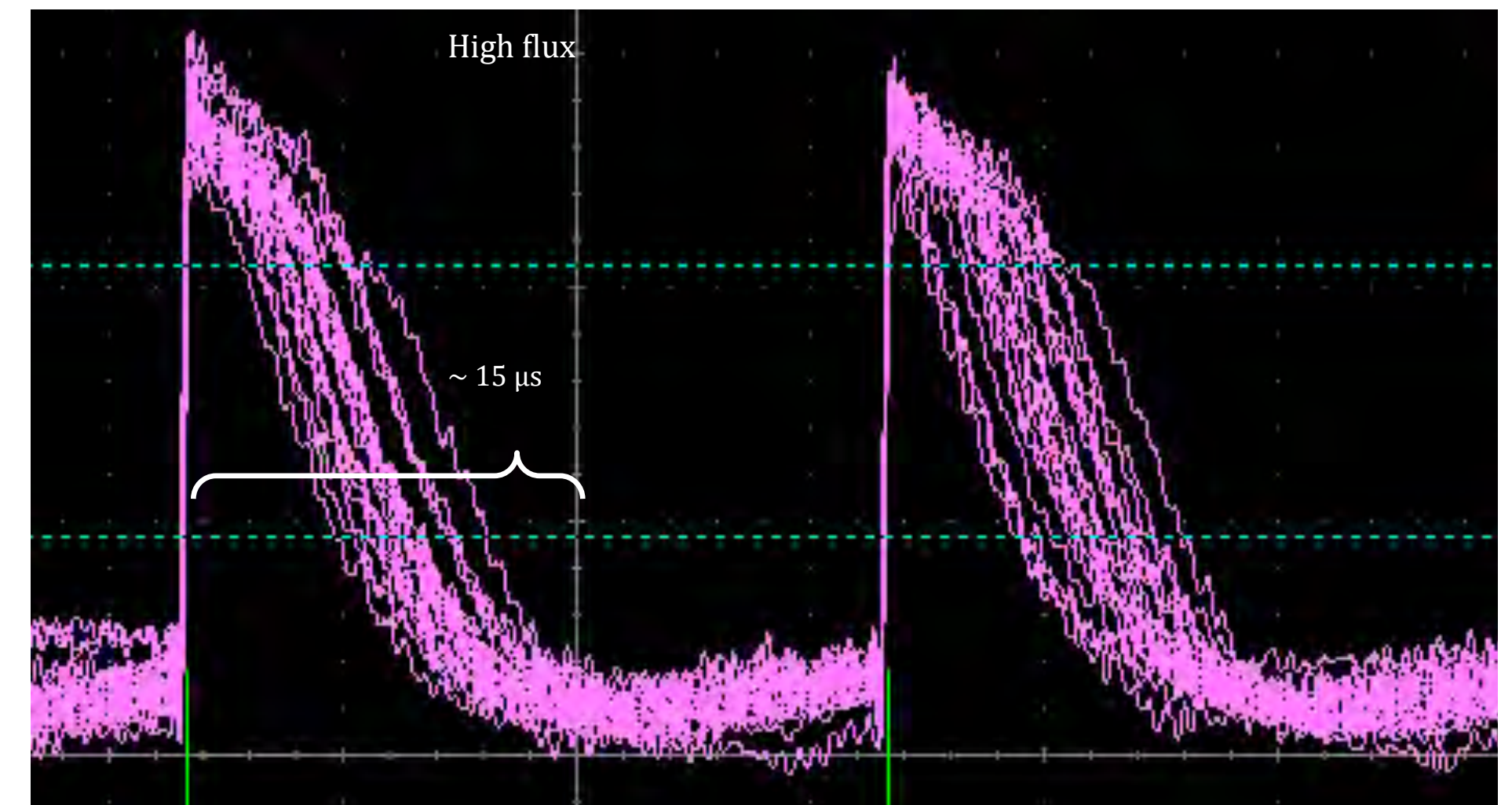
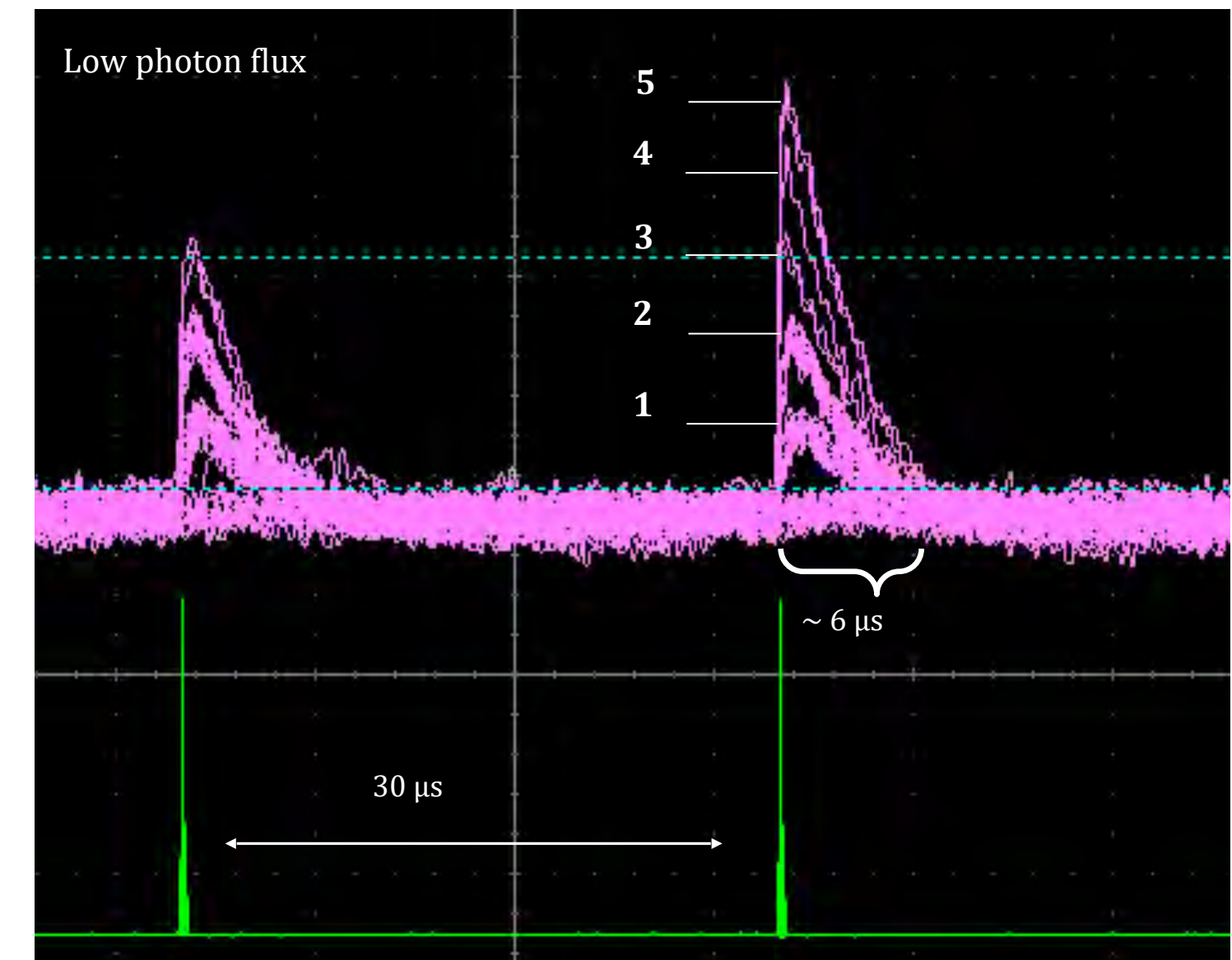
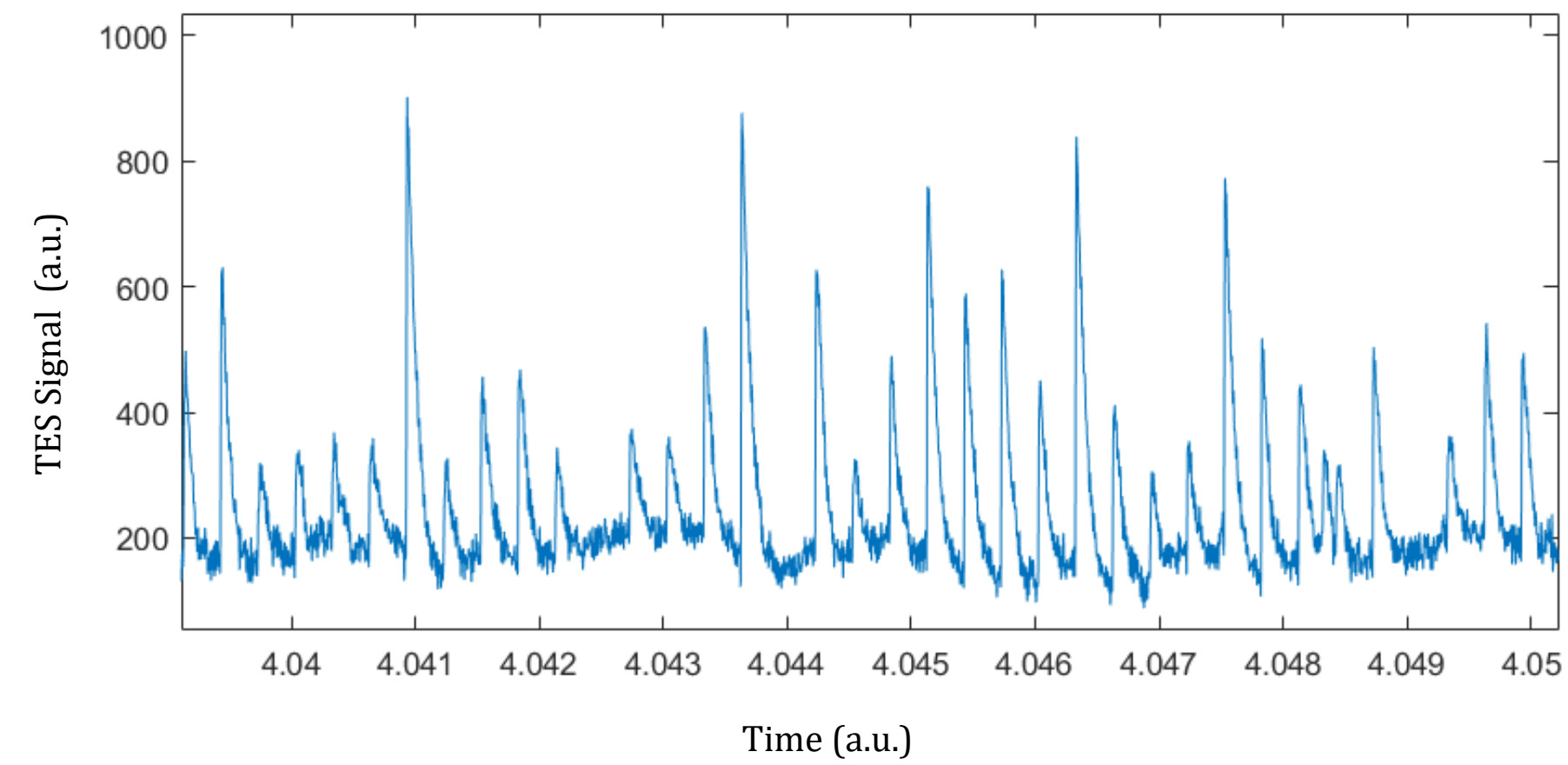
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Physics Today **71**, 8, 28 (2018)

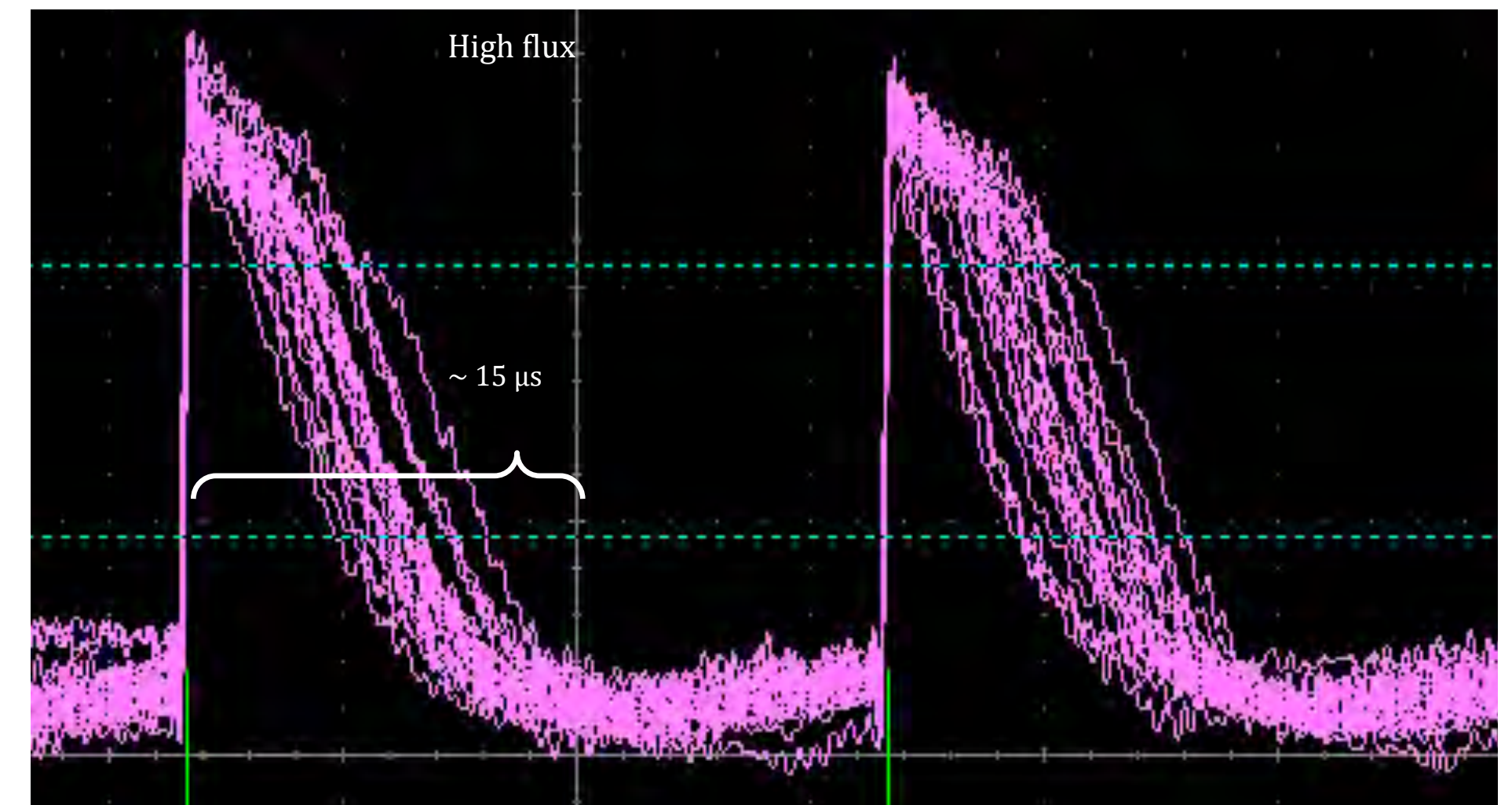
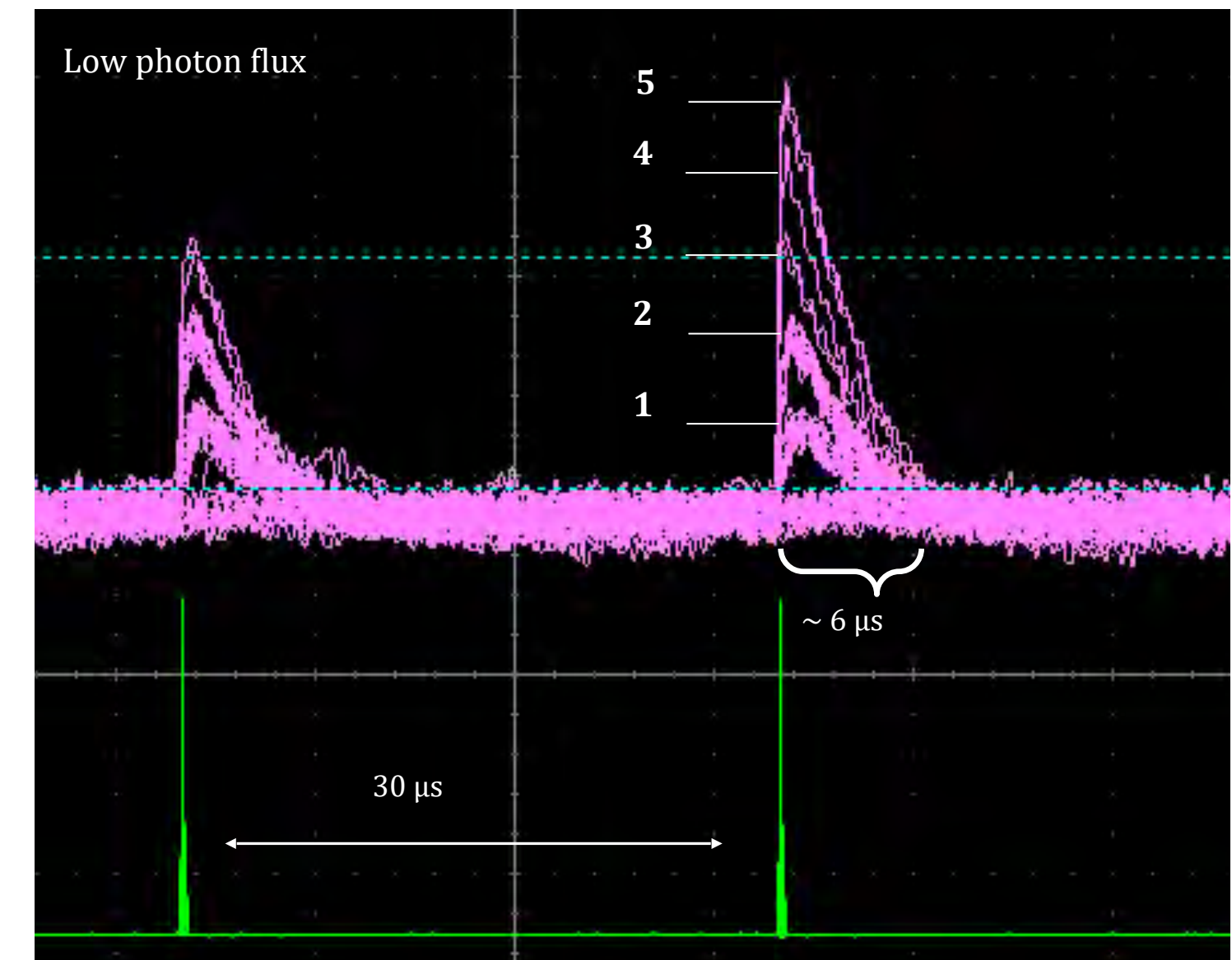
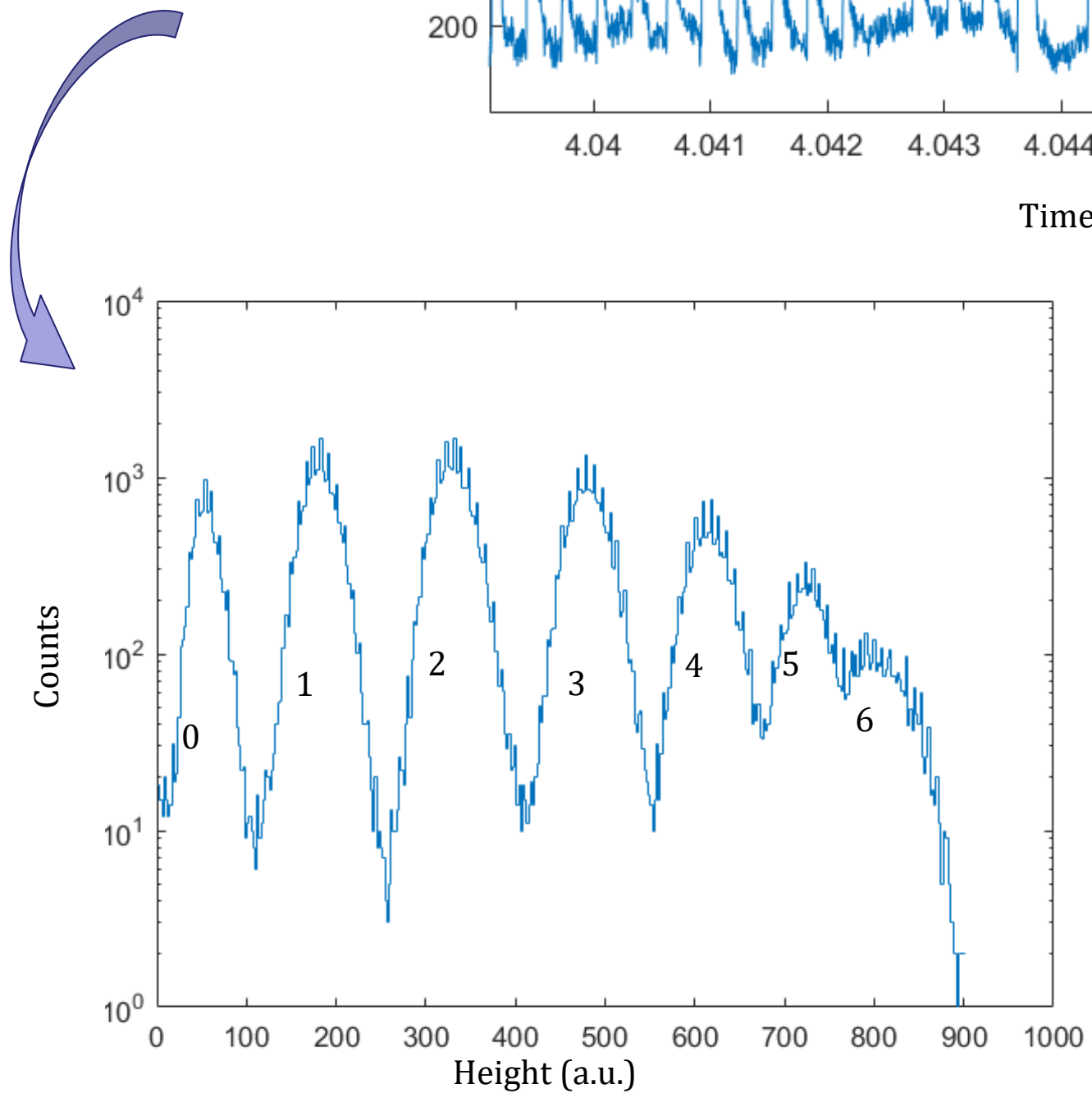
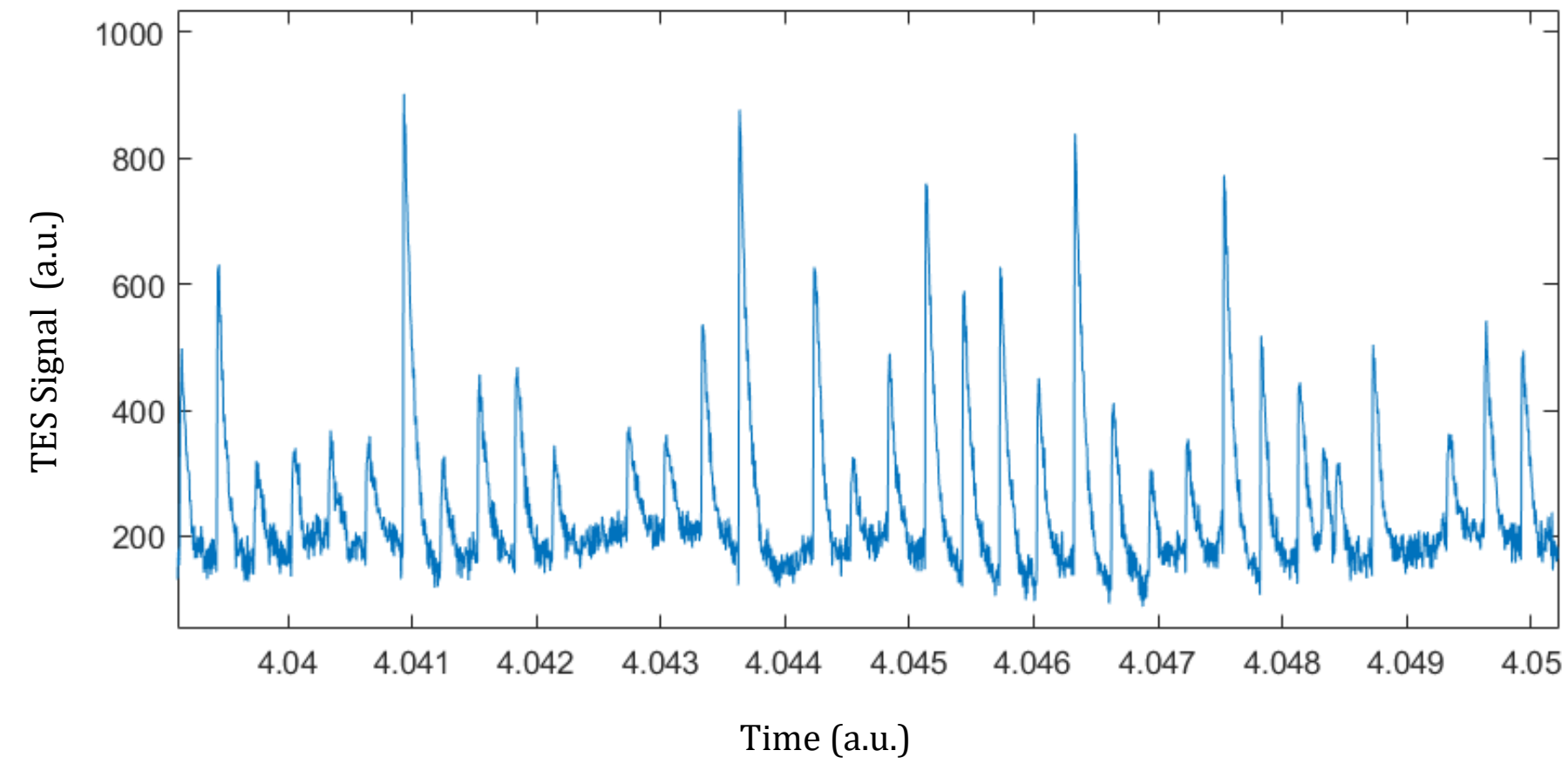


# Processing TES signal: PNR detection



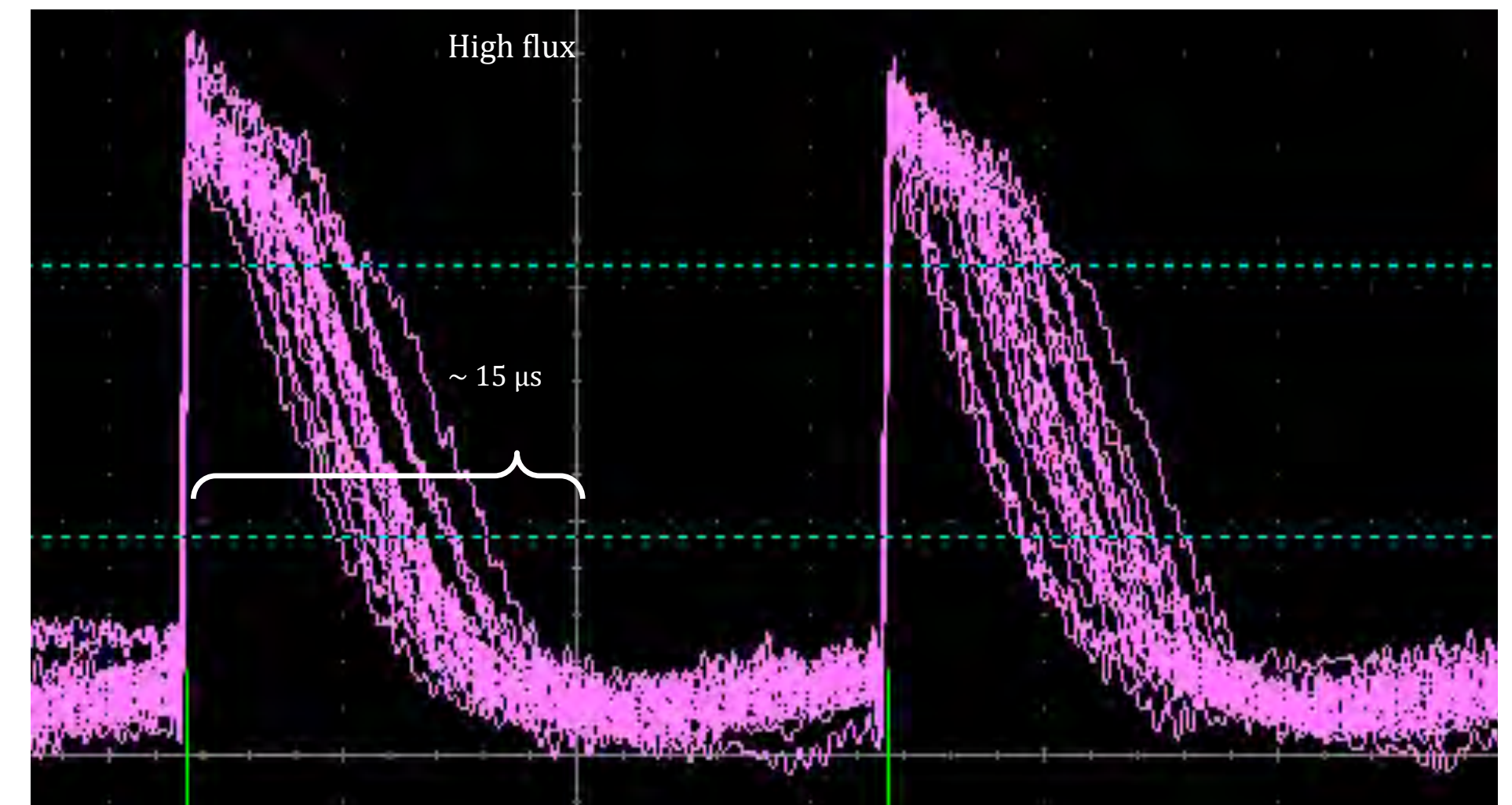
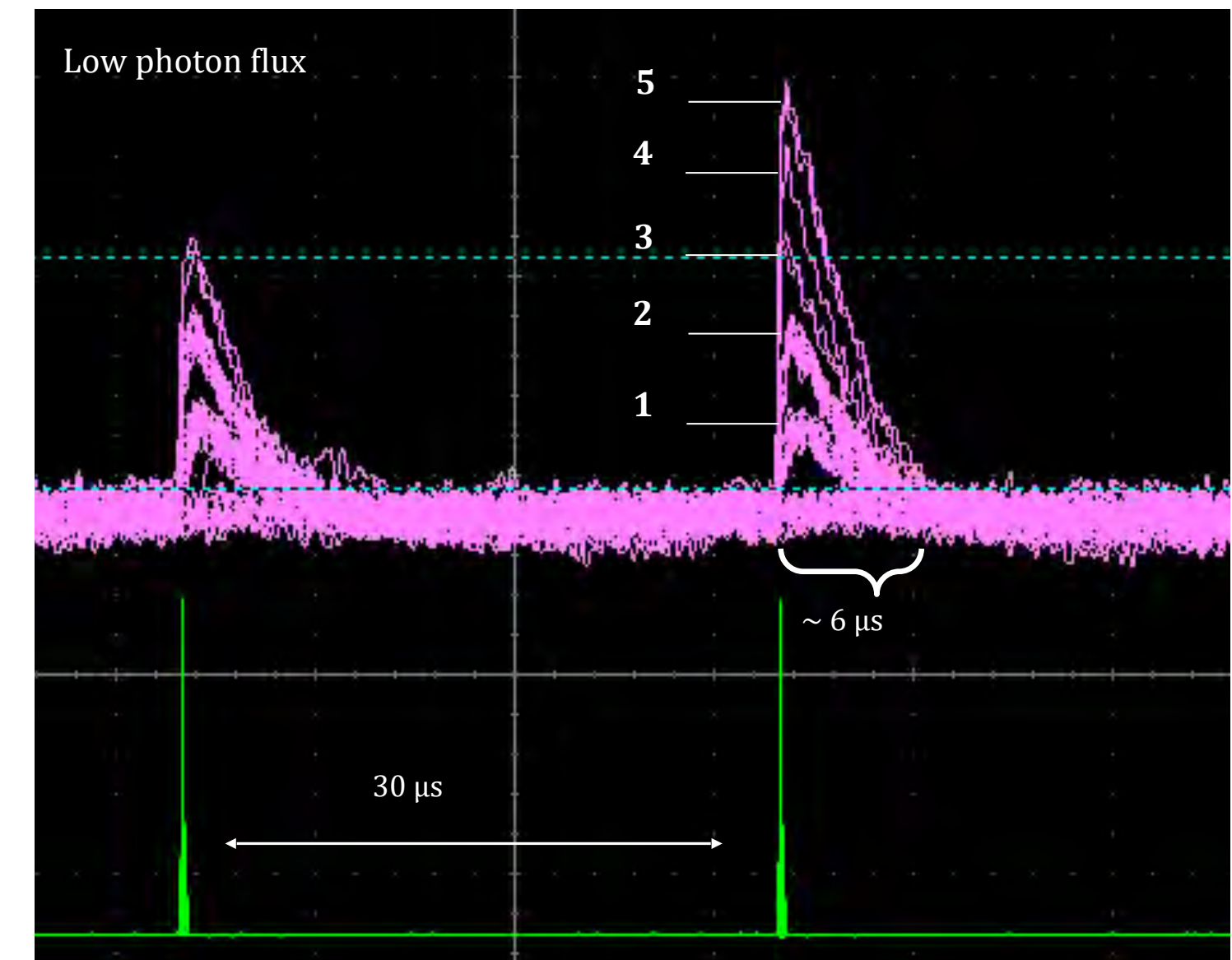
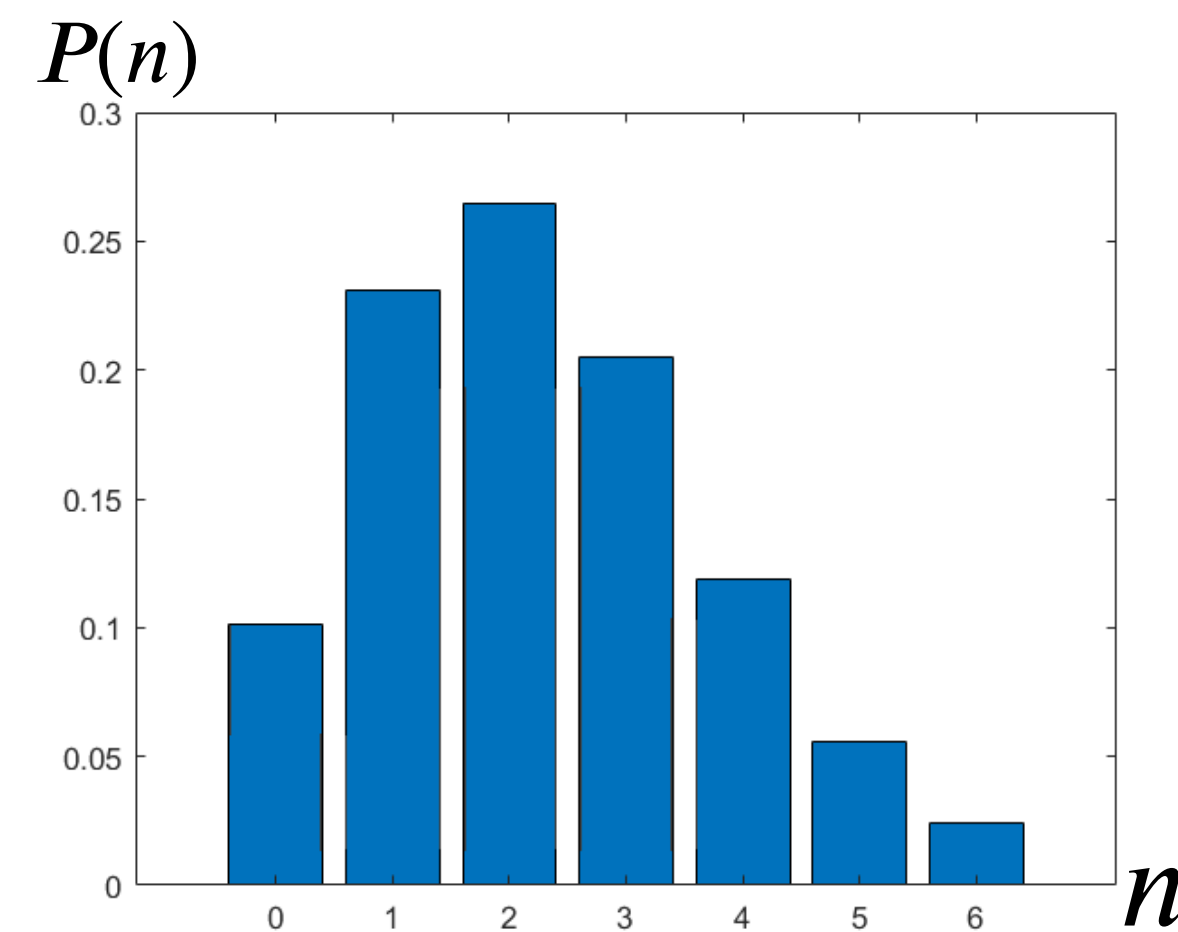
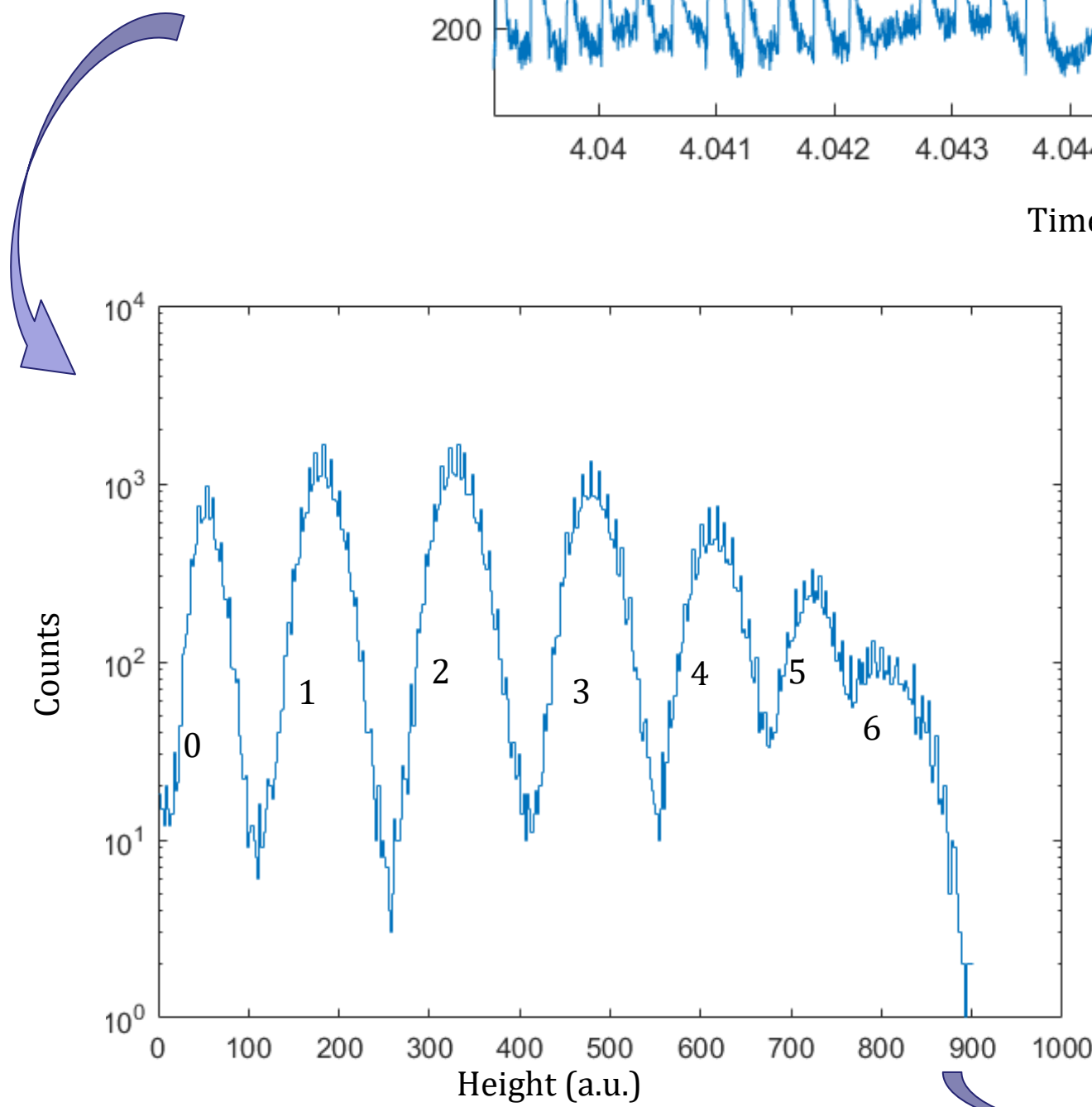
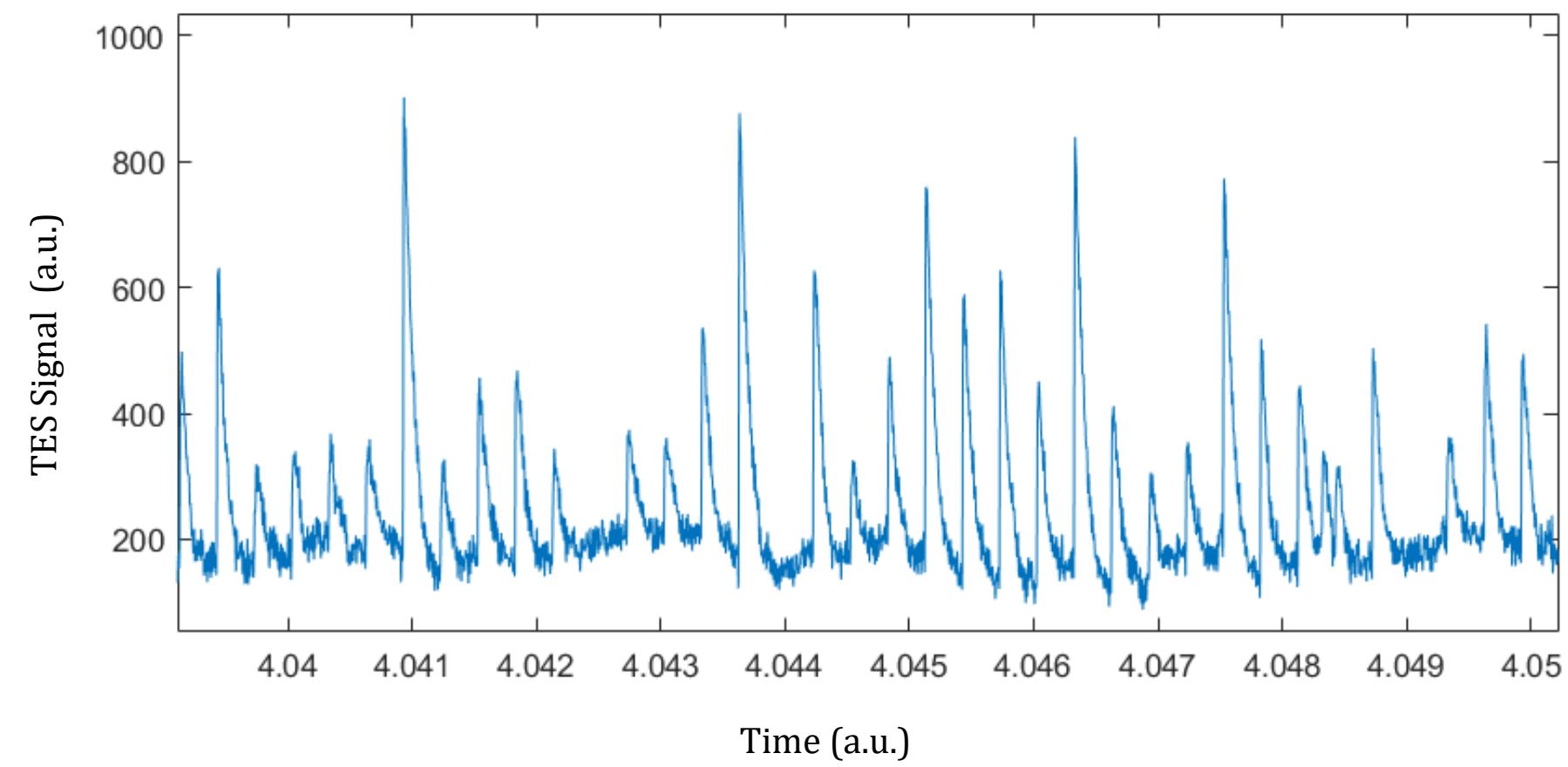


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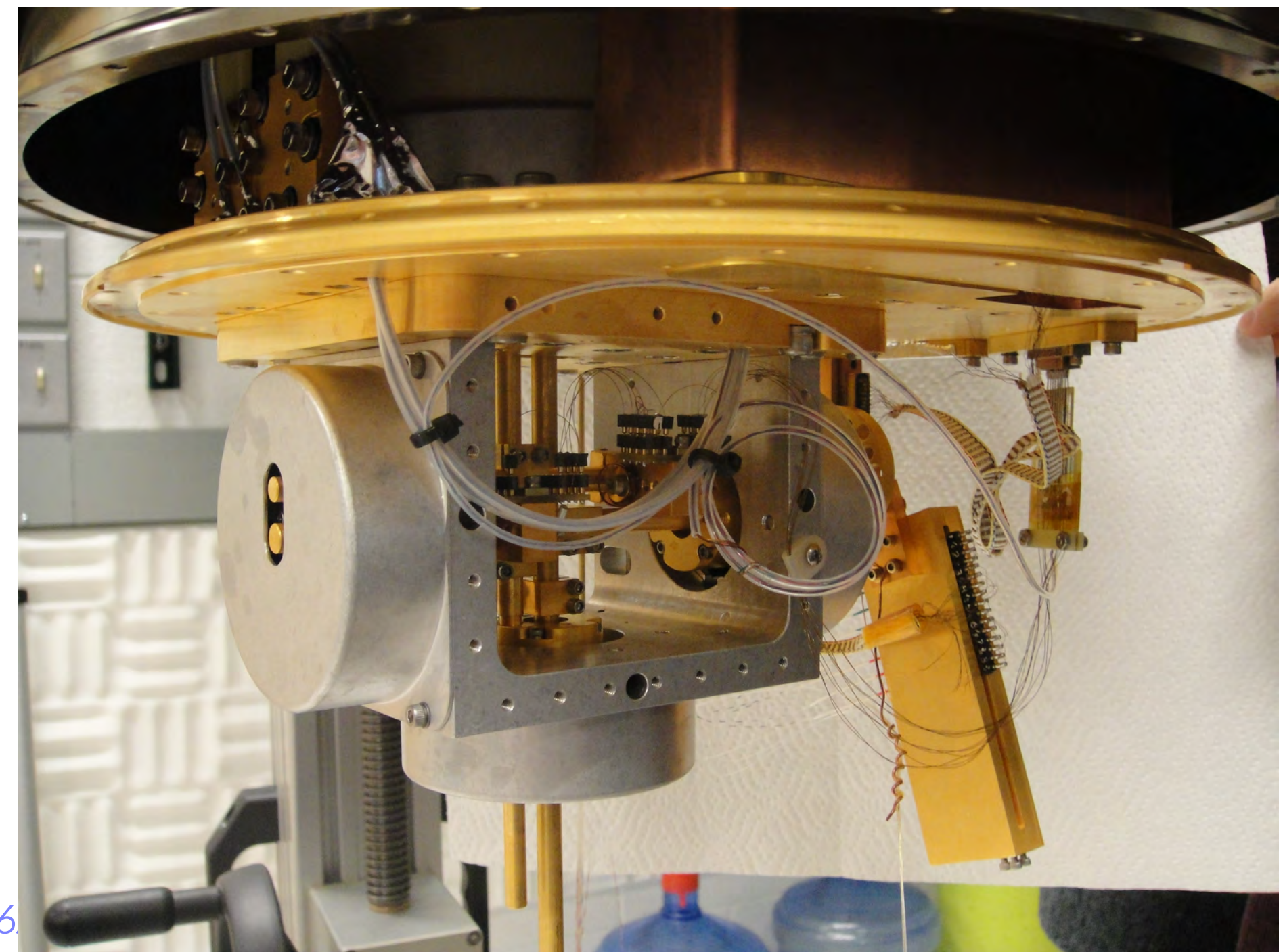
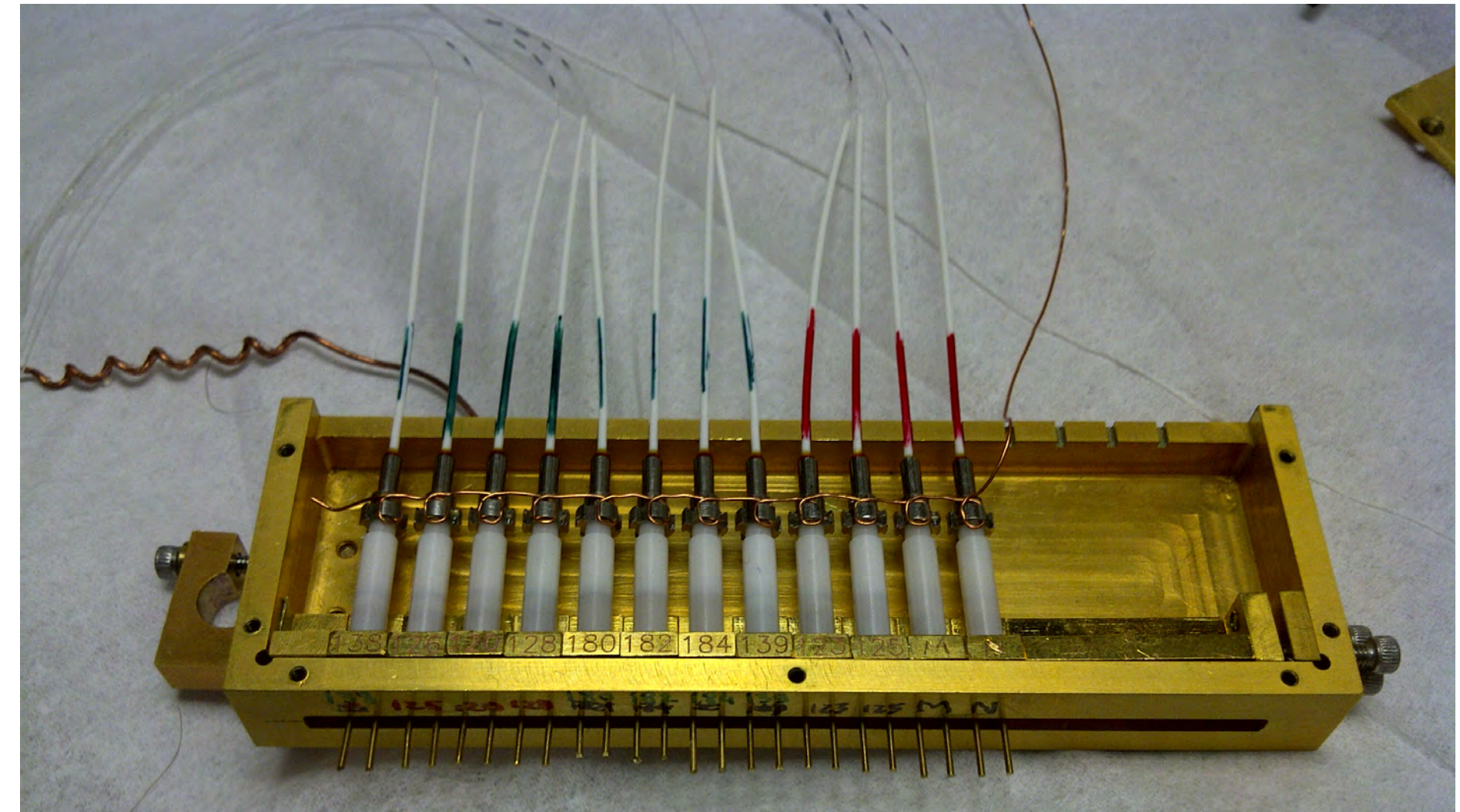
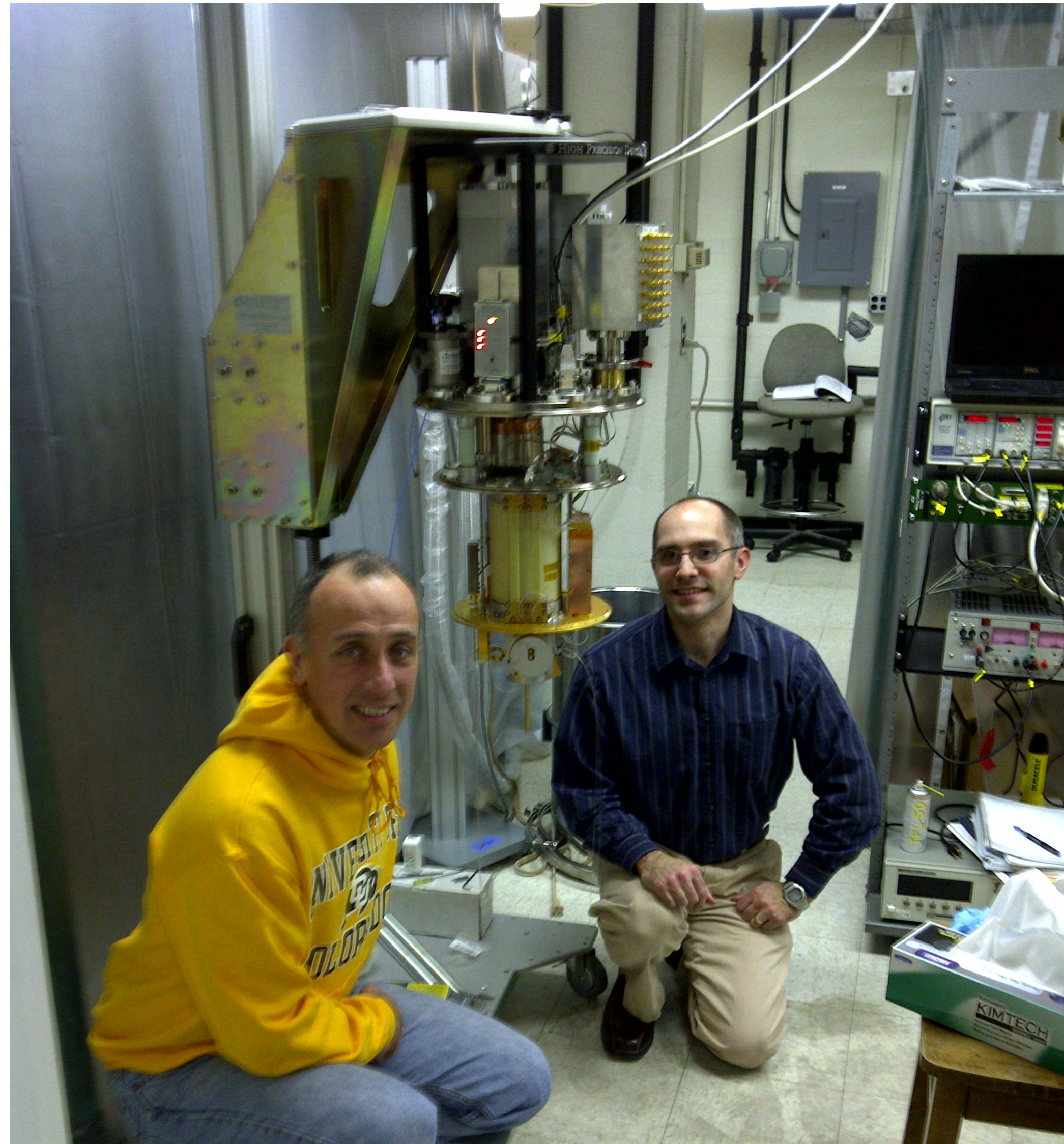


# Processing TES signal: PNR detection





# TES @ PfisterLabs





# Quantum tomography with photon counting

S. Wallentowitz & W. Vogel, PRA (1996)

K. Banaszek & W. Wódkiewicz, PRL (1996)

WIGNER FUNCTION

$$W(q, p) = \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{iyp} \left\langle q - \frac{y}{2} \left| \rho \right| q + \frac{y}{2} \right\rangle dy$$

- Only function whose marginals yield the quantum probability distributions
- Can be NONPOSITIVE (i.e., nonGaussian for a pure state)

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$$W_{\hat{\rho}}(\alpha) = \frac{1}{\pi} \text{Tr}[\hat{\rho} D(\alpha) (-1)^{\hat{n}} D^{\dagger}(\alpha)] \quad \alpha := q + ip$$

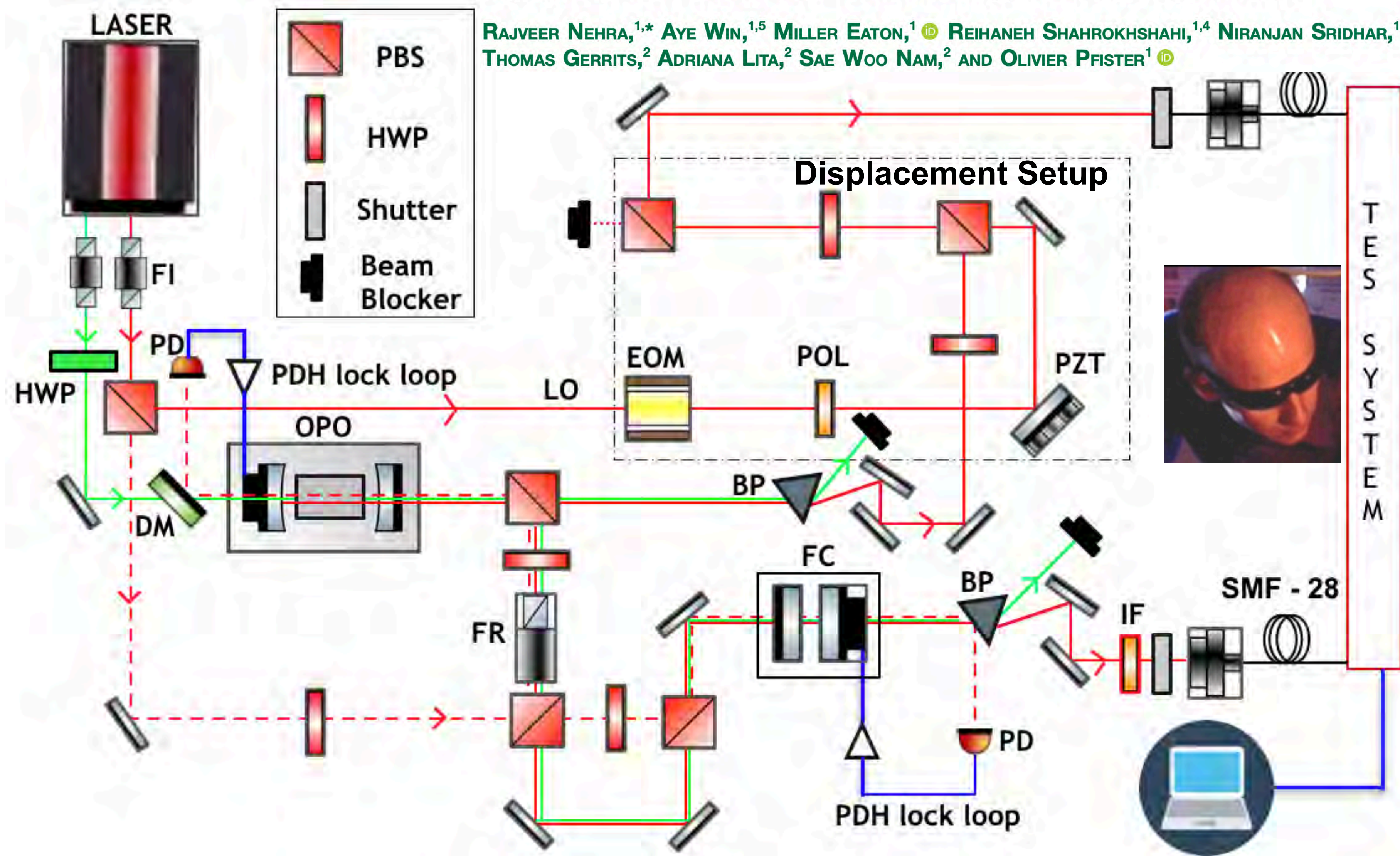
$$W_{\hat{\rho}}(\alpha) = \frac{1}{\pi} \text{Tr}[D^{\dagger}(\alpha) \hat{\rho} D(\alpha) (-1)^{\hat{n}}]$$

- Expectation value of the photon-number parity
- Easily measured directly with photon-number-resolving detection
- Raster scan of phase space by amplitude/phase shifts gives whole  $W(q,p)$



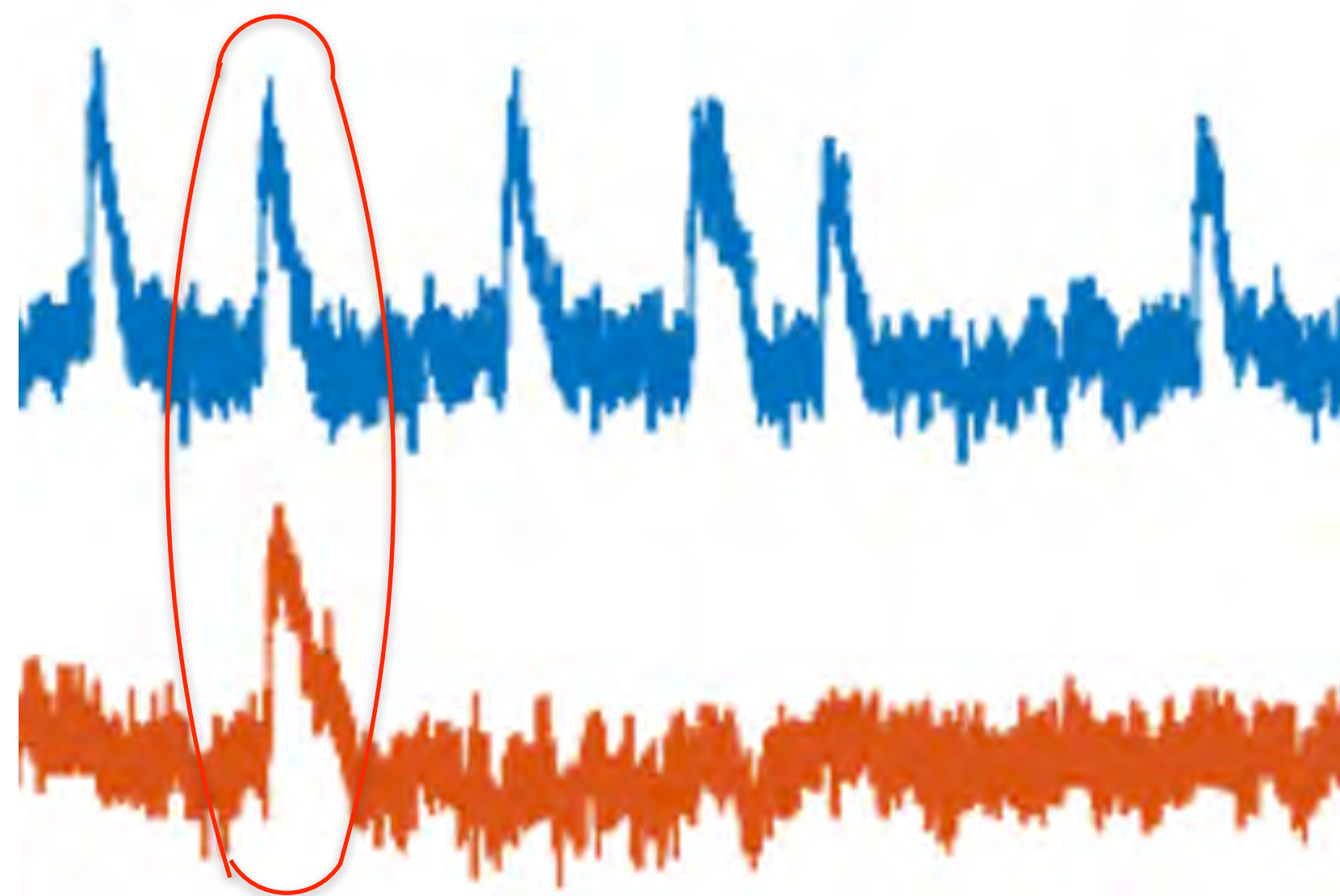
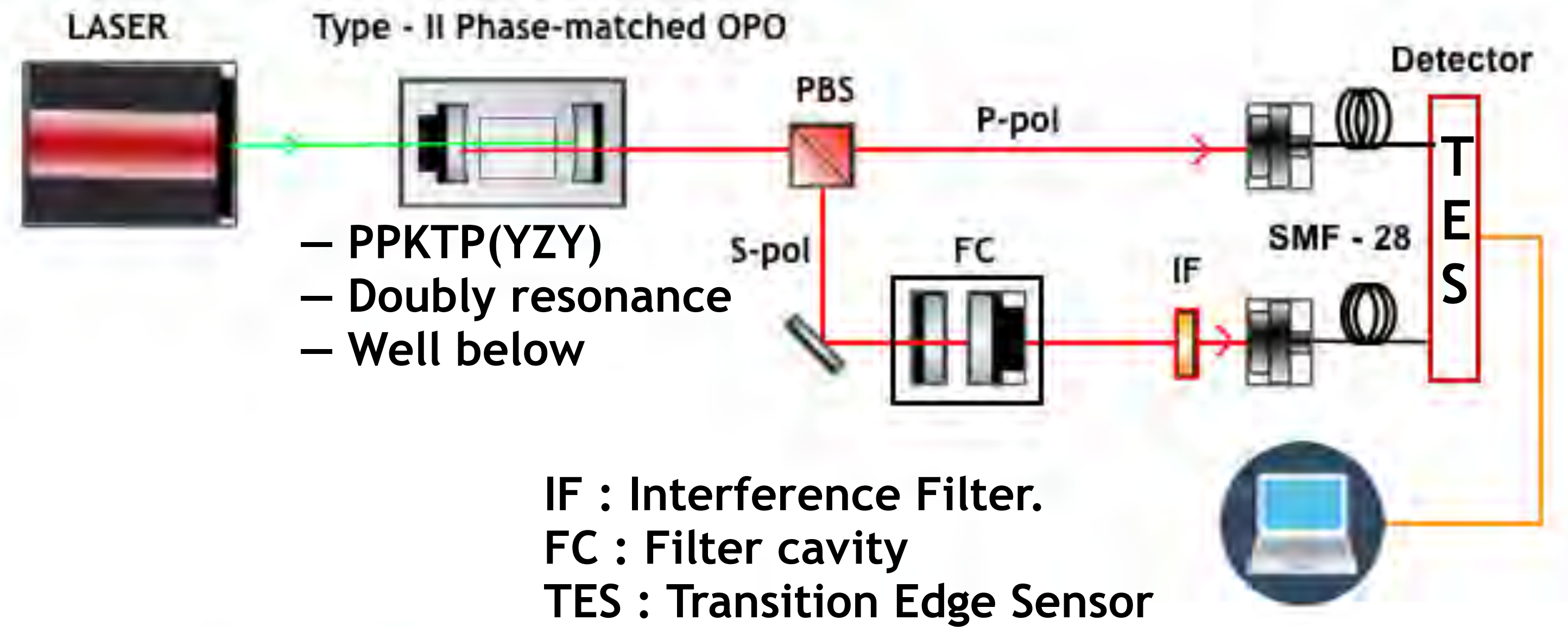
# State-independent quantum state tomography by photon-number-resolving measurements

RAJVEER NEHRA,<sup>1,\*</sup> AYE WIN,<sup>1,5</sup> MILLER EATON,<sup>1</sup> REIHANEH SHAHROKHSAMI,<sup>1,4</sup> NIRANJAN SRIDHAR,<sup>1,3</sup> THOMAS GERRITS,<sup>2</sup> ADRIANA LITA,<sup>2</sup> SAE WOO NAM,<sup>2</sup> AND OLIVIER PFISTER<sup>1</sup>





Nd: YAG  
532 nm,  
FWHM = 1kHz



Heralded(Signal)  
Channel

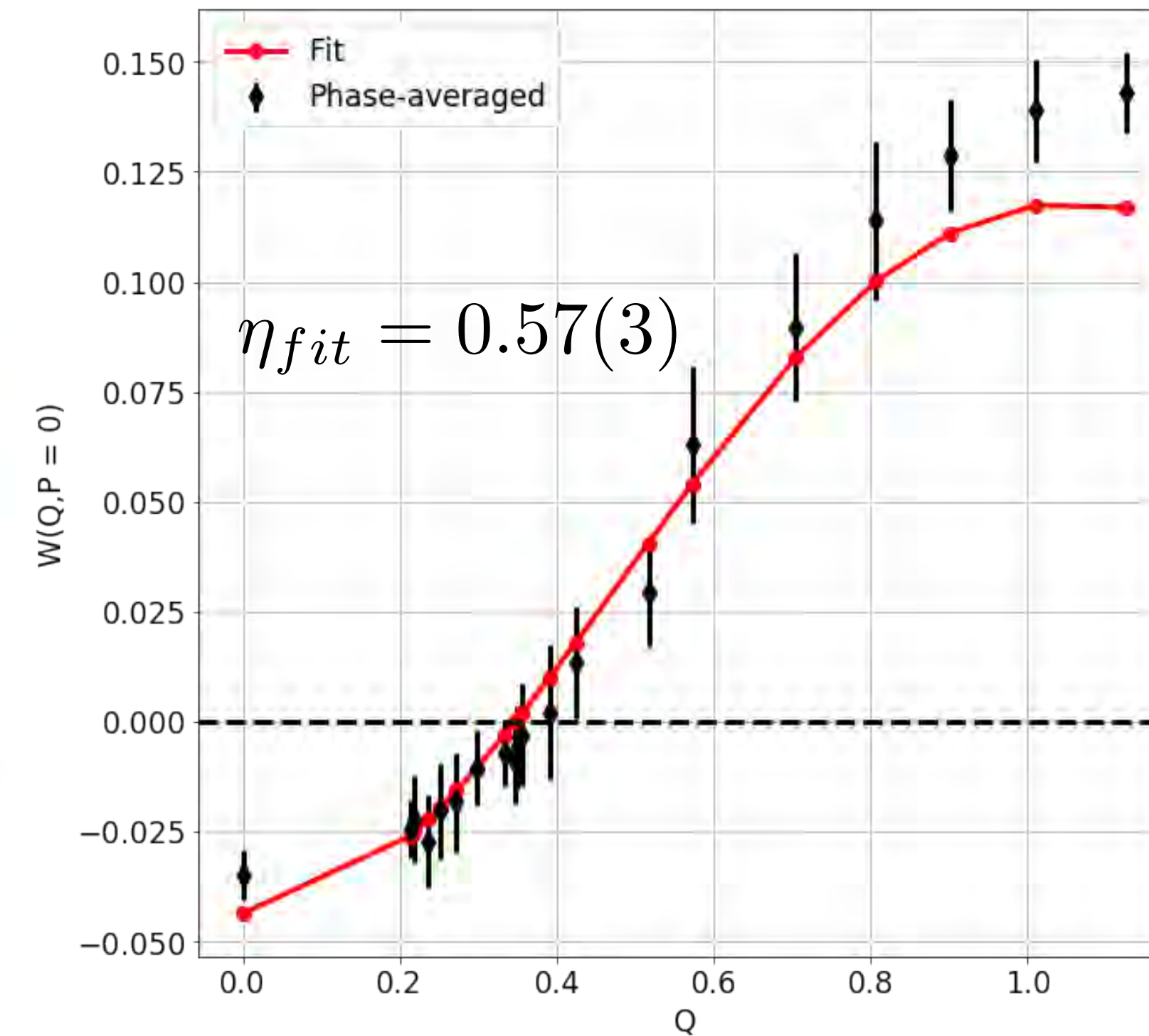
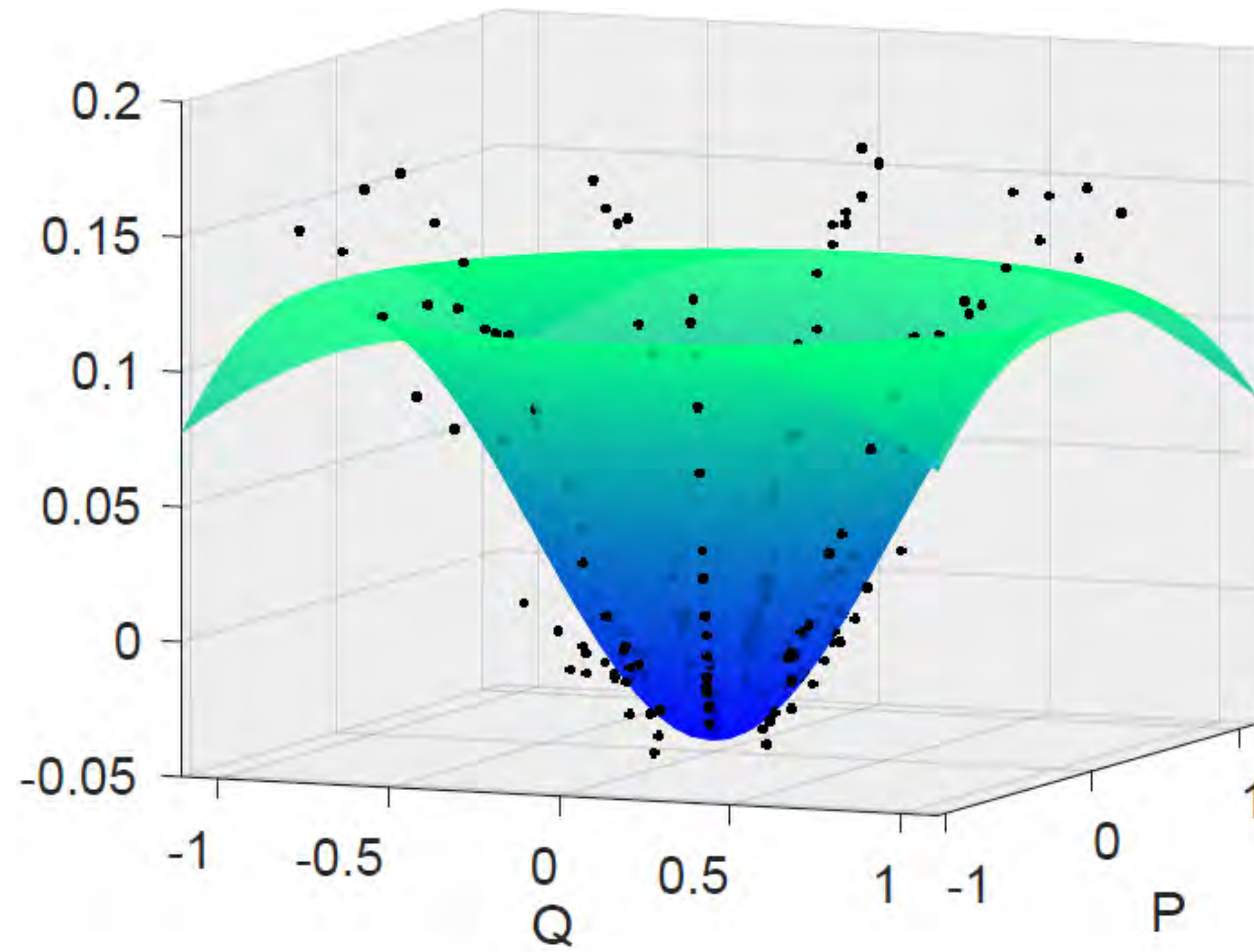
Heralding(idler)  
Channel

$$\eta_h = \frac{N_c}{N_i} = 0.58(2)$$

Overall efficiency  
of the signal path.



# Experimental results

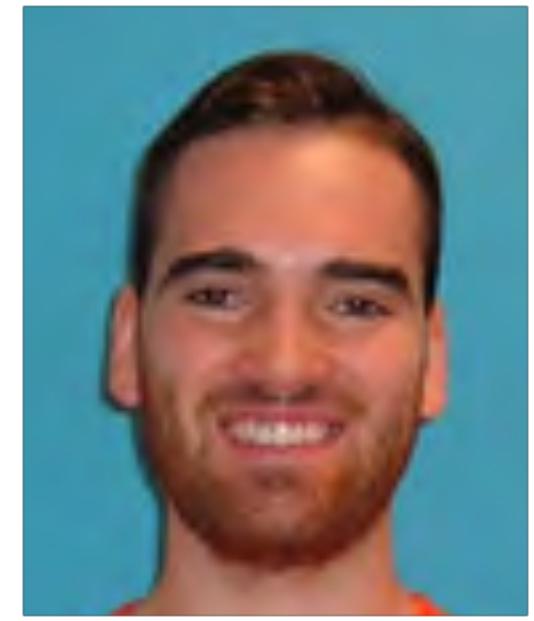


★ Negativity was observed in the raw data without any inference or correcting for losses.



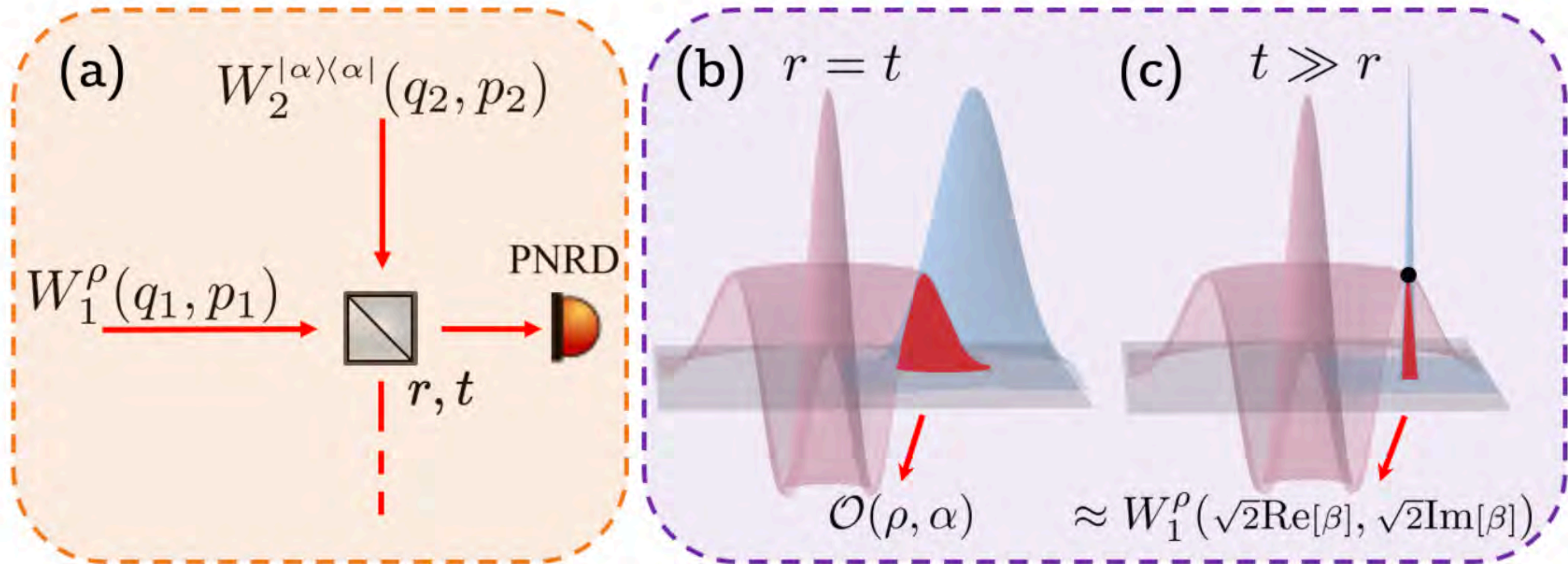


# An even better way to do it



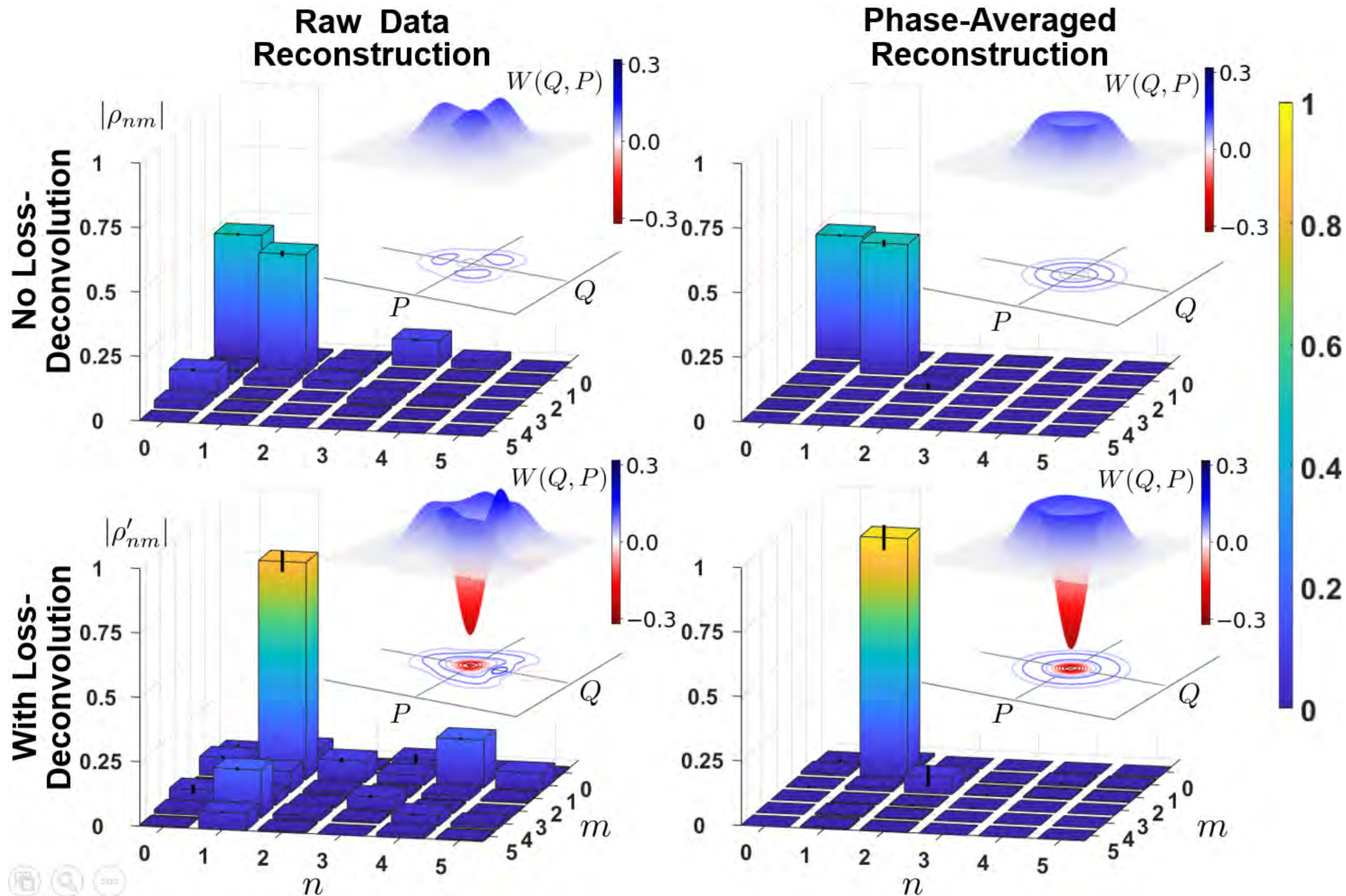
## Generalized overlap quantum state tomography

Rajveer Nehra,<sup>1,\*</sup> Miller Eaton,<sup>1,†</sup> Carlos González-Arciniegas,<sup>1</sup> M. S. Kim,<sup>2,3</sup>  
Thomas Gerrits,<sup>4</sup> Adriana Lita,<sup>4</sup> Sae Woo Nam,<sup>4</sup> and Olivier Pfister<sup>1,‡</sup>



- **Efficient reconstruction** using semidefinite programming





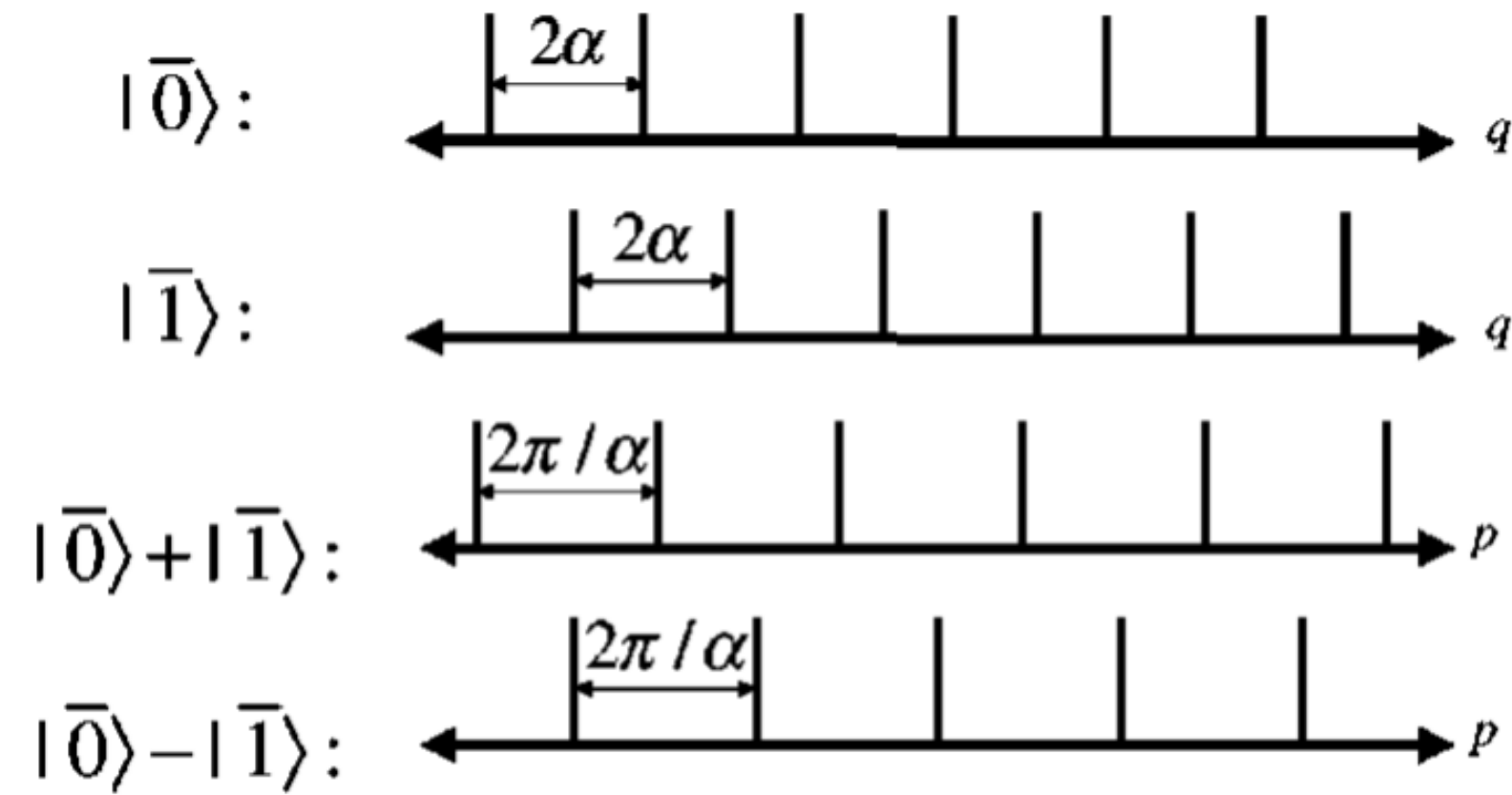


# Non-Gaussian resources => quantum error correction

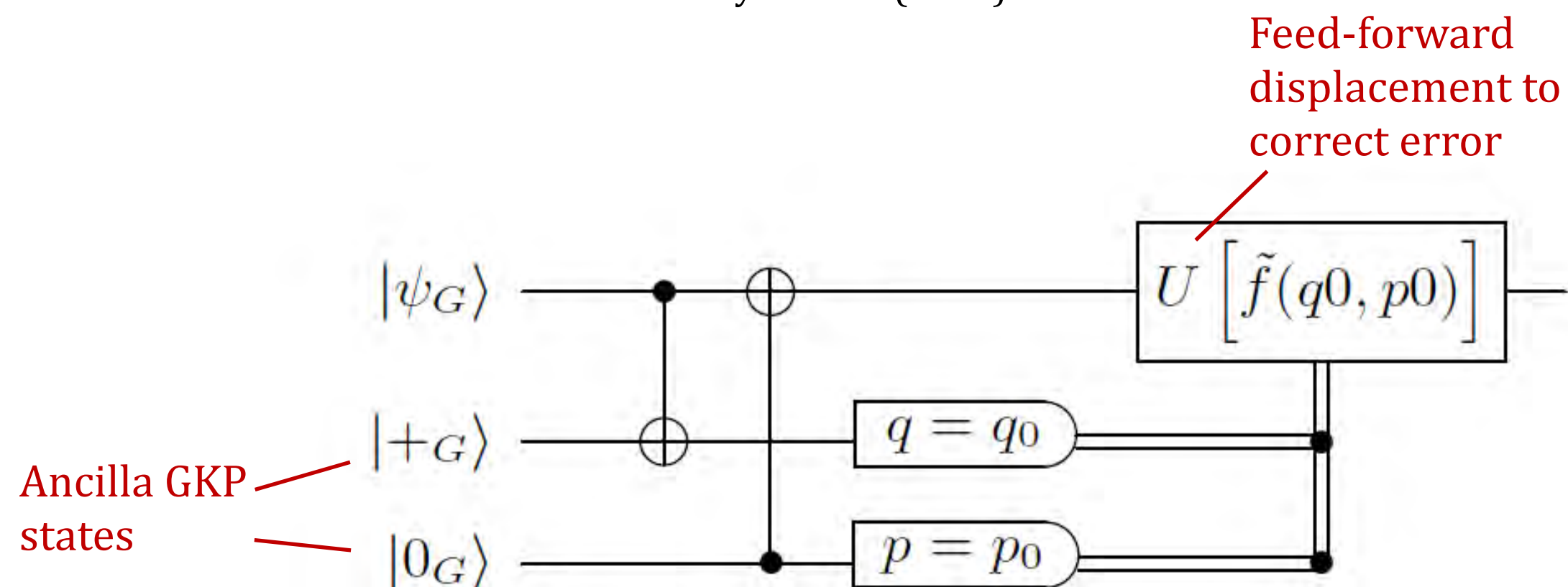
## Gottesman-Kitaev-Preskill (GKP) States

Simultaneous eigenstates of

$$X = e^{-i\hat{P}\alpha}, \quad Z = e^{\frac{i\pi}{\alpha}\hat{Q}}$$



Gottesman et al. Phys. Rev. A (2001)



Tzitrin et al. arXiv preprint arXiv:1910.03673 (2019).

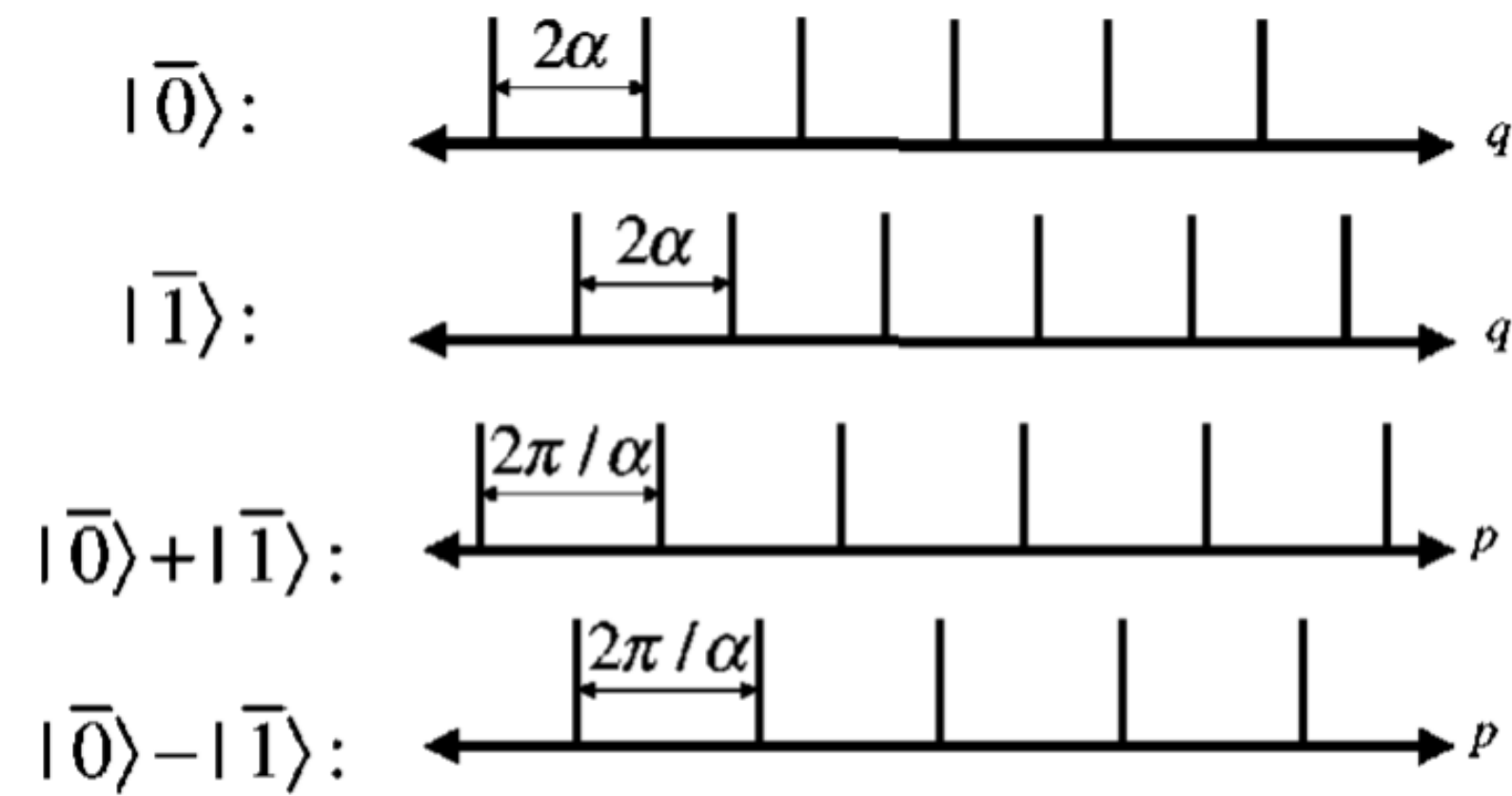


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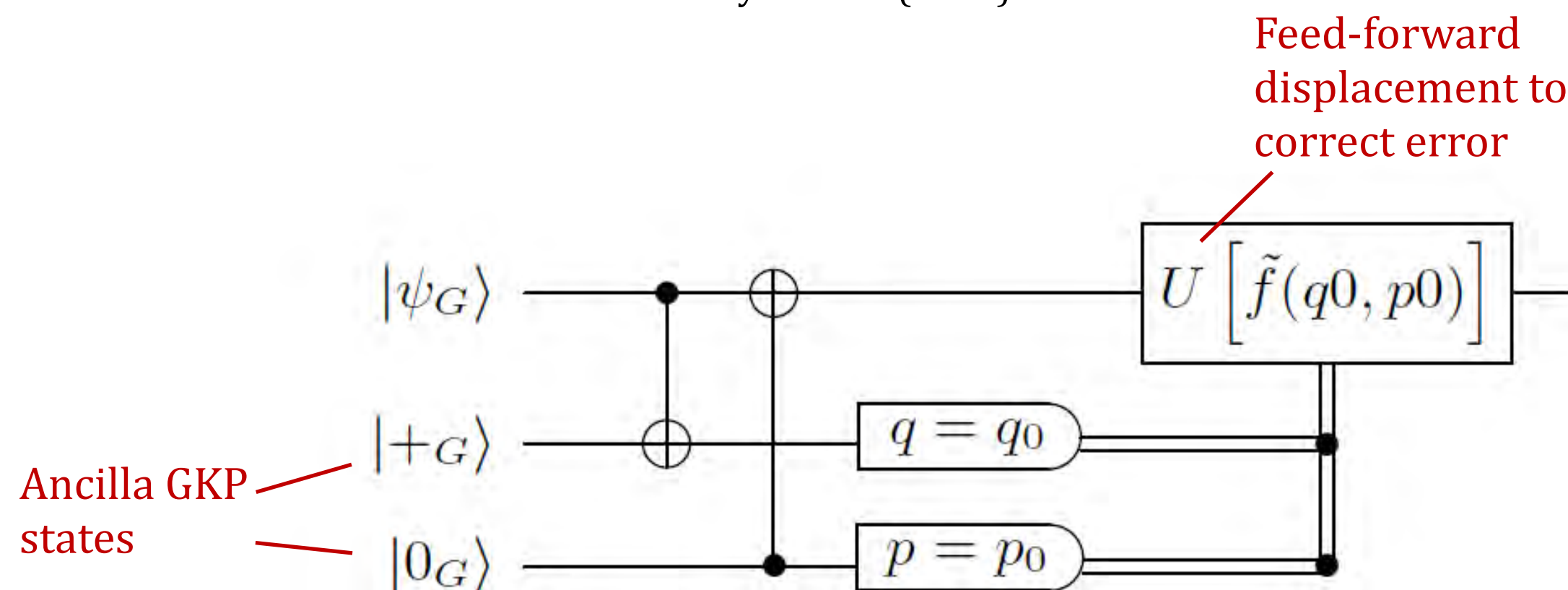
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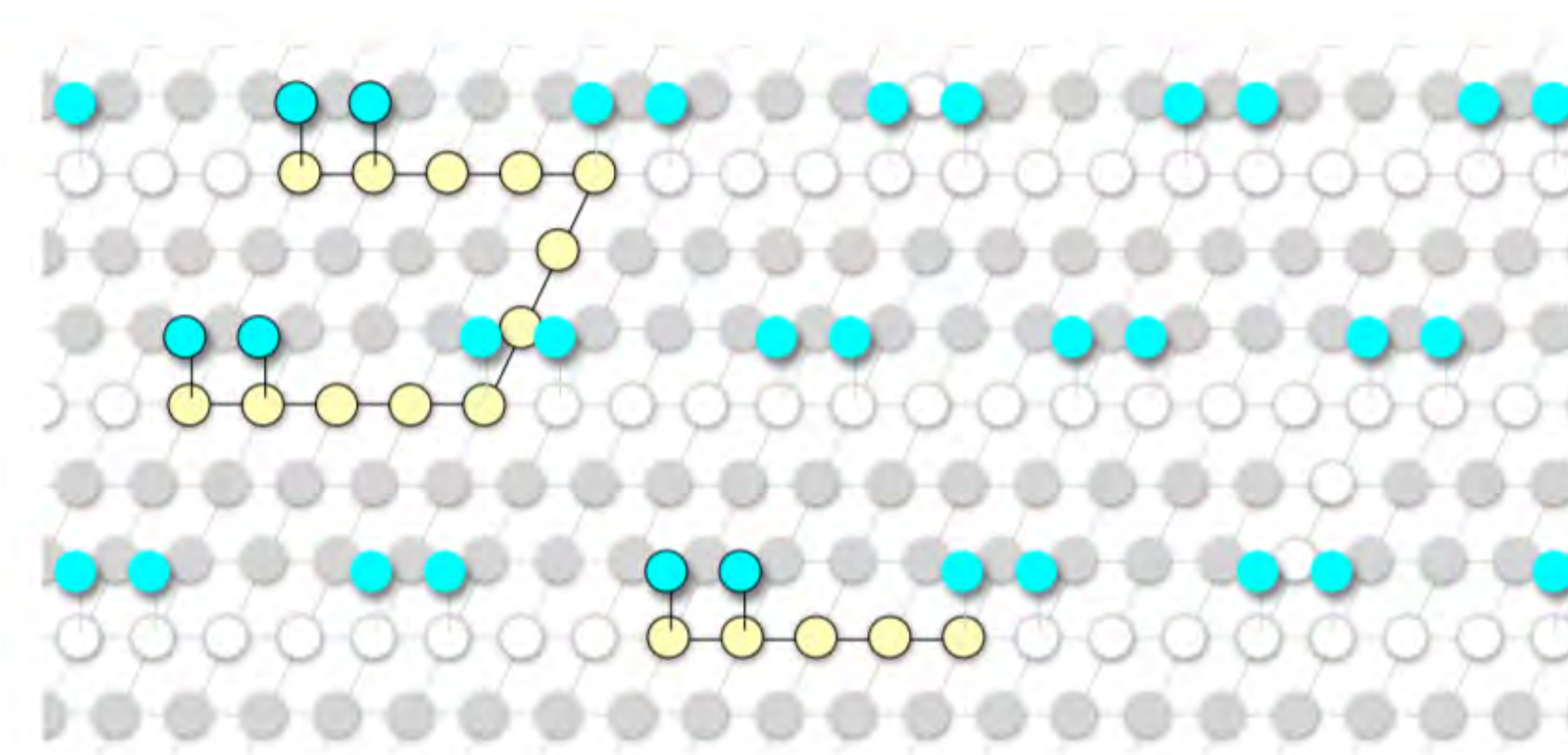
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Tzitrin et al. arXiv preprint arXiv:1910.03673 (2019).



- Cluster state nodes
- GKP states
- Error corrected

Walshe, et al. *Physical Review A* (2019).

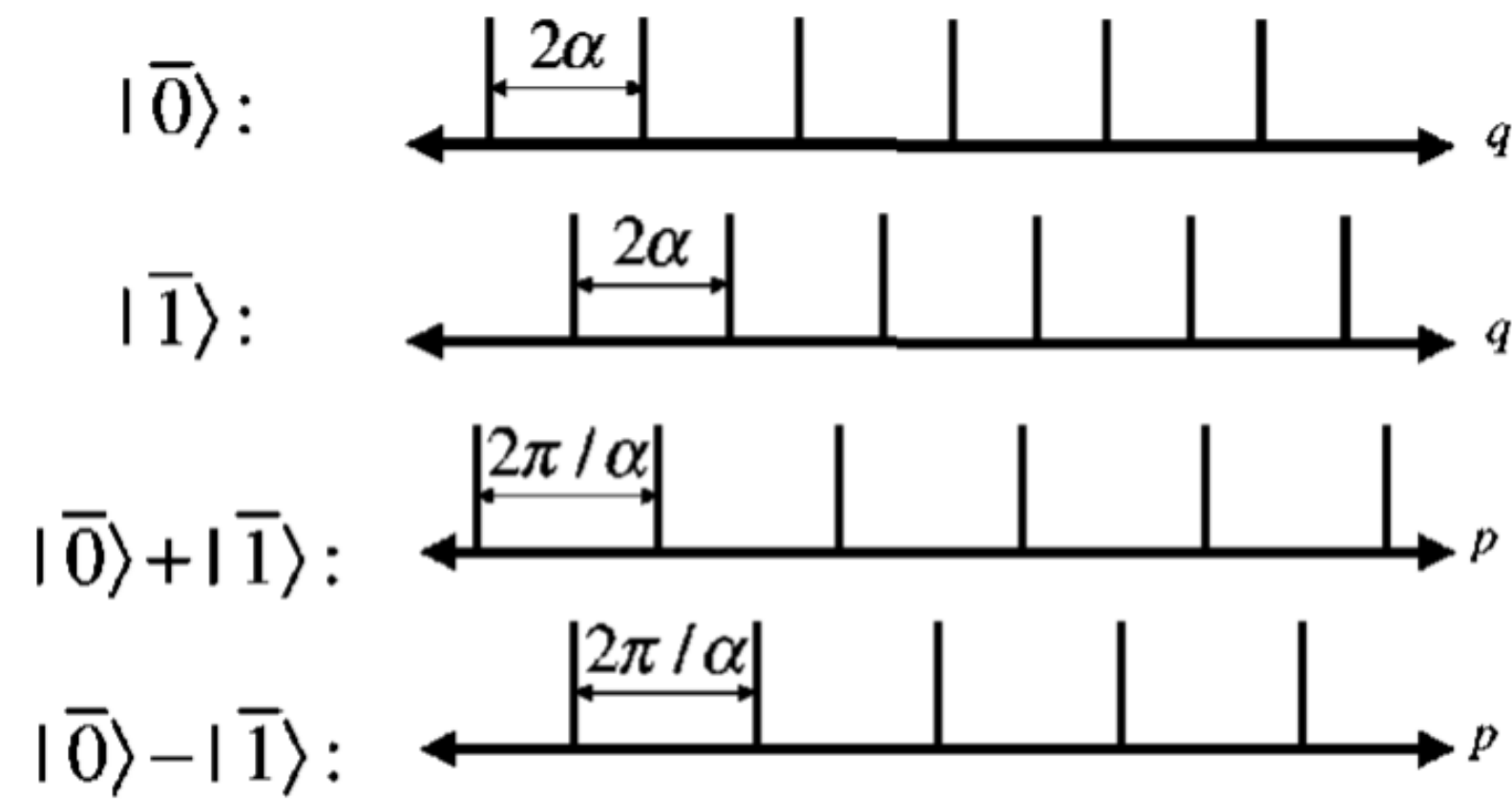


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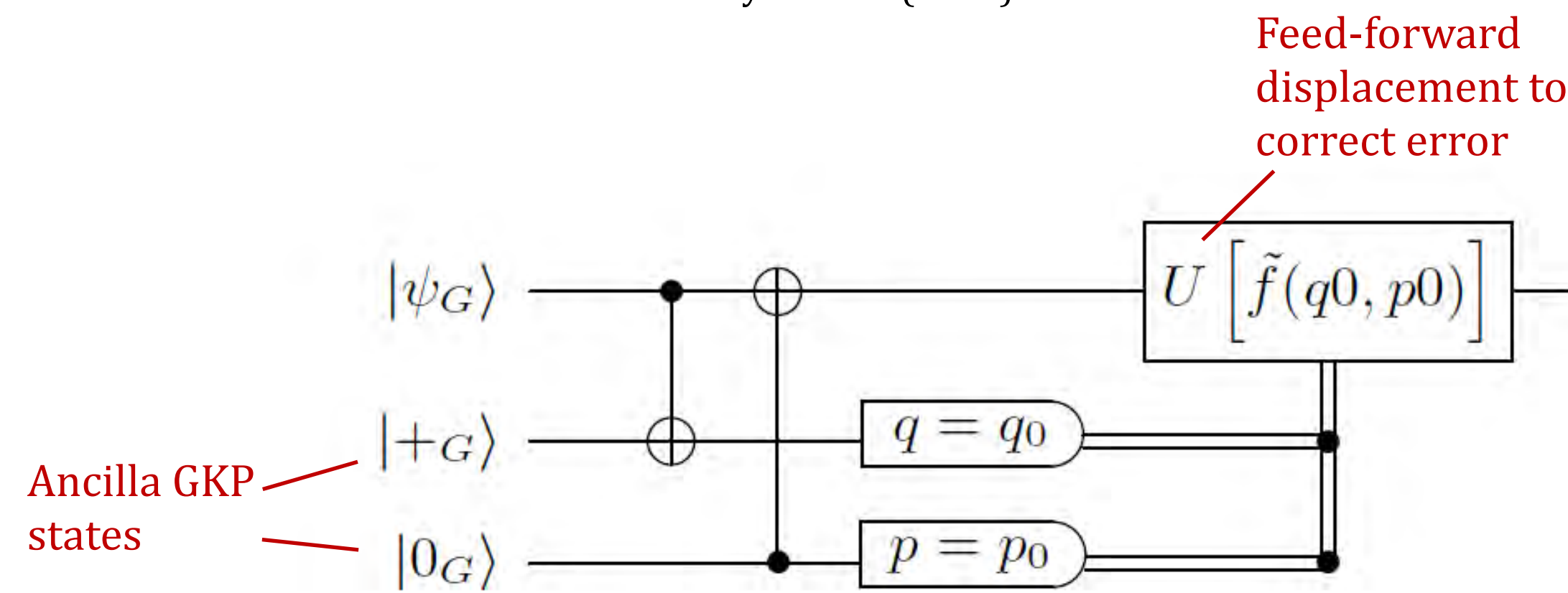
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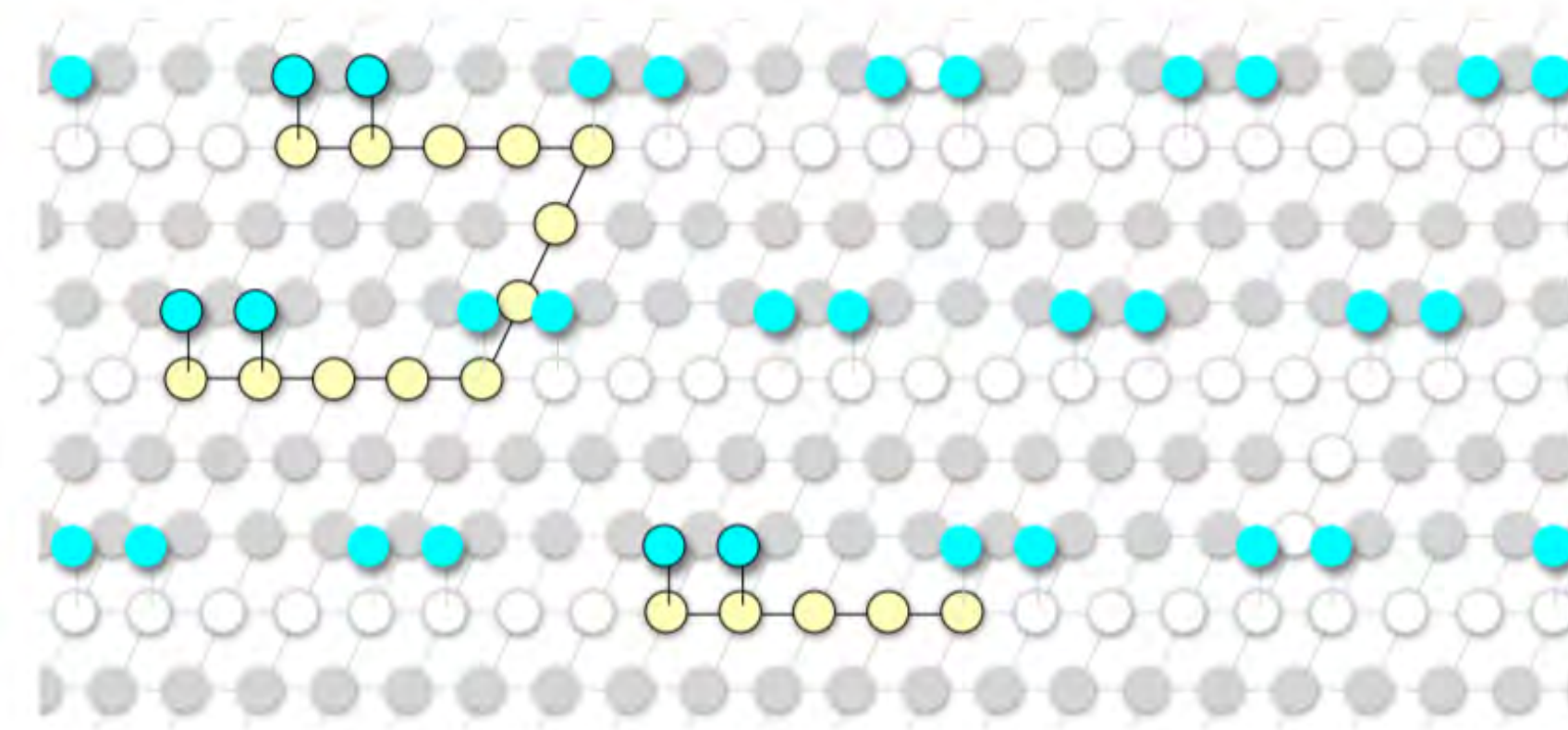
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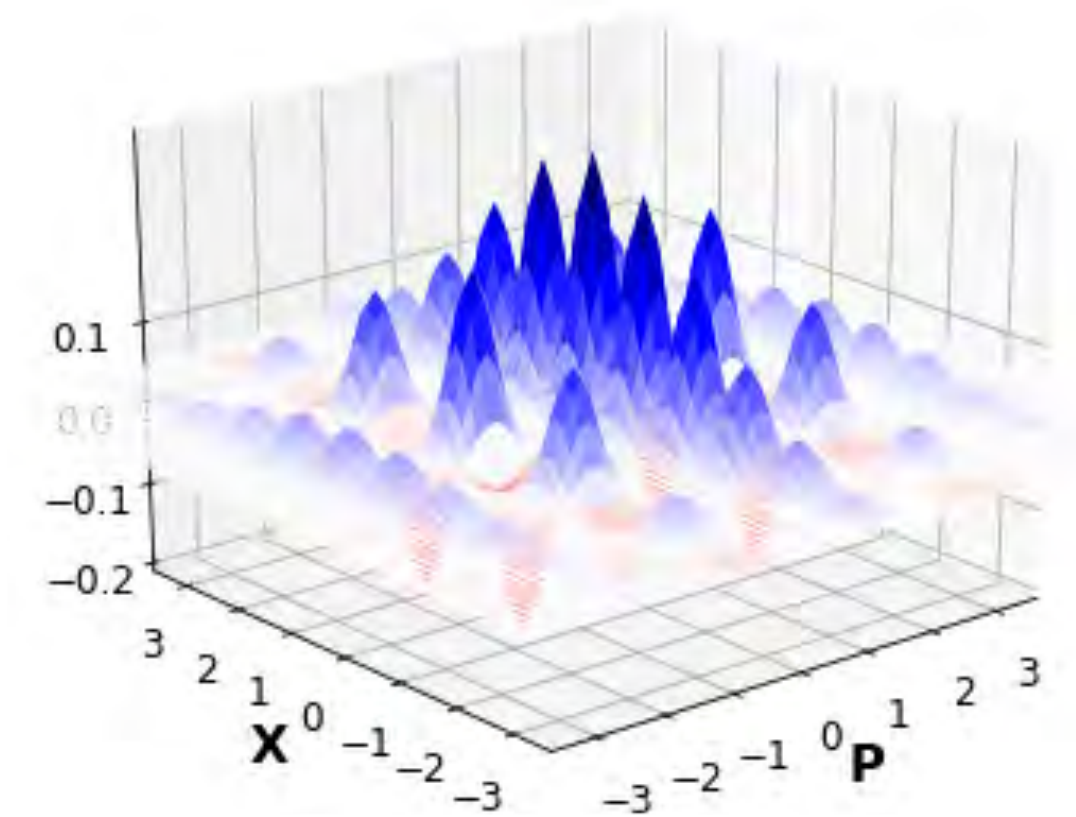
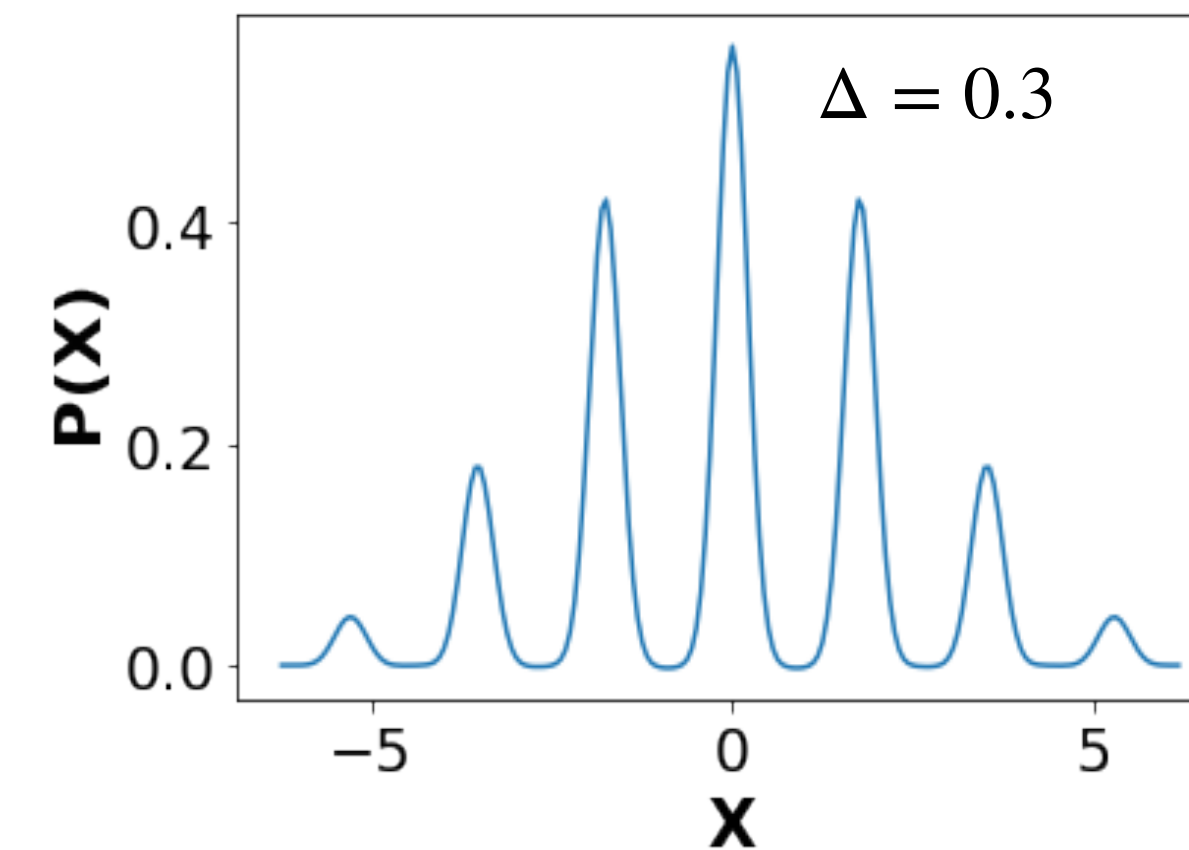
Cluster state  
 nodes

GKP states

Error corrected

Walshe, et al. *Physical Review A* (2019).

Realistic implementation: Delta spikes  $\rightarrow$  peak width  $\Delta$

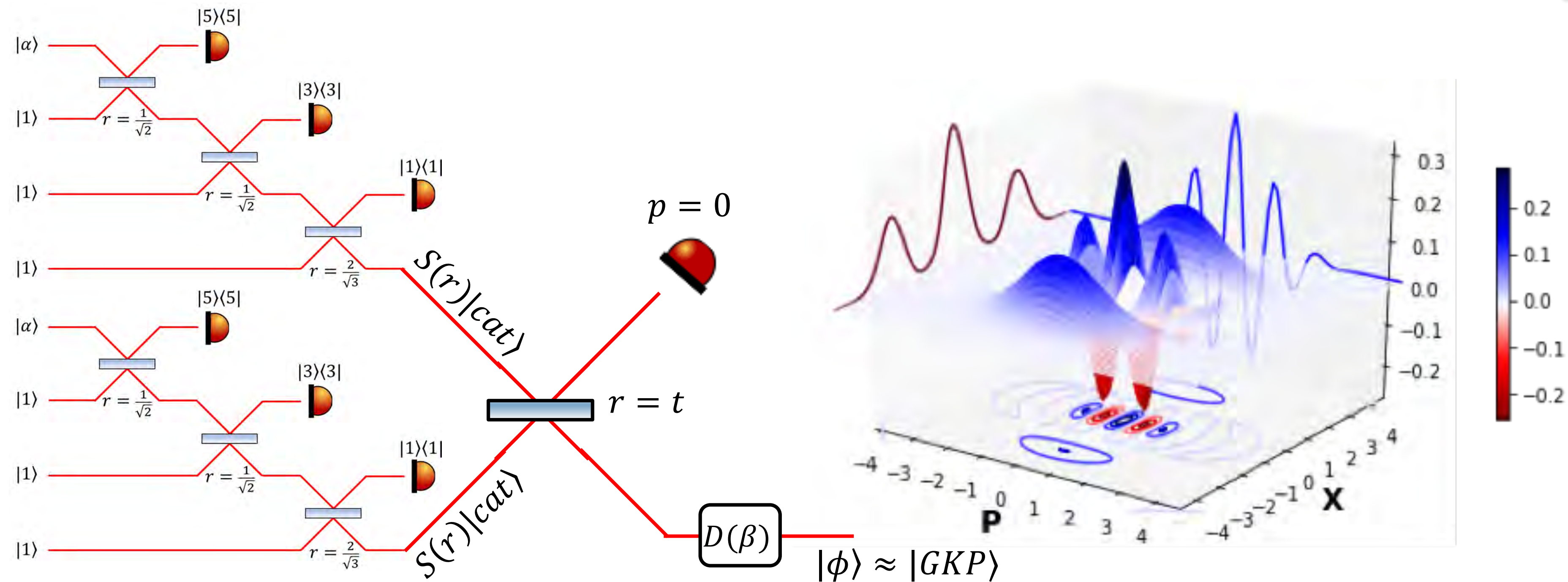


$$|\bar{0}\rangle \propto e^{\Delta^2 \hat{N}} \sum_{n_1, n_2 = -\infty}^{\infty} e^{-in_1 \hat{P} \sqrt{\pi}} e^{in_2 \hat{Q} \sqrt{\pi}} |vac\rangle$$





# Quantum engineering of GKP states for CV quantum error correction



H. M. Vasconcelos, L. Sanz, and S. Glancy, "All-optical generation of states for 'Encoding a qubit in an oscillator,'" *Opt. Lett.*, *OL*, vol. 35, no. 19, pp. 3261–3263, Oct. 2010.

M. Eaton, R. Nehra, and O. Pfister, **Non-Gaussian and Gottesman-Kitaev-Preskill state preparation by photon catalysis**, *New Journal of Physics* **21**, 113034 (2019).



# BIG PICTURE: beating lattice-gauge QCD calculations on classical supercomputers



## Quantum Algorithms for Quantum Field Theories

Stephen P. Jordan,<sup>1\*</sup> Keith S. M. Lee,<sup>2</sup> John Preskill<sup>3</sup>  
1 JUNE 2012 VOL 336 SCIENCE www.sciencemag.org

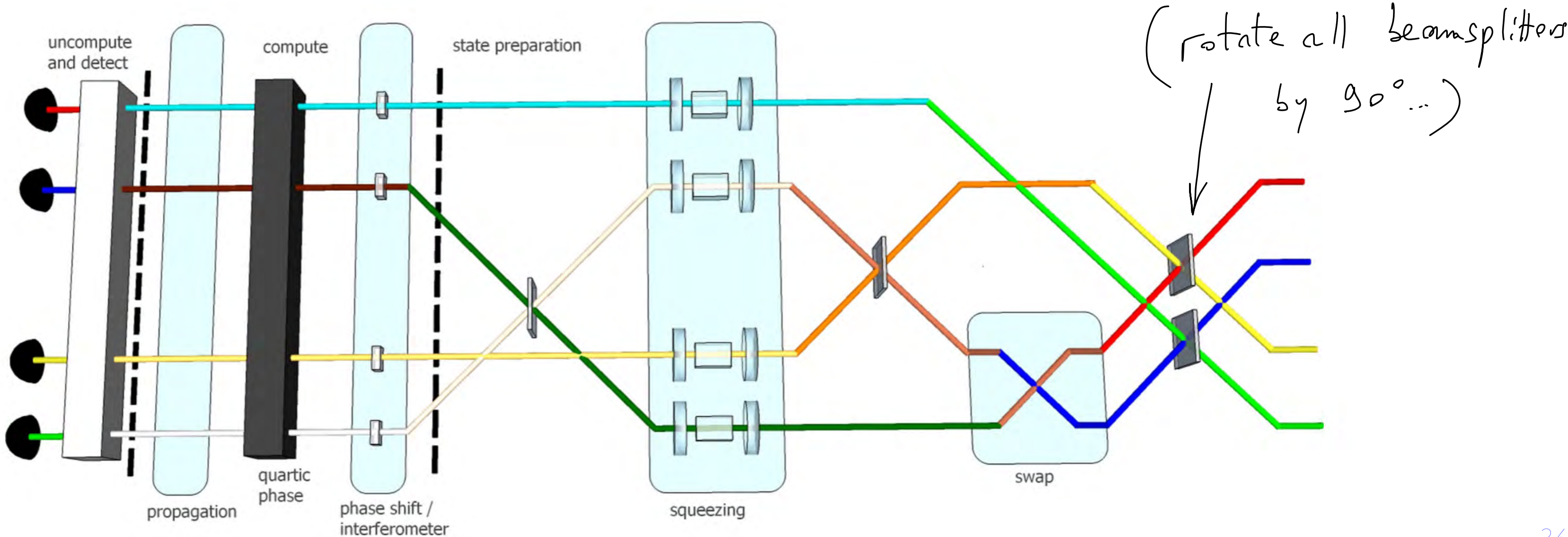
$$\mathcal{A} = \langle \text{out} | T \exp \left\{ i \int_{-T}^T dt [H_{\text{int}}(t) + H_{\text{c.t.}}(t)] \right\} | \text{in} \rangle$$

$\phi^4$  term

PHYSICAL REVIEW A 92, 063825 (2015)

## Quantum simulation of quantum field theory using continuous variables

Kevin Marshall,<sup>1</sup> Raphael Pooser,<sup>2,3</sup> George Siopsis,<sup>3,\*</sup> and Christian Weedbrook<sup>4</sup>





# Conclusion

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**The quantum optical frequency comb is a viable platform for universal quantum computing**



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- **LARGE-SCALE** Gaussian entanglement
  - No postselection
  - Universal measurement-based QC, strictly equivalent to qubit model

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  - Universal measurement-based QC, strictly equivalent to qubit model
- **FAULT TOLERANT** non-Gaussian technology for universal QC
- **INTEGRABLE** in photonic circuits
- **TRANSLATABLE** to any bosonic mode:
  - microwave cavity photons
  - phonons
  - transverse spatial modes (Hermite- or Laguerre-Gauss)
  - temporal modes



# A lesson in research from Ted Hänsch...

